

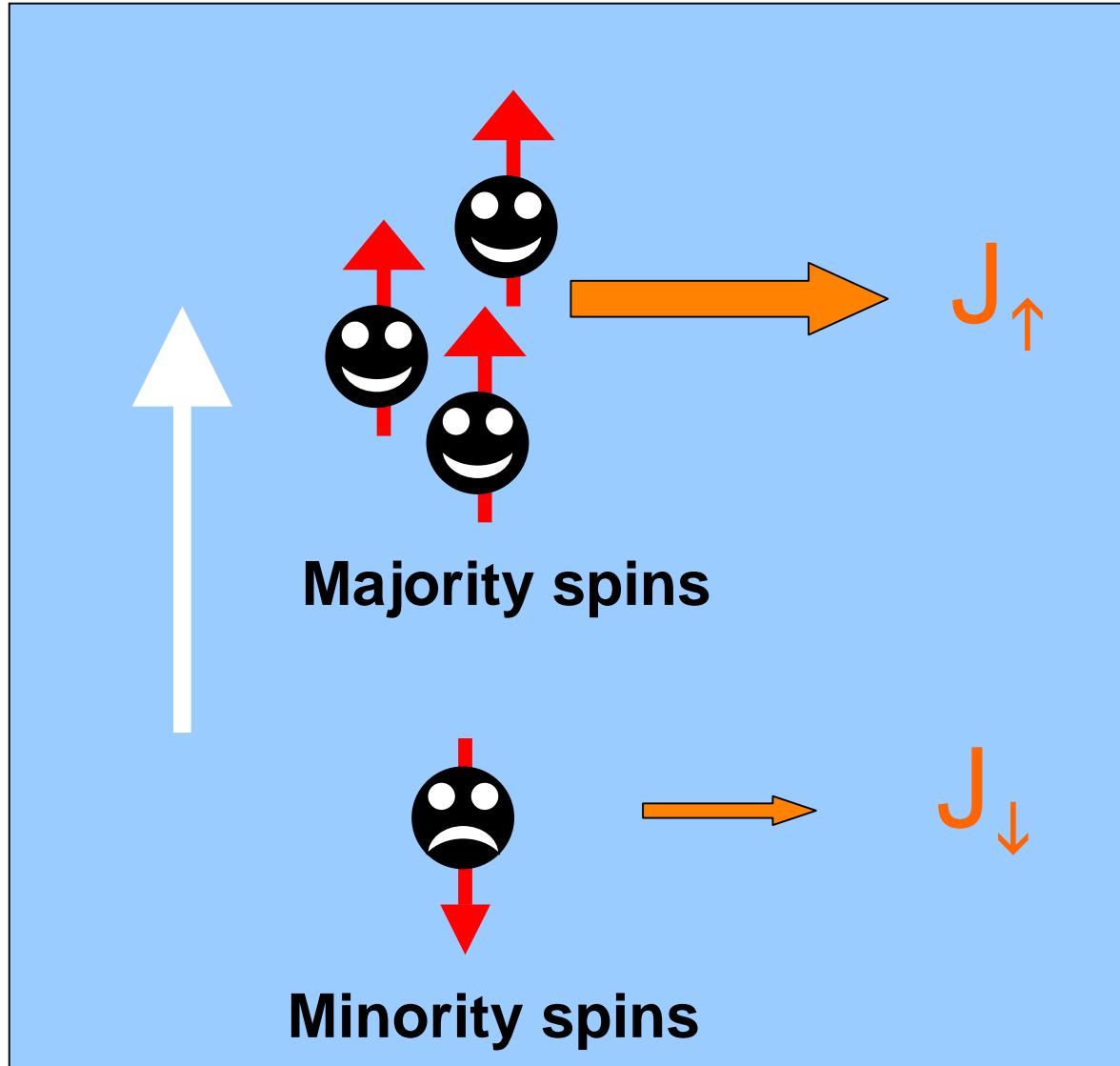
Spin transfer torque and thermally assisted FMR in magnetic tunnel junctions

C. Baraduc

S. Petit, N. de Mestier, C. Thirion, Y. Liu, M. Li, P. Wang, B. Diény

- Introduction to spin transfer torque (STT)
 - longitudinal (in-plane) torque
 - transverse (out-of-plane) torque
 - phase diagram
- Influence of STT on magnetic fluctuations
 - Model
 - Experimental result

