

***The effect of
cross-sectional geometry and size on
magnetostatic modes in nanorods***

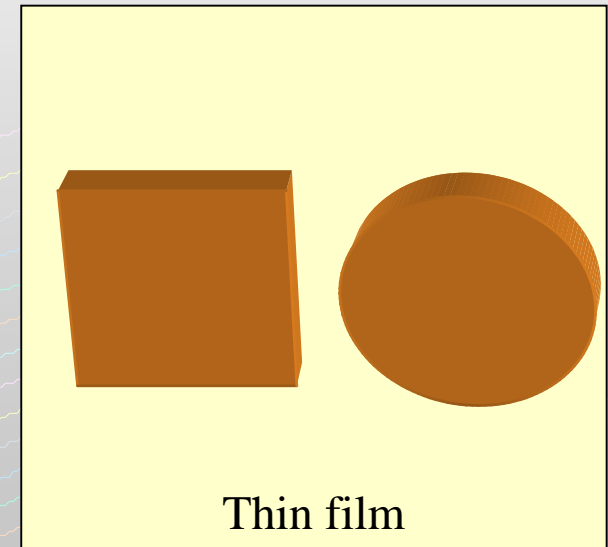
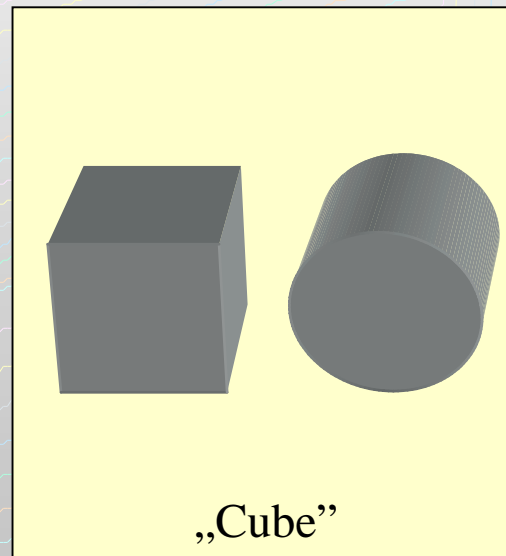
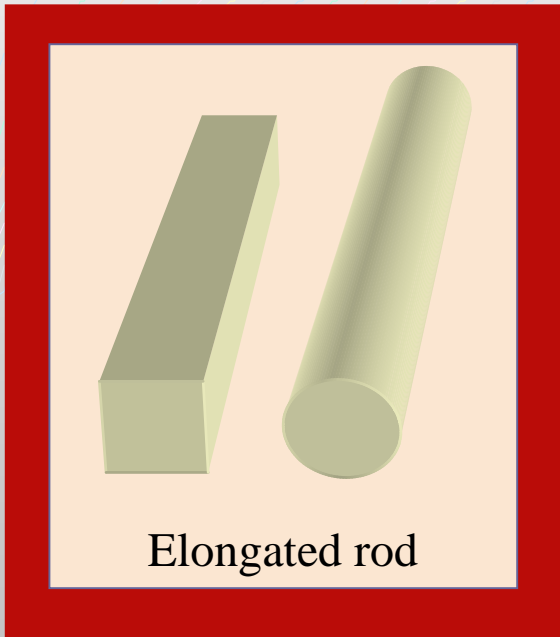
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OUTLINE

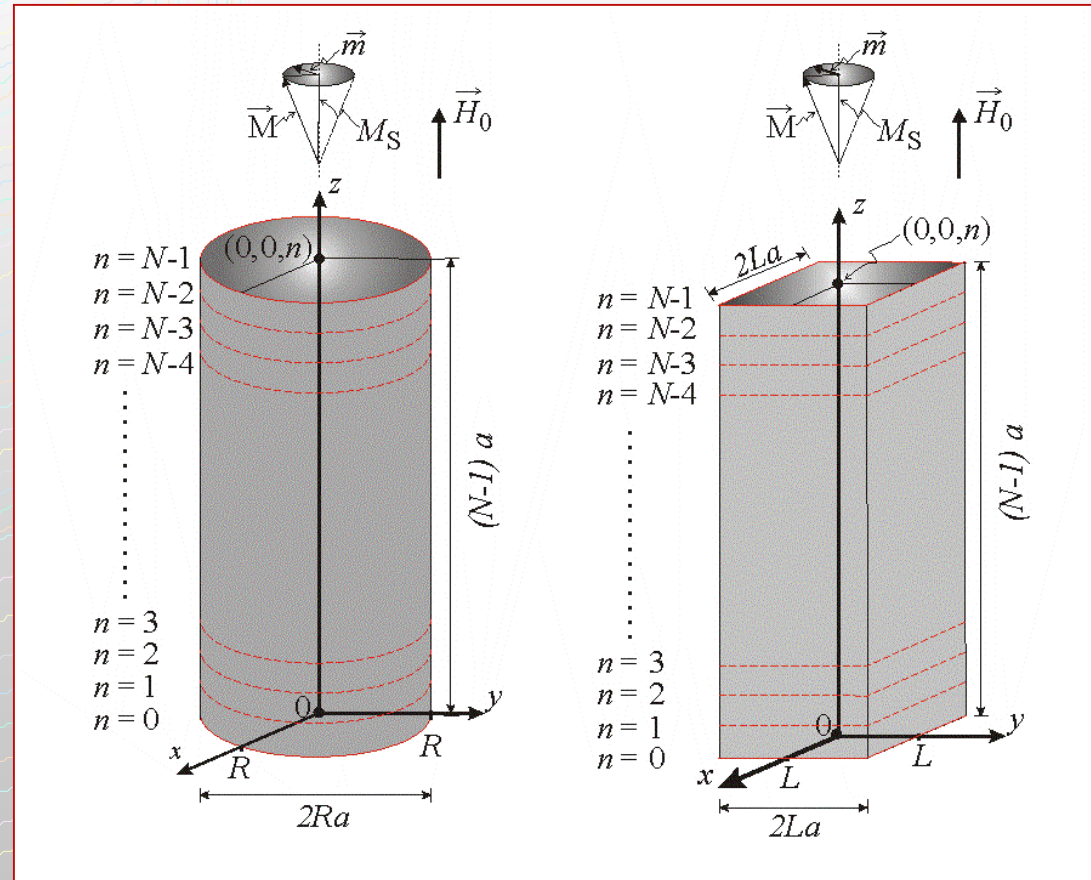
- ❑ The model of calculations
- ❑ Magnetostatic modes in nanorods
- ❑ The effect of cross-sectional geometry and size on
 - A) magnetostatic modes
 - B) effective dipolar field



Model assumptions:

- magnetic dipoles in planar arrangements,
- uniform magnetization,
- propagation along external magnetic field,
- uniform precession in each plane.

Central axis approximation.