Spin polarized current



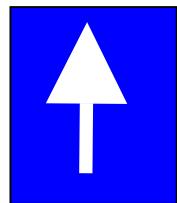
Charge current

$$J = e (J_{\uparrow} + J_{\downarrow})$$

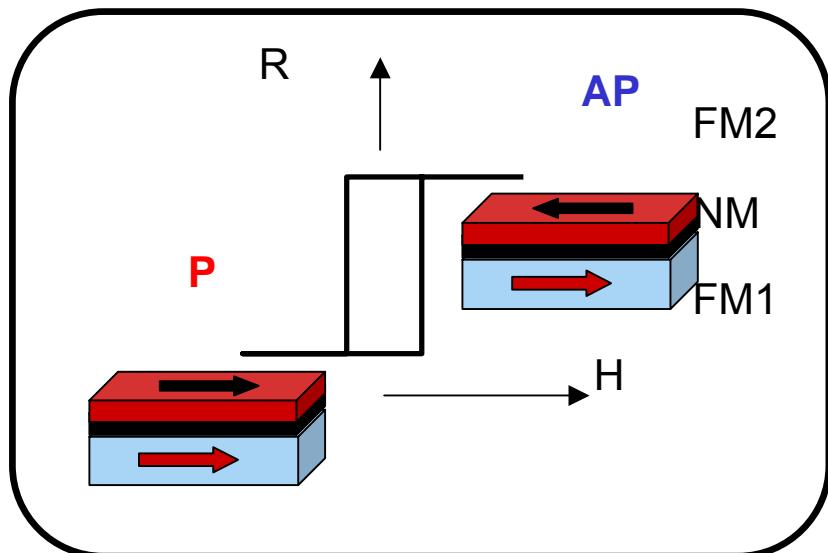
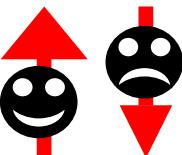
Spin current

$$J_p = J_{\uparrow} - J_{\downarrow}$$

Spin filtering

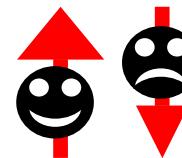


acts on

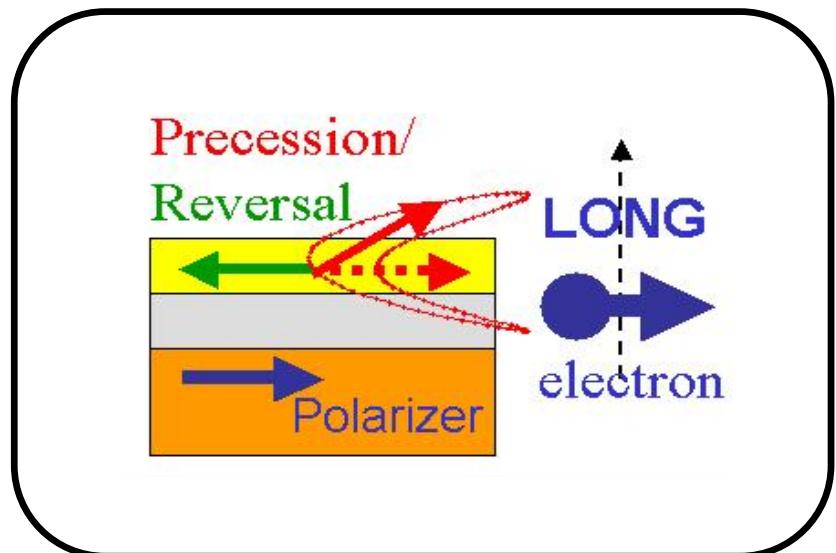
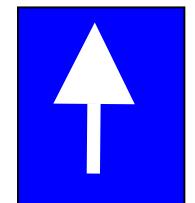


Magnetization controls current

Spin transfer torque



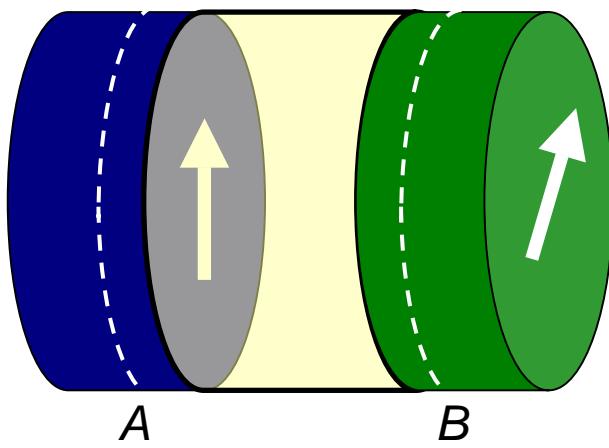
acts on



Current controls magnetization

Longitudinal (in-plane) spin transfer torque

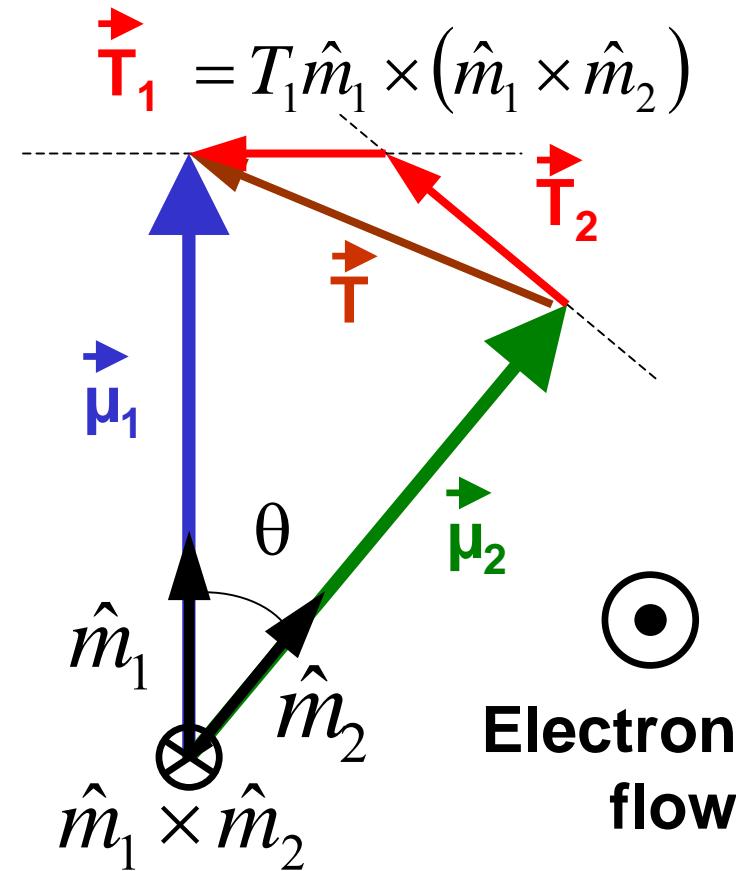
F_1 F_2



Electron flow

Conservation of length

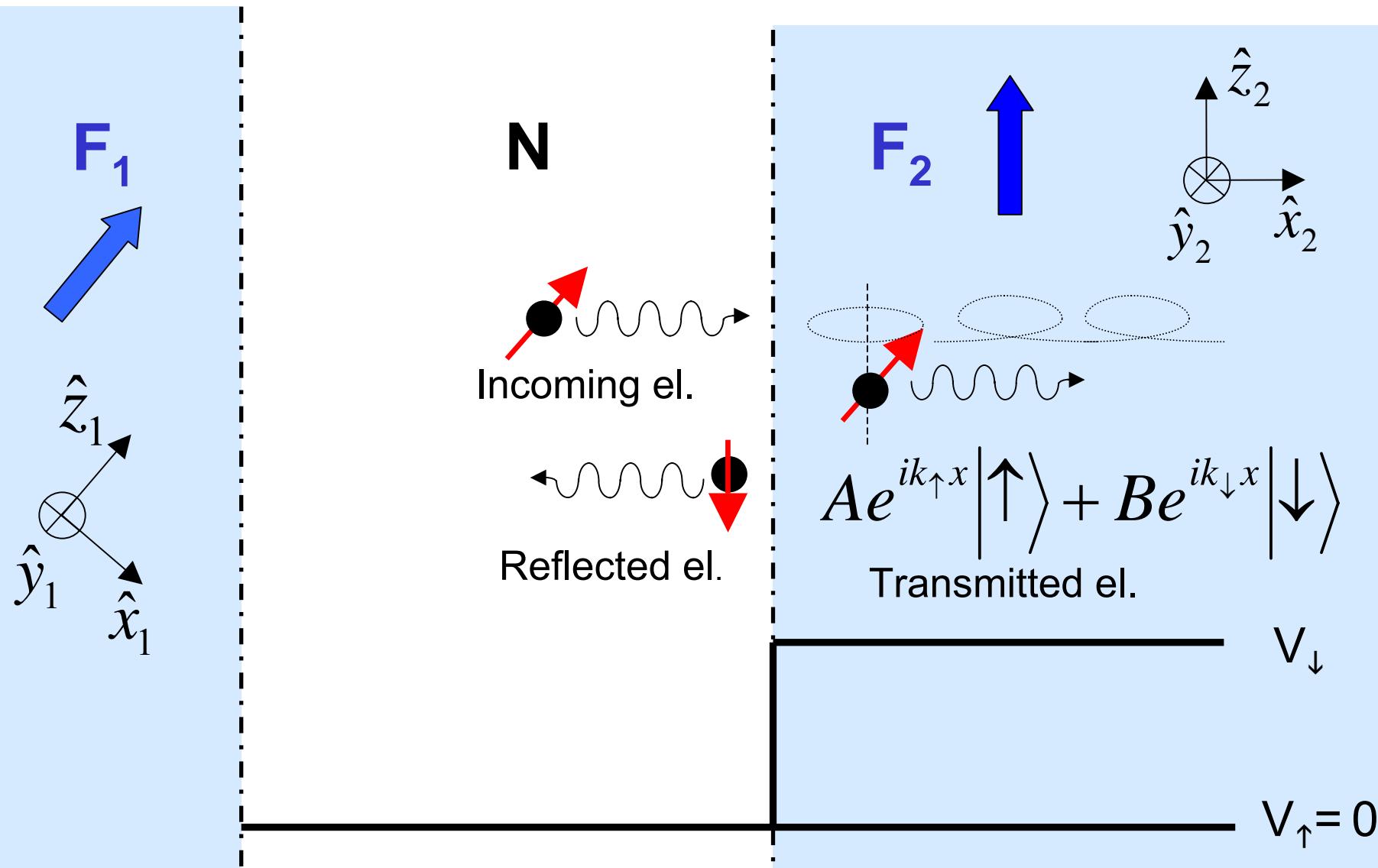
$$\frac{dM^2}{dt} = 2\vec{M} \cdot \frac{d\vec{M}}{dt} = 0$$