Equations:

$$\frac{\partial \vec{M}_{n'}}{\partial t} = \gamma \mu_0 \vec{M}_{n'} \times \left(\vec{H}_0 + \vec{h}_{n'} \right)$$

Dipolar field in the central axis approximation:

$$\vec{h}_{n'} = \frac{1}{4\pi} \sum_n D_{n,n'} \left[\hat{i} M_n^x + \hat{j} M_n^y - 2\hat{k} M_n^z \right]$$

is obtained from a standard formulae:

$$\vec{h}_{\vec{\rho}} = \frac{1}{4\pi} \sum_{\vec{r} \neq \vec{\rho}} \frac{3(\vec{r} - \vec{\rho})(\vec{\mu}_{\vec{r}} \cdot (\vec{r} - \vec{\rho})) - \vec{\mu}_{\vec{r}} |\vec{r} - \vec{\rho}|^2}{|\vec{r} - \vec{\rho}|^5}$$

We define **dipolar matrix**:

$$\delta = n - n'$$

$$D_\delta \equiv \sum_{p,q} \frac{\frac{1}{2}(p^2 + q^2) - \delta^2}{[p^2 + q^2 + \delta^2]^{\frac{5}{2}}}$$

$$\hat{D} = \begin{pmatrix} D_0 & D_1 & D_2 & \cdots & D_{N-2} & D_{N-1} \\ D_1 & D_0 & D_1 & \cdots & D_{N-3} & D_{N-2} \\ D_2 & D_1 & D_0 & \cdots & D_{N-4} & D_{N-3} \\ & & & \ddots & & \\ & & & & \ddots & \\ & & & & & \ddots \\ D_{N-2} & D_{N-3} & D_{N-4} & \cdots & D_0 & D_1 \\ D_{N-1} & D_{N-2} & D_{N-3} & \cdots & D_1 & D_0 \end{pmatrix}$$

In linear approximation we obtain eigenvalue problem

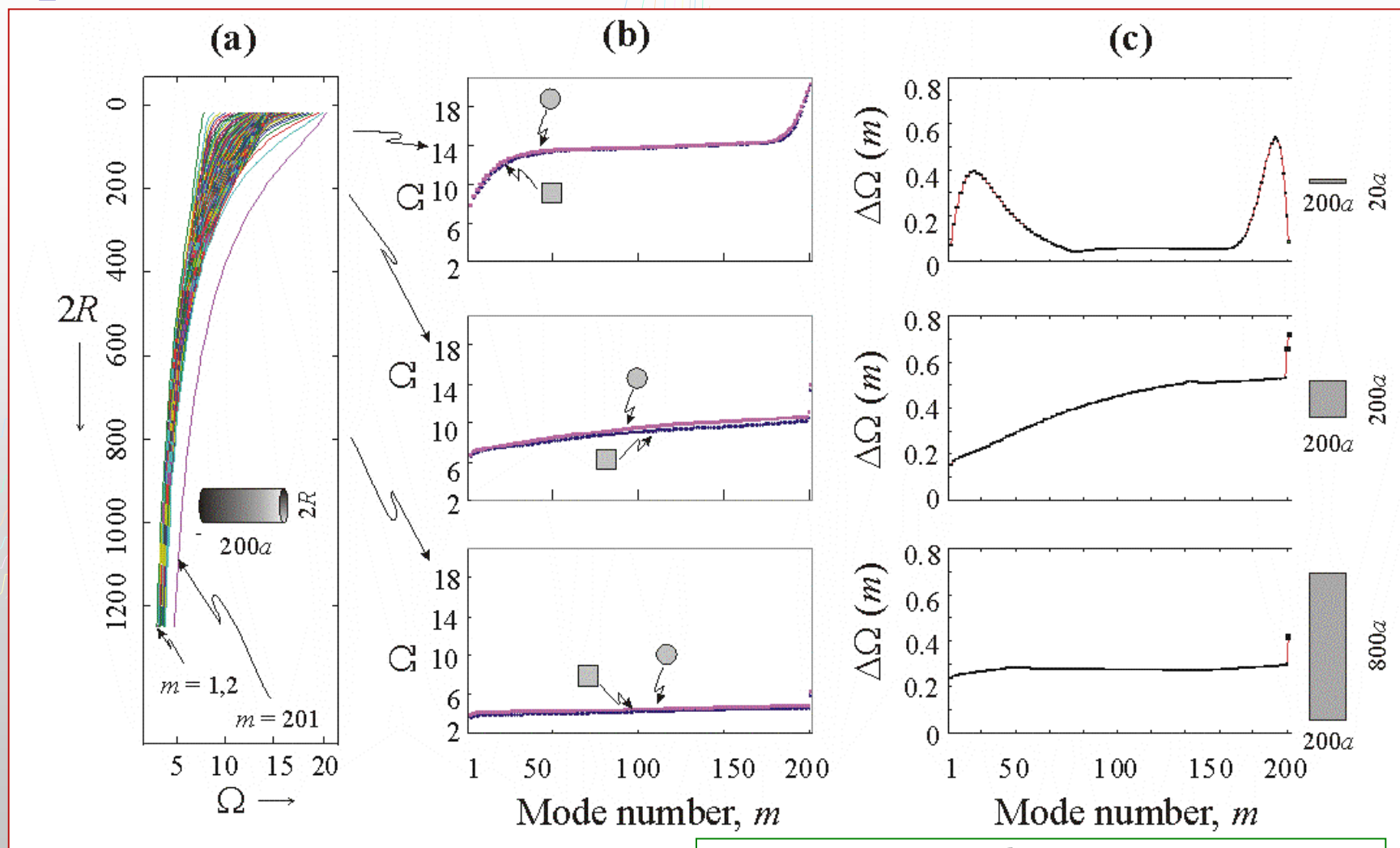
$$\hat{H} \vec{m}^+ = \Omega \vec{m}^+$$

$$\begin{pmatrix} H(0) & D_1 & D_2 & \cdots & D_{N-2} & D_{N-1} \\ D_1 & H(1) & D_1 & \cdots & D_{N-3} & D_{N-2} \\ D_2 & D_1 & H(2) & \cdots & D_{N-4} & D_{N-3} \\ & & & \cdots & & \\ & & & & \cdots & \\ D_{N-2} & D_{N-3} & D_{N-4} & \cdots & H(N-2) & D_1 \\ D_{N-1} & D_{N-2} & D_{N-3} & \cdots & D_1 & H(N-1) \end{pmatrix} \begin{pmatrix} m_0^+ \\ m_1^+ \\ m_2^+ \\ \cdot \\ \cdot \\ m_{N-2}^+ \\ m_{N-1}^+ \end{pmatrix} = \frac{4\pi\omega}{\gamma\mu_0 M_S} \begin{pmatrix} m_0^+ \\ m_1^+ \\ m_2^+ \\ \cdot \\ \cdot \\ m_{N-2}^+ \\ m_{N-1}^+ \end{pmatrix}$$

Local dipolar field is defined as:

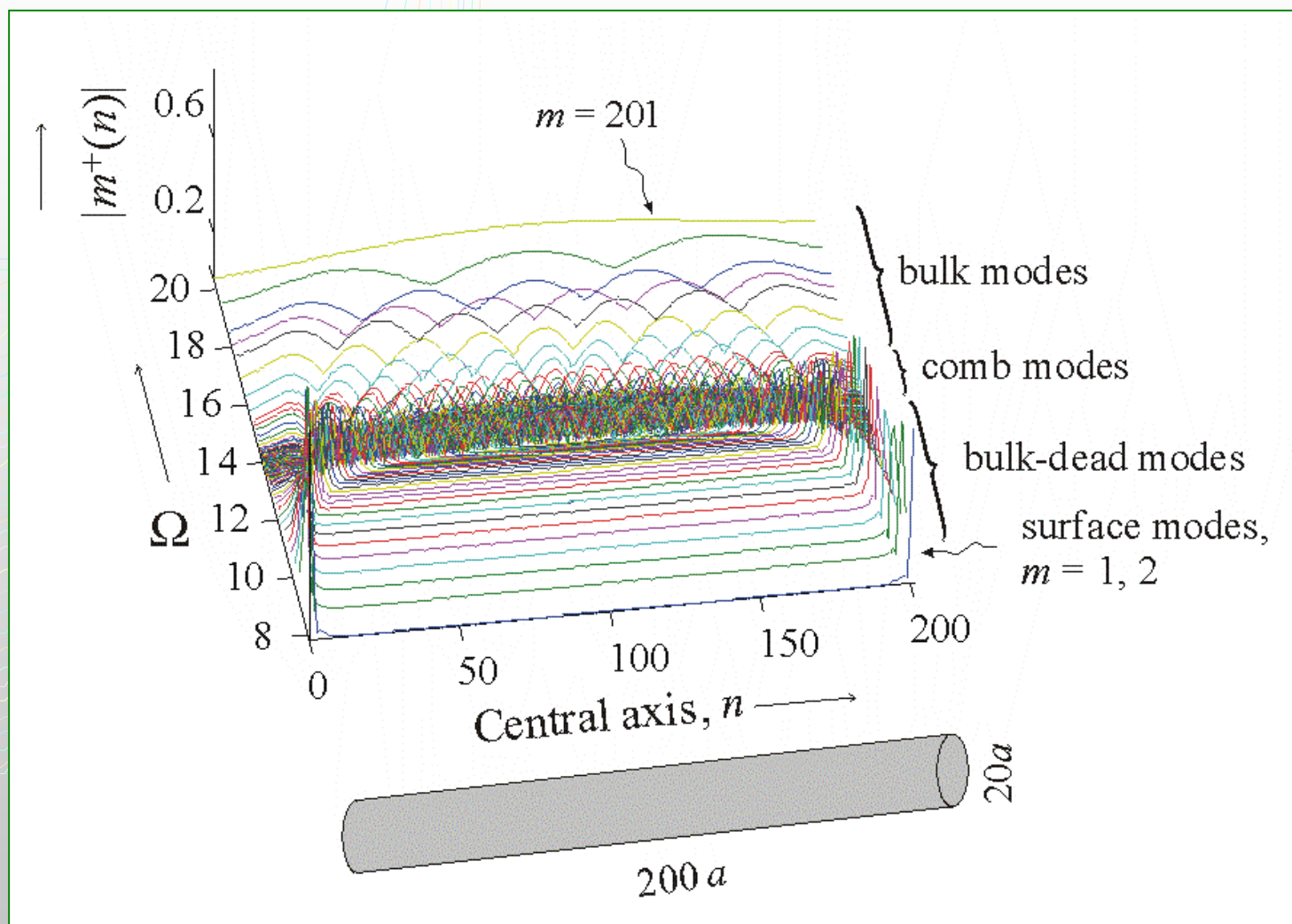
$$H(n') \equiv \Omega_H - D_0 - 2 \sum_{\delta=0,\pm 1,\dots}^{n'} D_\delta$$