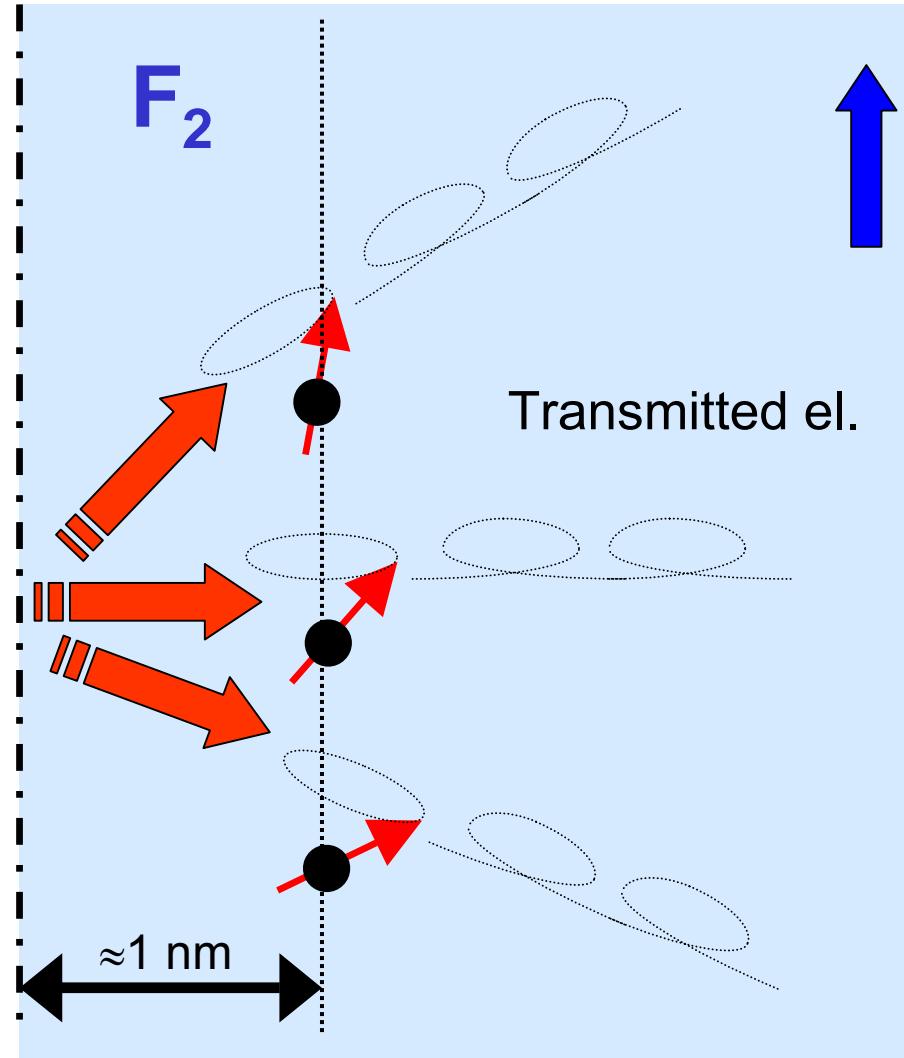
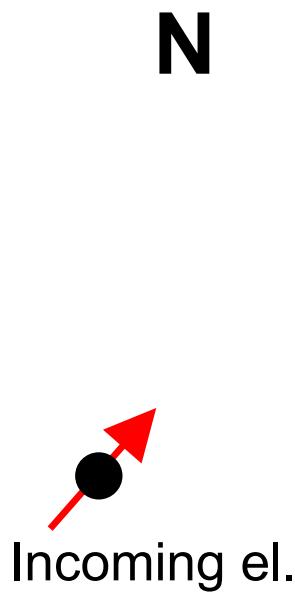
Electron flow

$$\vec{T}_{//} = \gamma_0 \frac{a_j}{M_s} \vec{M} \times (\vec{M} \times \vec{p})$$

Microscopic picture



Microscopic picture (2)



Classical dephasing → transverse component is transferred

Transverse (out of plane) torque

Inter-layer coupling energy

$$E = J_{ex} \vec{M}_1 \cdot \vec{M}_2$$

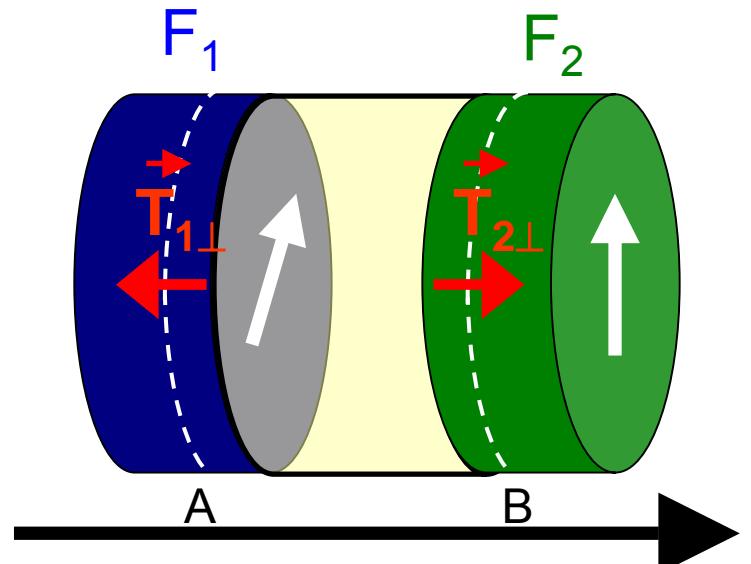
Equivalent to a **field** in the direction of the magnetization of the other layer