Magnetostatic modes dependence on cross-sectional shape and size

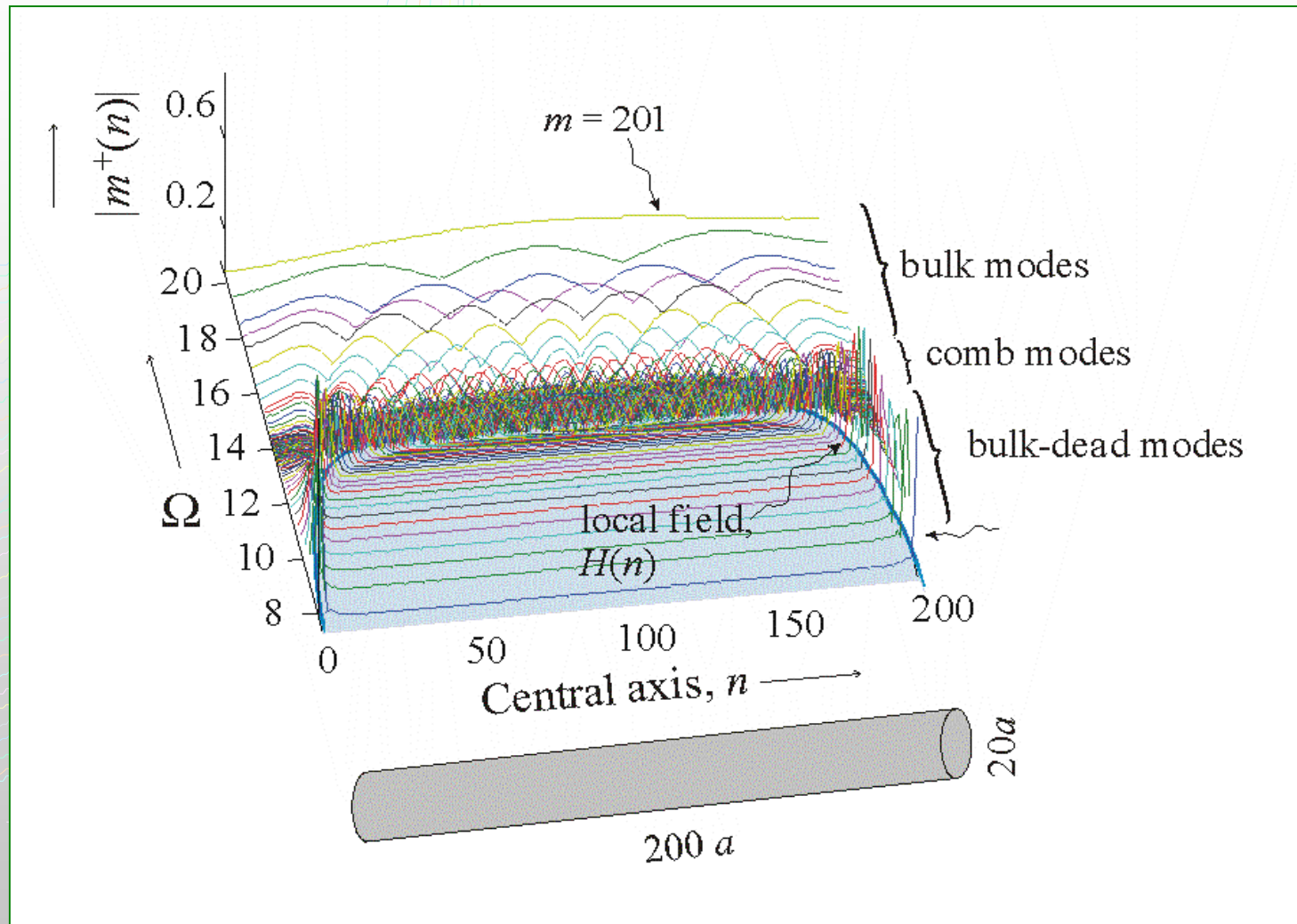


$$\Delta\Omega(m) = \Omega^{\text{circle}}(m) - \Omega^{\text{square}}(m)$$

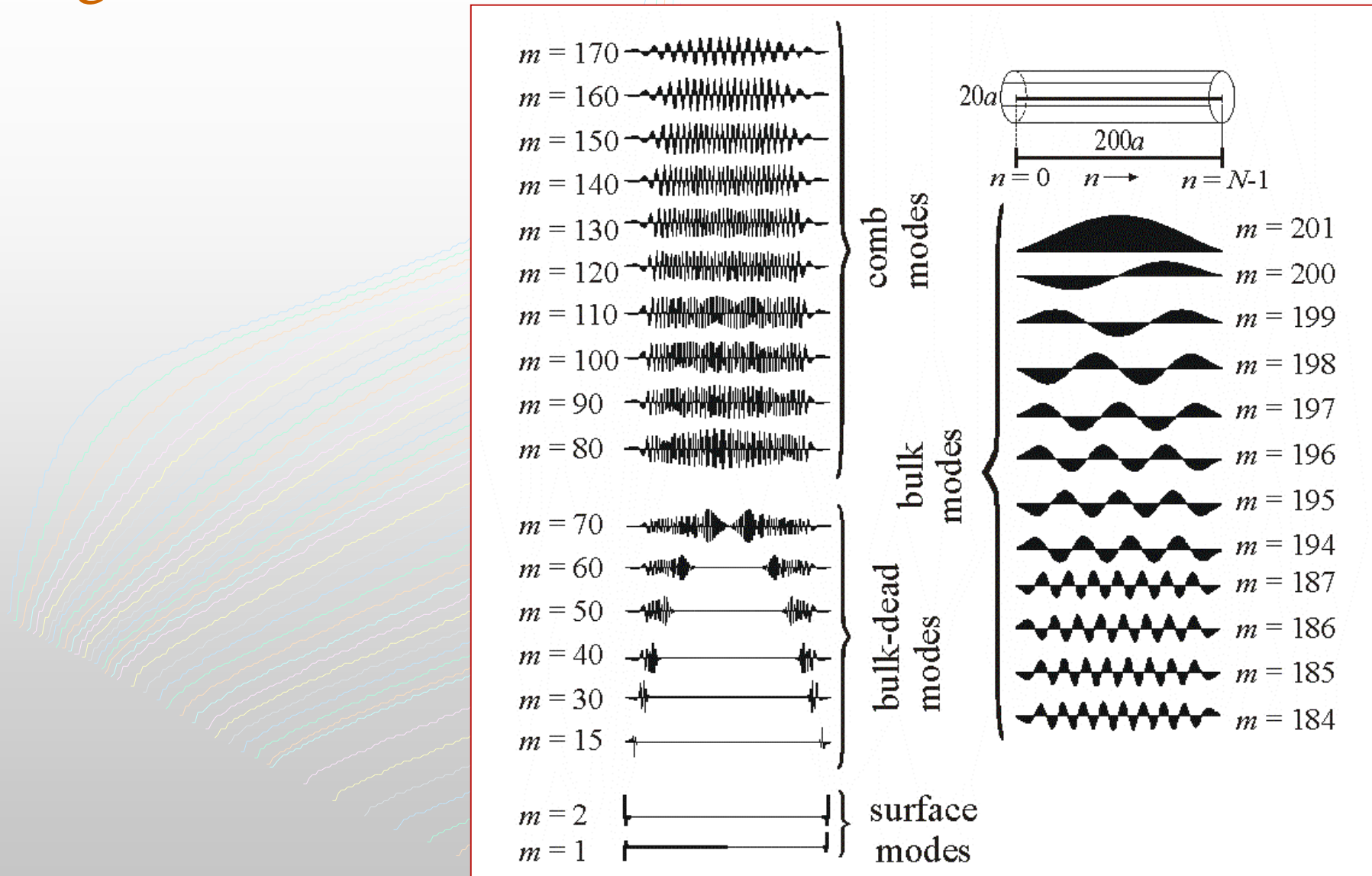
Magnetostatic modes



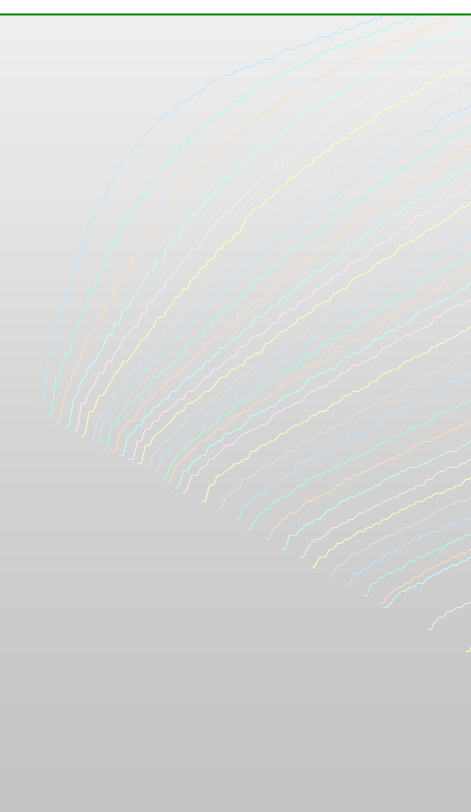
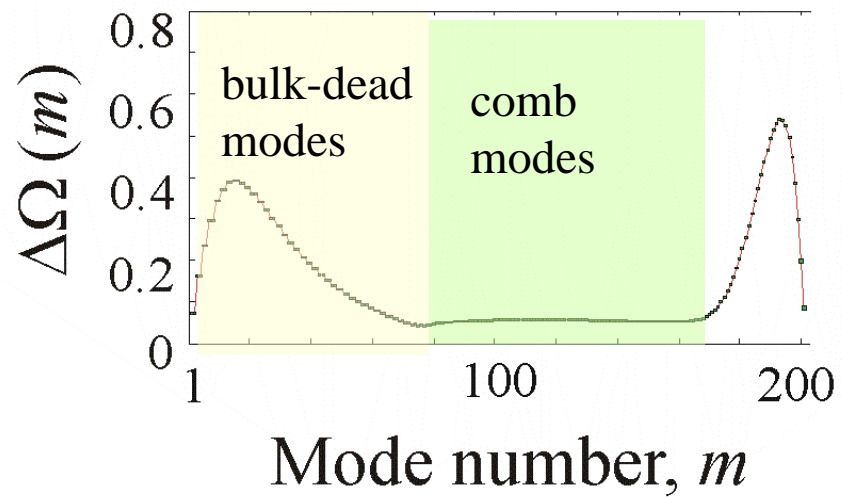
Magnetostatic modes vs. local dipolar field $H(n)$ in nanorods



Magnetostatic modes in nanorods

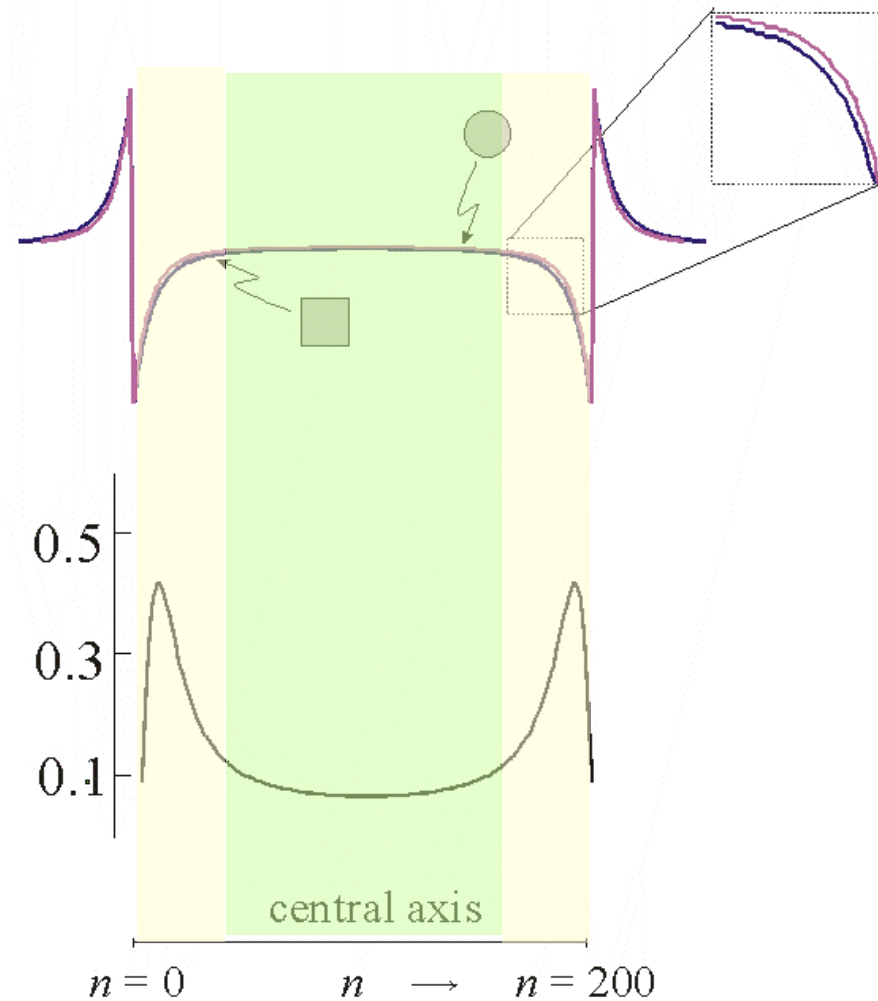


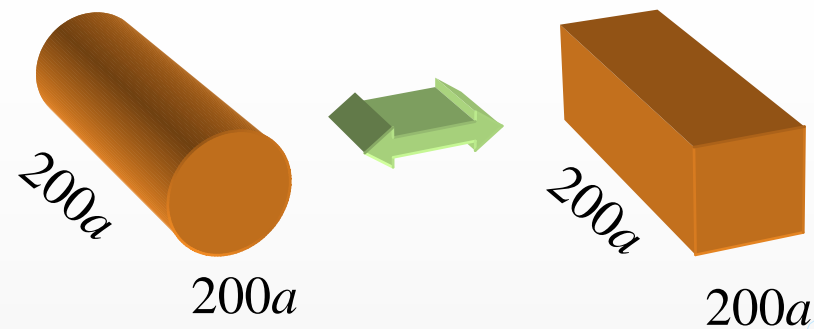
H. Puszkarski, M. Krawczyk, and J.C.S. Lévy, J. Appl. Phys. **101**, 024326 (2007).



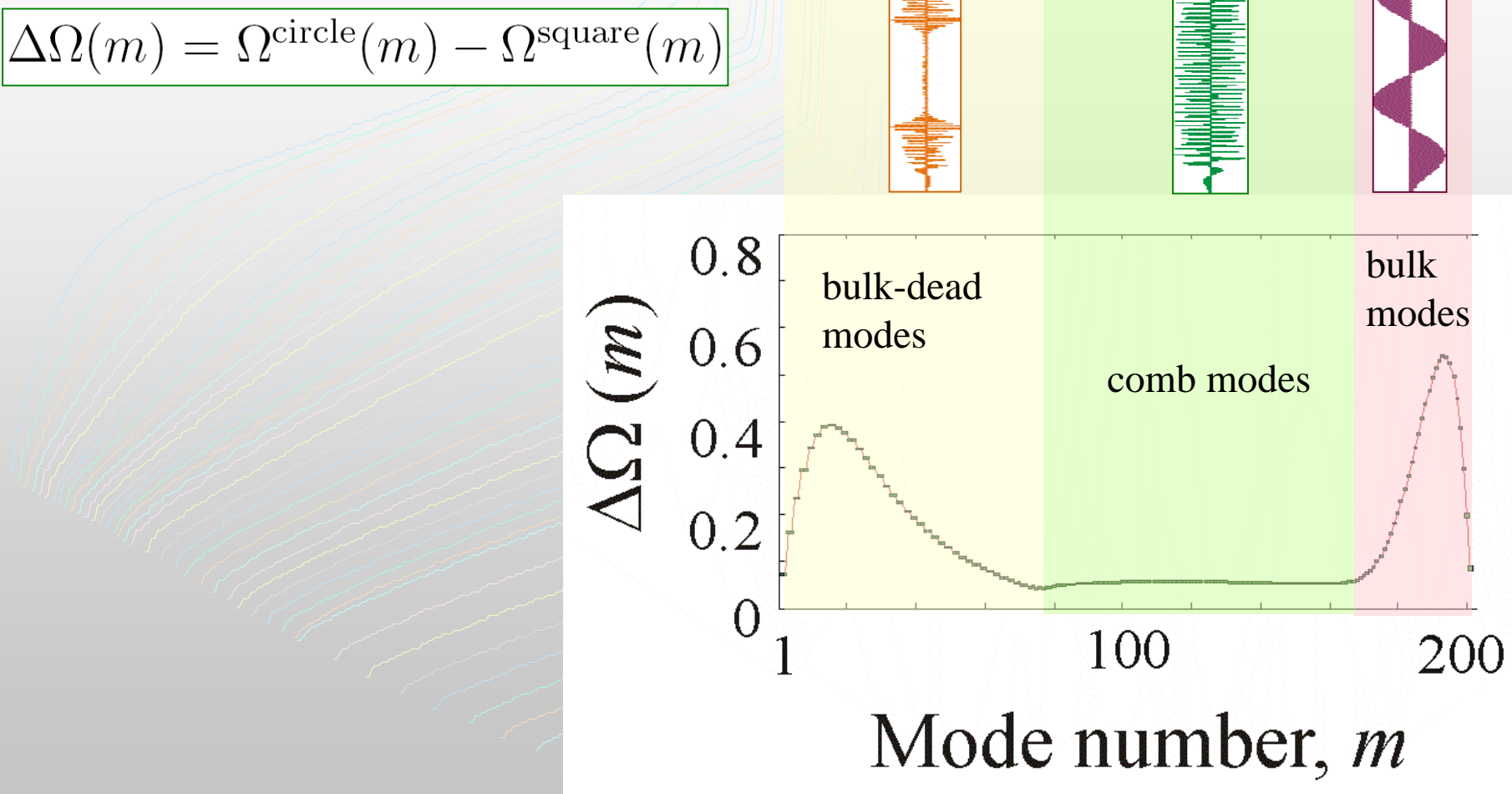
local field, $H(n)$

local field difference, $\Delta H(n)$



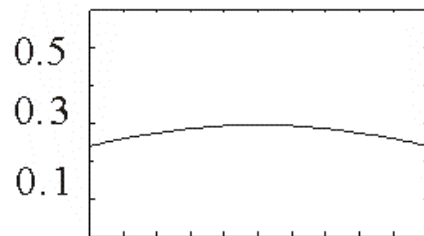
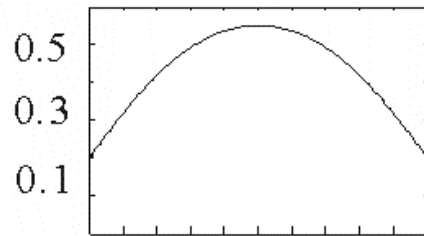
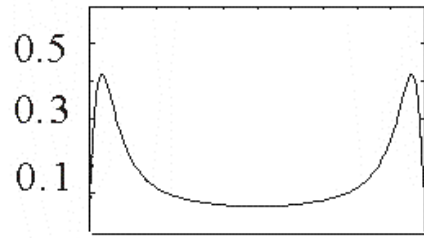


$$\Delta\Omega(m) = \Omega^{\text{circle}}(m) - \Omega^{\text{square}}(m)$$



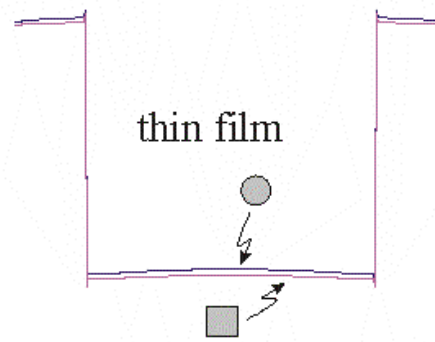
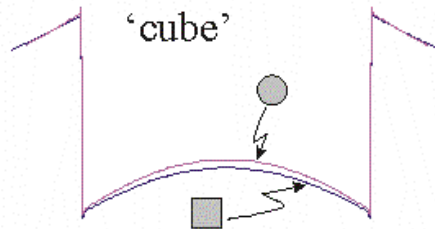
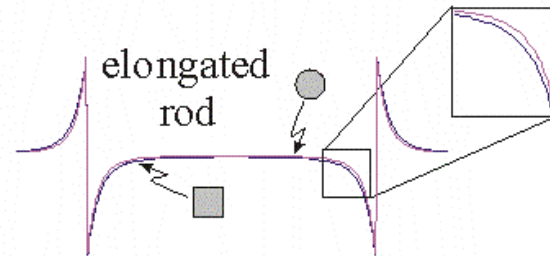
Effective dipolar field dependence on cross-sectional shape and size