$$\vec{H}_1 = -\partial E / \partial \vec{M}_1 \propto \vec{M}_2$$

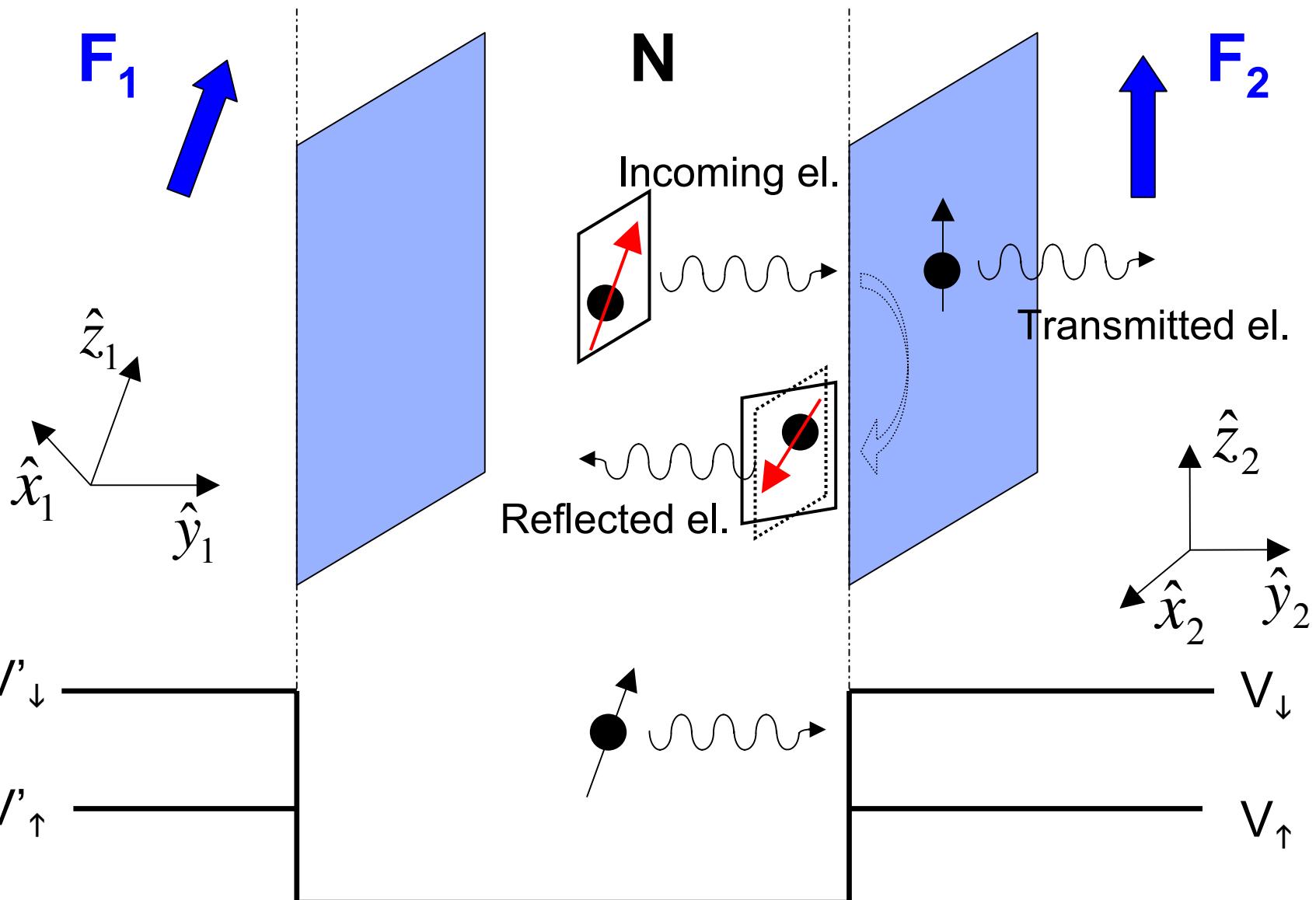
Equivalent to a torque

$$\vec{T}_\perp \propto \vec{M}_1 \times \vec{M}_2$$

$$\vec{T}_\perp = \gamma_0 b_j \vec{M} \times \vec{p}$$



Microscopic picture



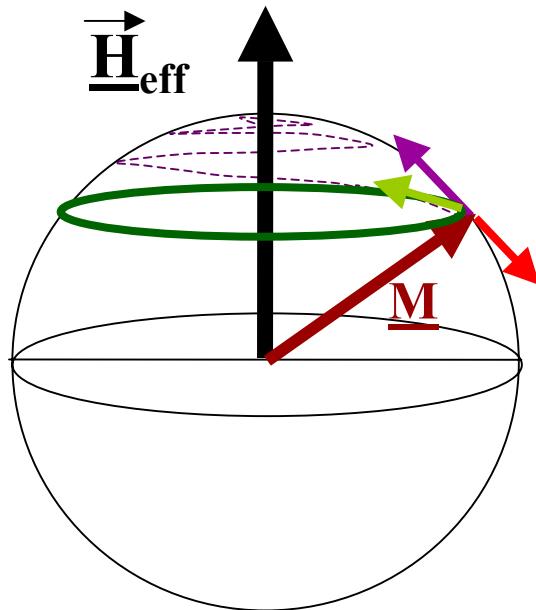
Magnetization dynamics

$$\frac{d\vec{M}}{dt} = -\gamma_0 \vec{M} \times \vec{H}_{eff} + \frac{\alpha}{M_s} \vec{M} \times \frac{d\vec{M}}{dt} + \frac{\gamma_0 a_j}{M_s} \vec{M} \times (\vec{M} \times \vec{P})$$