local field difference,
 $\Delta H(n)$



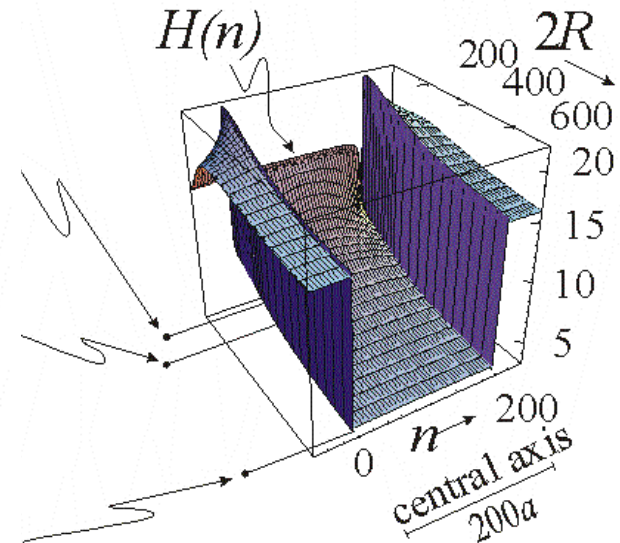
central axis
 $n = 0$ $n \rightarrow$ $n = 200$

local field, $H(n)$

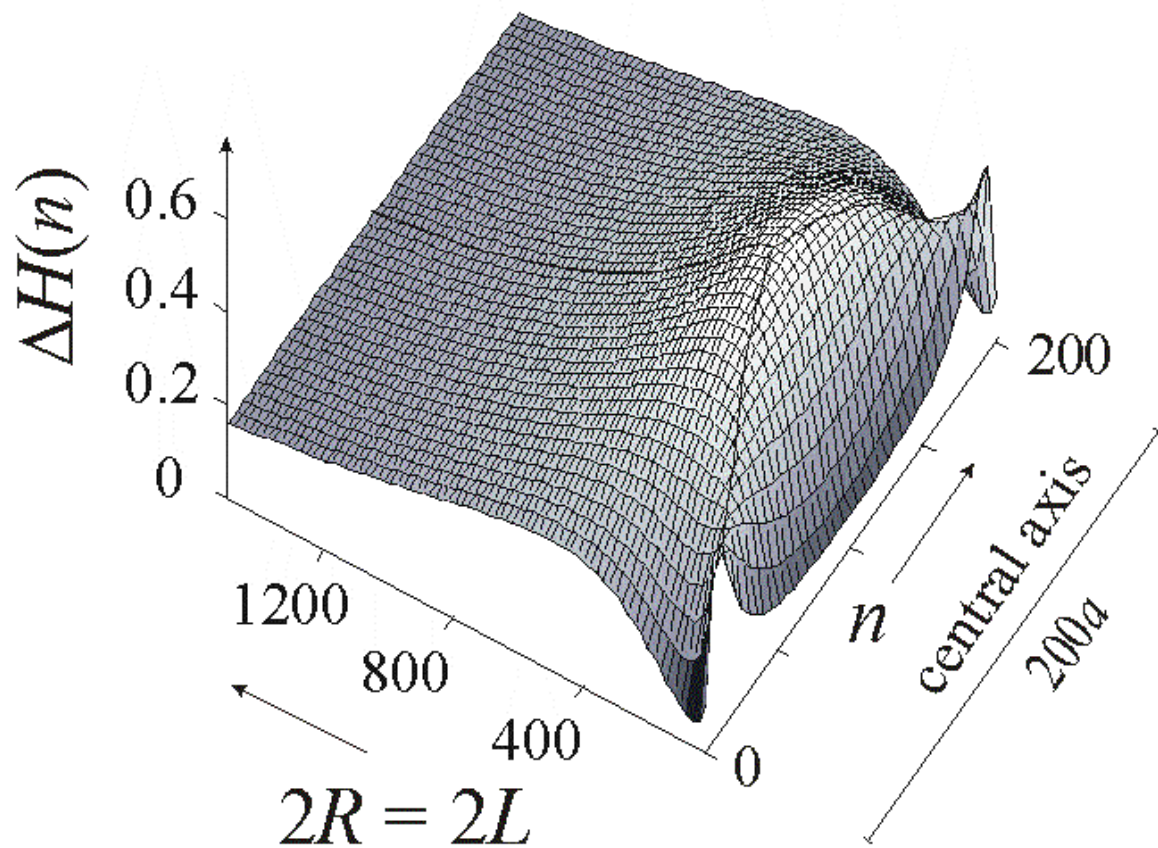


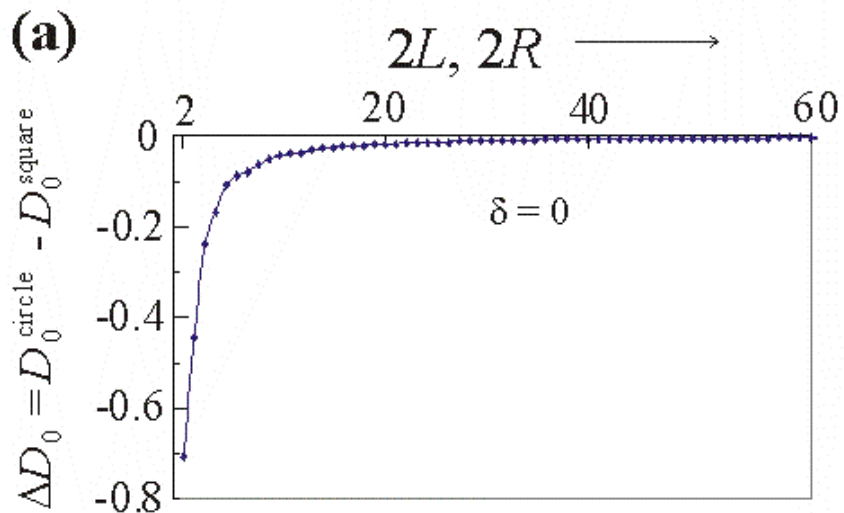
central axis
 $n = 0$ $n \rightarrow$ $n = 200$