Larmor
precession

Damping
(Gilbert)

Longitudinal spin torque
(Slonczewski)



Low current
→ damped motion



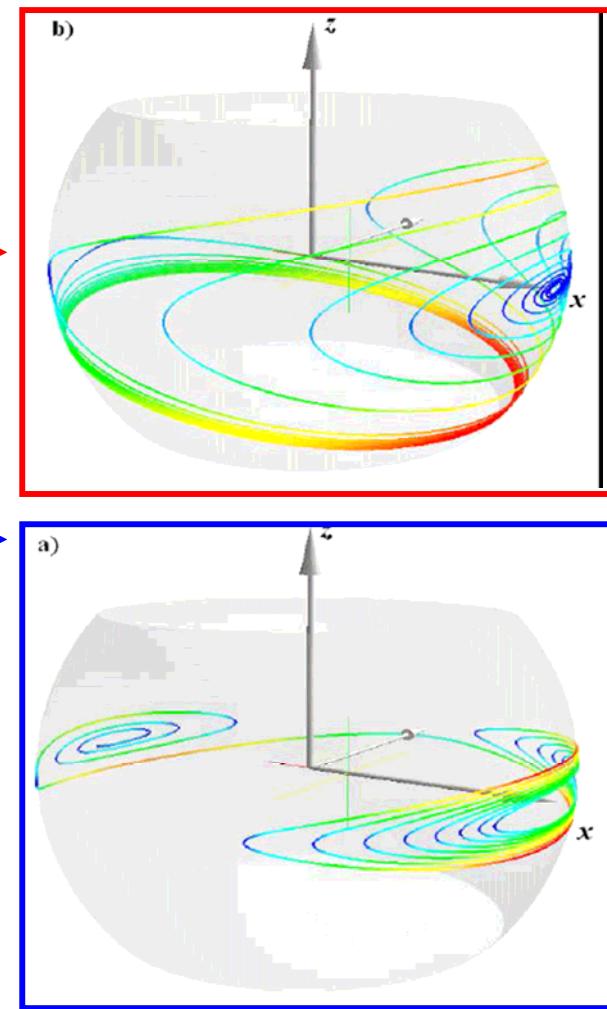
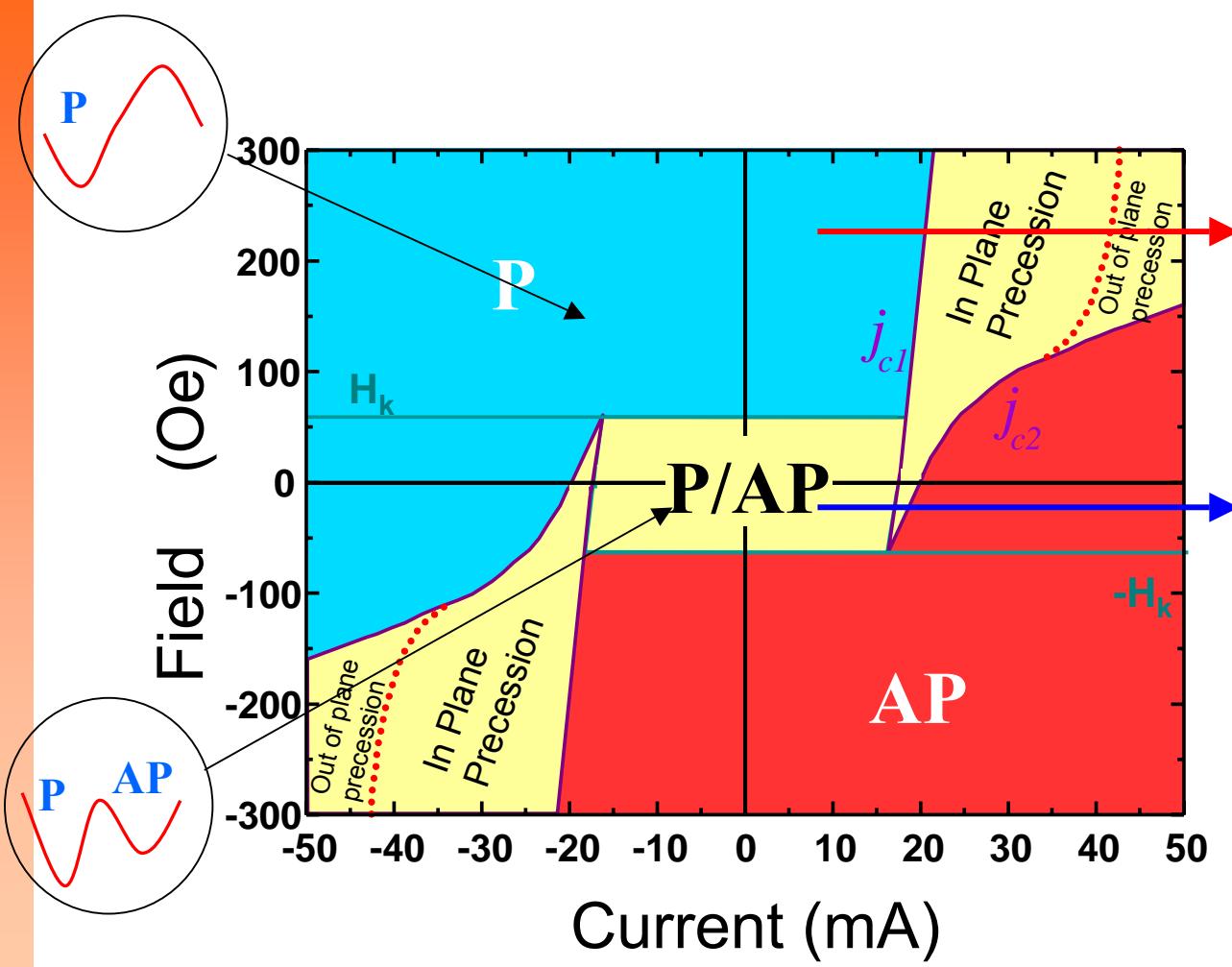
High current
→ precession



High current
→ switching

Ralph and Stiles, JMMM(2008)

Phase diagram

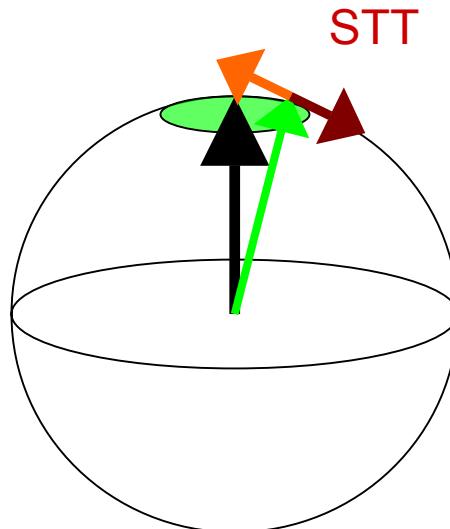


Stiles and Miltat, *Spin Dynamics in Confined Magnetic Structures III*, 225-308, (2006)

STT at equilibrium

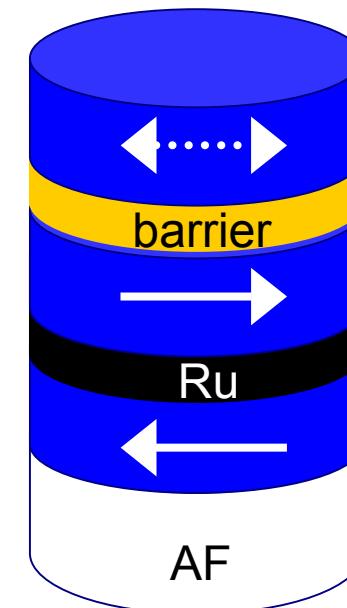
Influence of STT on mag. fluctuations

- Linear regime
- Stabilizing/Destabilizing torque



TMR Read-heads:

Tunnel junction {



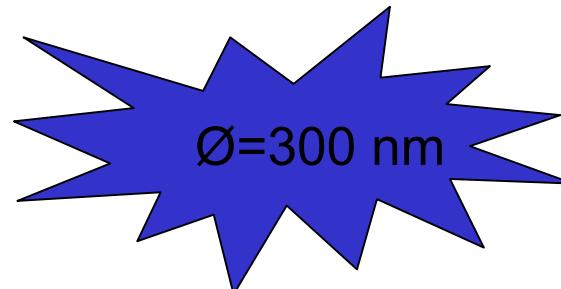
Free Layer

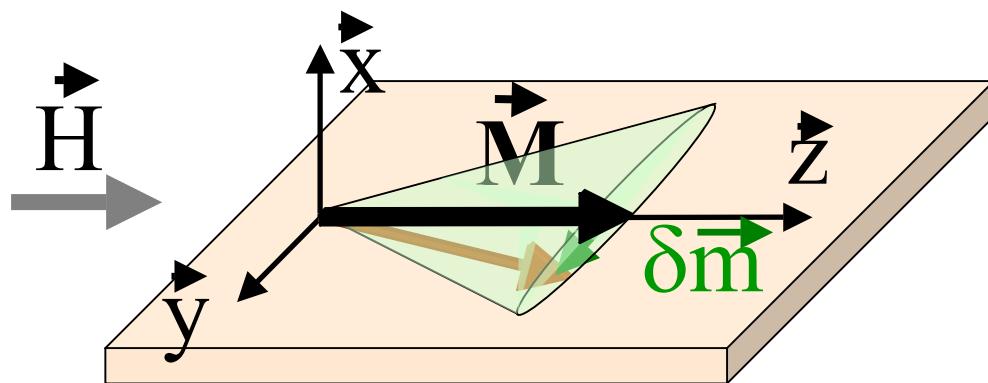
Reference Layer

Fixed Layer



Headway
Technologies





$$\delta\vec{m}(\omega) = \underline{\chi}(\omega) \delta\vec{h}_T(\omega)$$

$$\frac{d\vec{M}}{dt} = -\gamma_0 \vec{M} \times (\vec{H}_{eff} + \boxed{\delta\vec{h}_T}) + \frac{\alpha}{Ms} \vec{M} \times \frac{d\vec{M}}{dt}$$