cylindrical rod

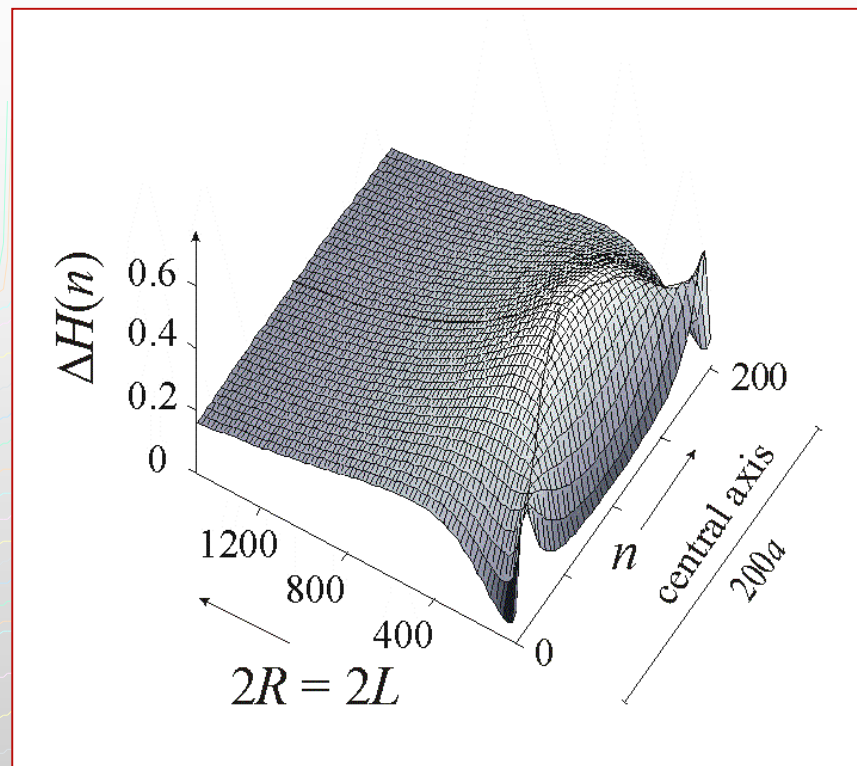
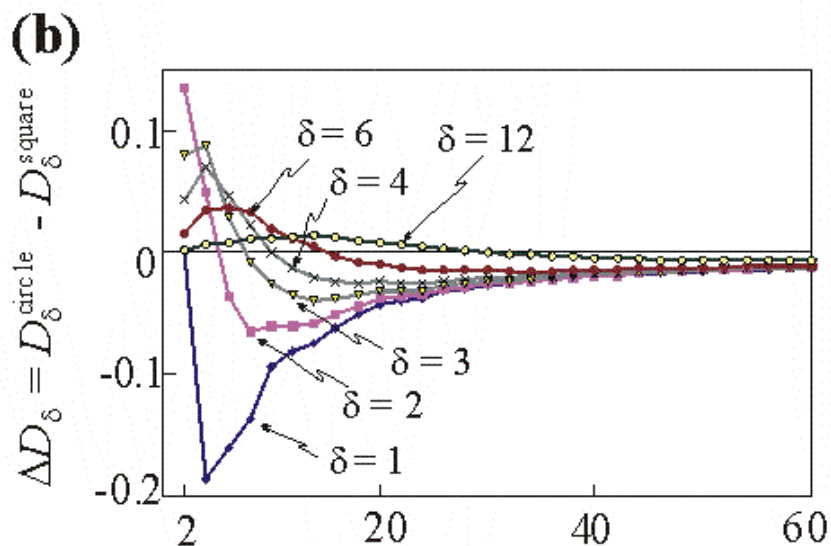


$$\Delta H(n) \equiv H^{\text{circle}}(n) - H^{\text{square}}(n) = -3\Delta D_0 - 2 \sum_{\delta=\pm 1, \dots}^n \Delta D_\delta$$





$$D_\delta \equiv \sum_{p,q} \frac{\frac{1}{2}(p^2 + q^2) - \delta^2}{[p^2 + q^2 + \delta^2]^{\frac{5}{2}}}$$



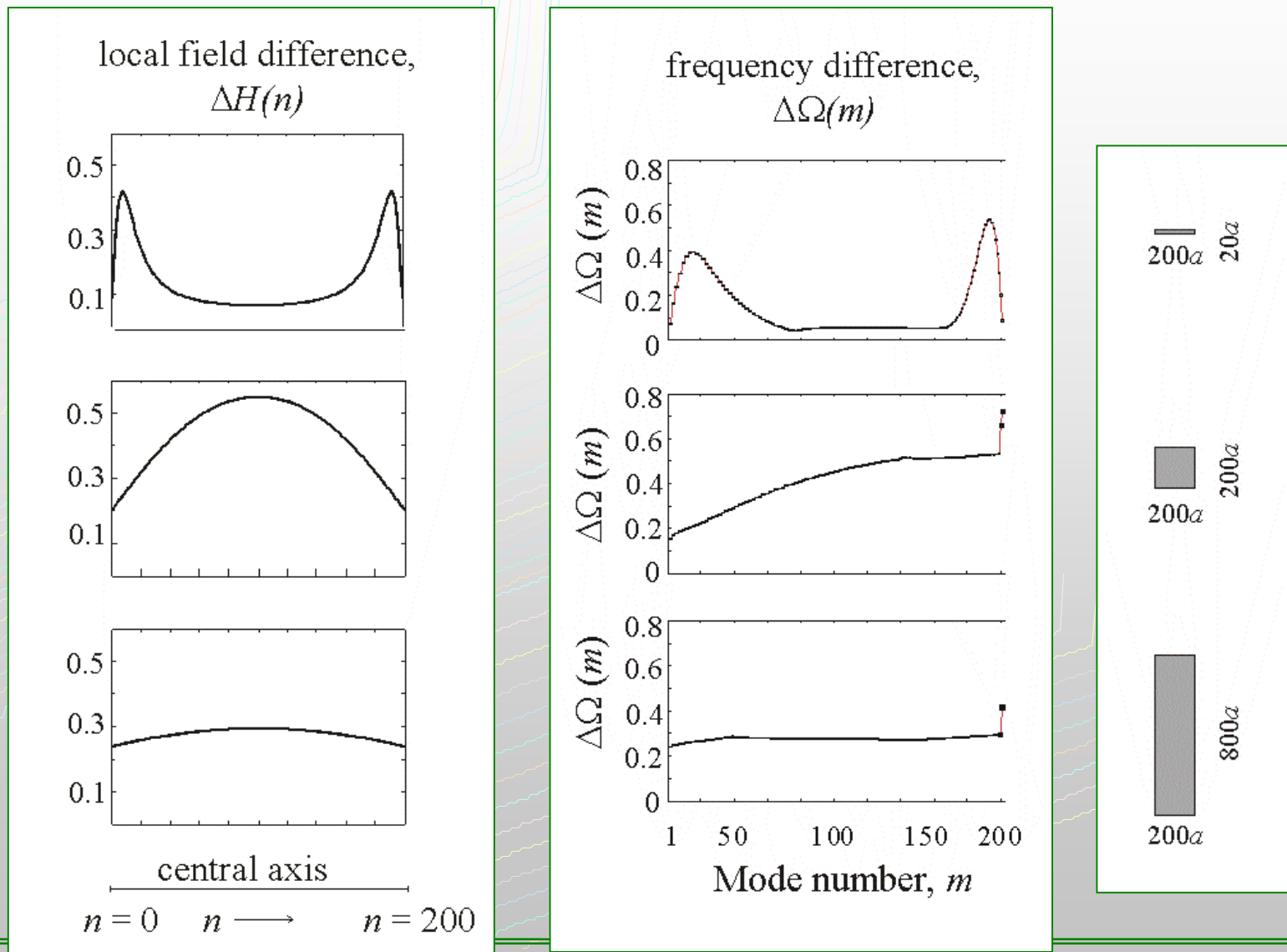
$2L, 2R \longrightarrow$

$$\Delta H(n) \equiv H^{\text{circle}}(n) - H^{\text{square}}(n) = -3\Delta D_0 - 2 \sum_{\delta=\pm 1, \dots} n \Delta D_\delta$$

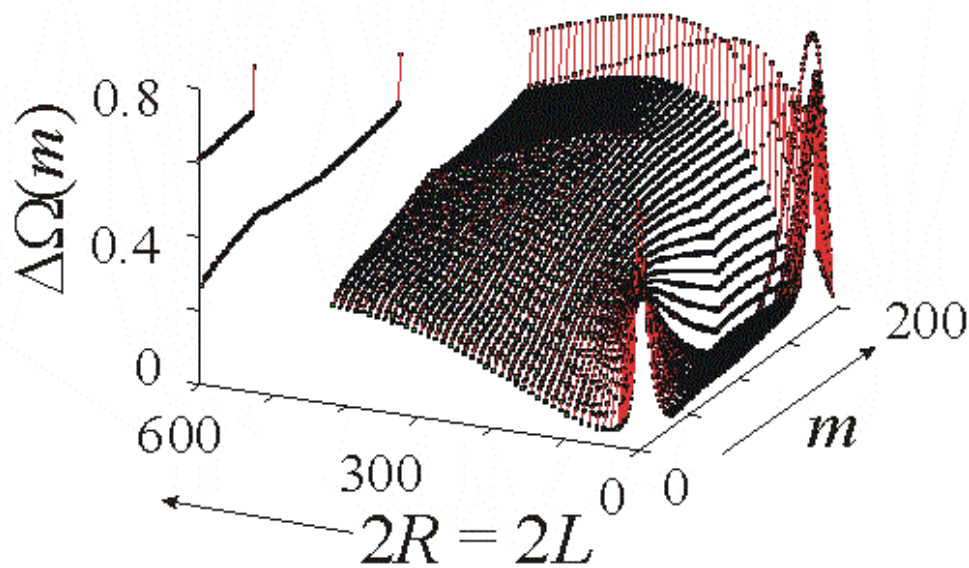
Conclusion

Magnetostatic modes localized in the wells of local dipolar field, so-called *bulk-dead modes*, are most sensitive to the cross-sectional shape of nanorods.

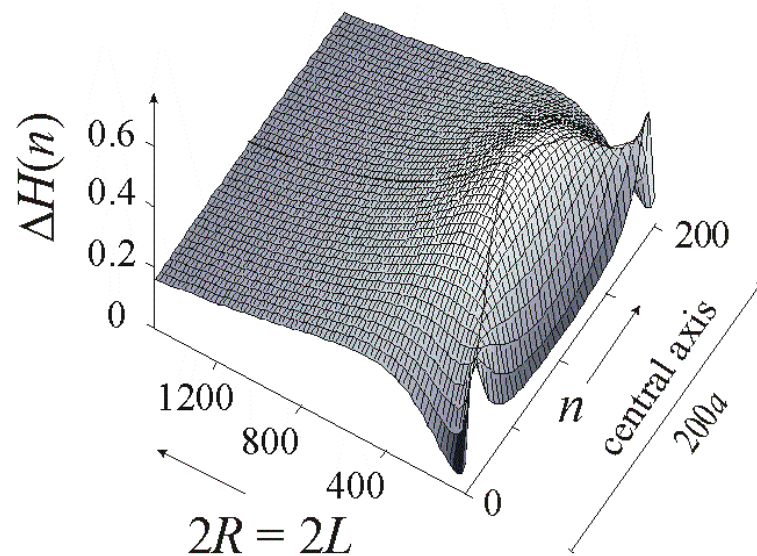
Differences of local dipolar field and frequencies of magnetostatic modes



Difference of the frequencies



Difference of the local fields



Dipolar matrix dependence on cross-sectional shape and size

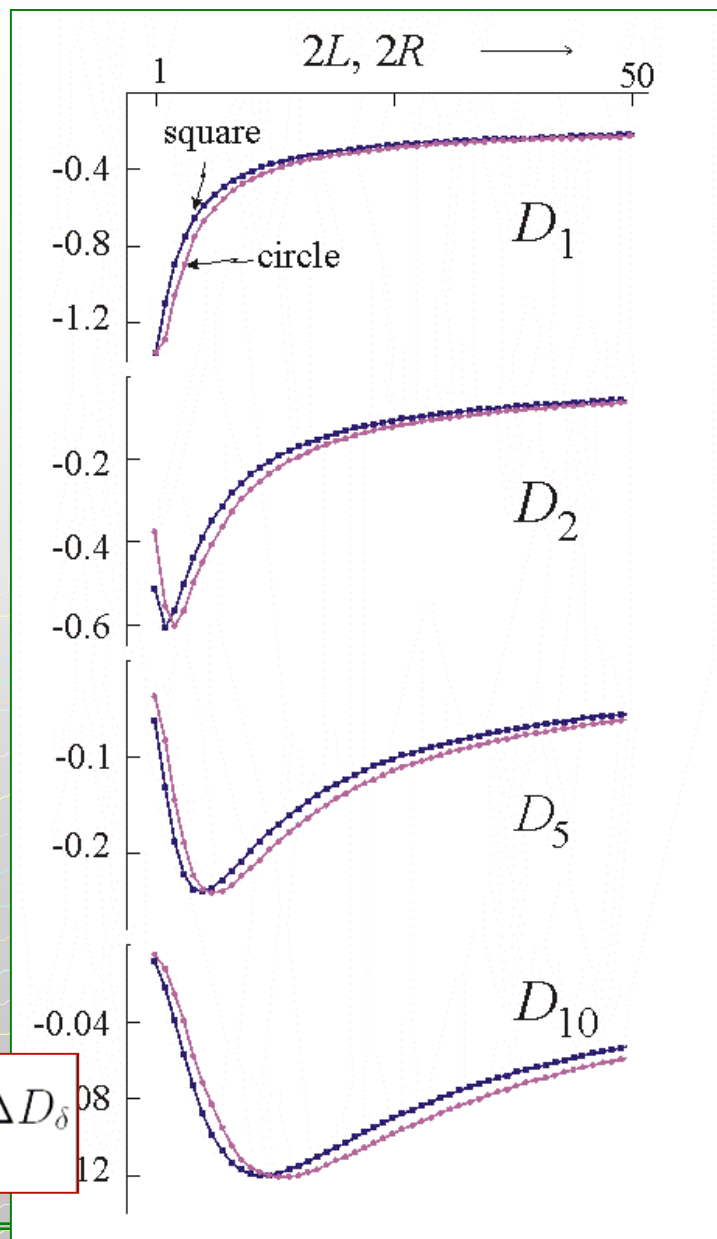
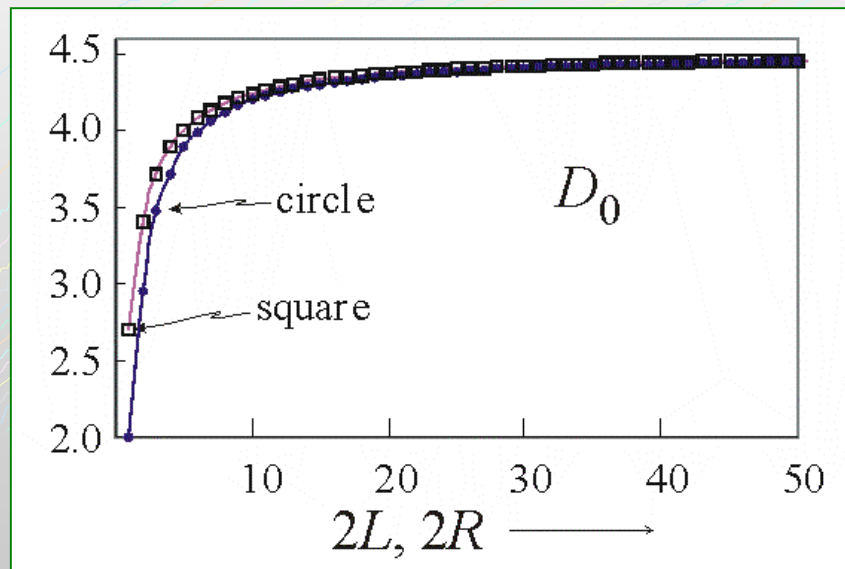
$$D_\delta \equiv \sum_{p,q} \frac{\frac{1}{2}(p^2 + q^2) - \delta^2}{[p^2 + q^2 + \delta^2]^{\frac{5}{2}}}$$

square

circle

$$-L \leq p, q \leq L$$

$$p^2 + q^2 \leq R^2$$



$$\Delta H(n) \equiv H^{\text{circle}}(n) - H^{\text{square}}(n) = -3\Delta D_0 - 2 \sum_{\delta=\pm 1, \dots} n \Delta D_\delta$$