$$+ \frac{\gamma_0 a_J}{Ms} \vec{M} \times (\vec{M} \times \hat{p}) + \gamma_0 b_J \vec{M} \times \hat{p}$$

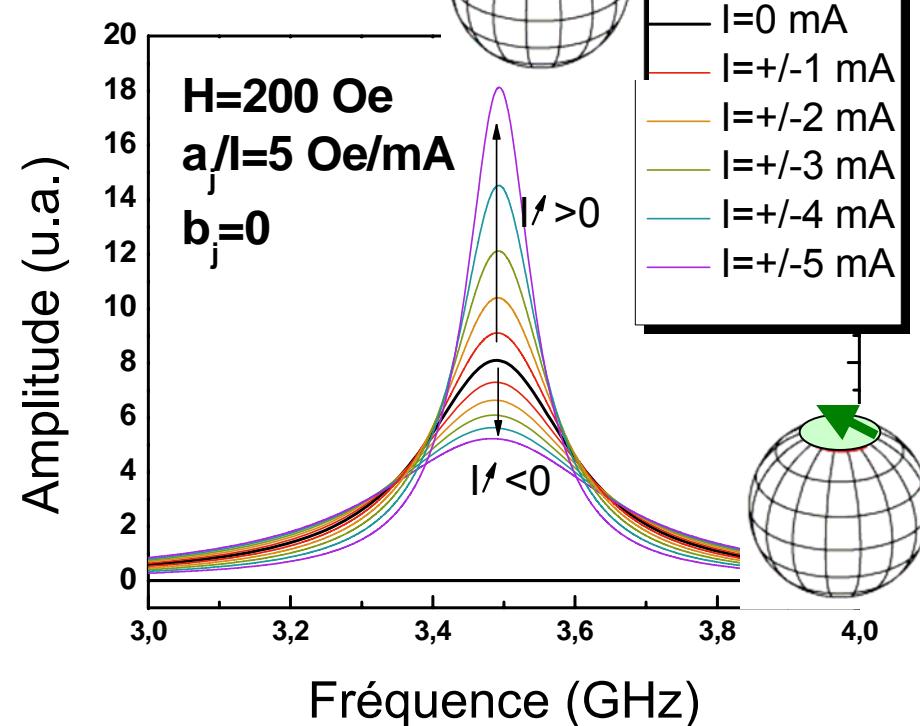
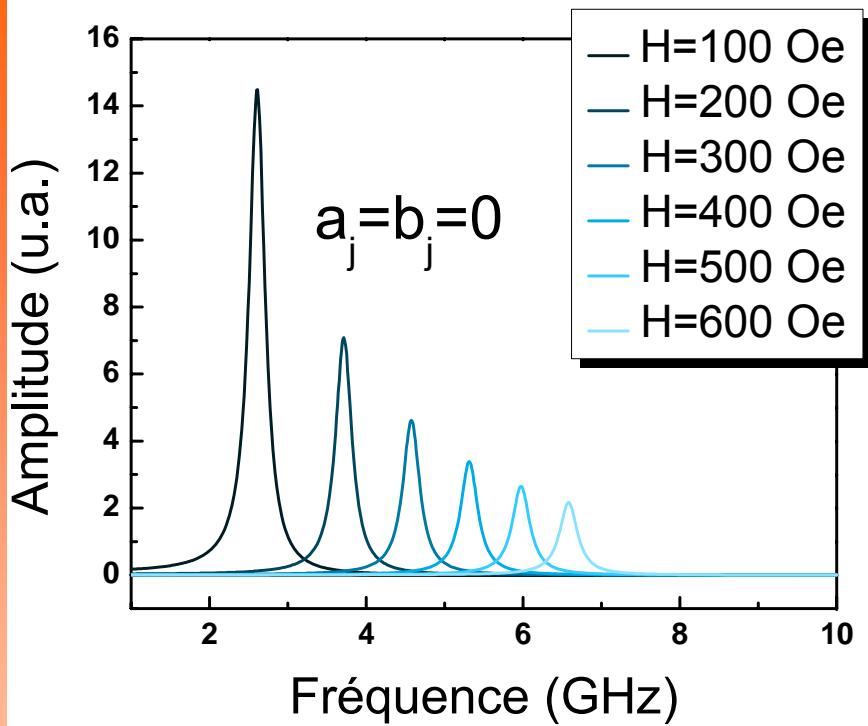
Spin torque $T//$

$T\perp$

Fluctuation-dissipation theorem gives the magnetization Power Spectral Density (PSD) :

$$S_{M_y} = \frac{4kT}{\mu_0 V} \frac{\chi''_{yy}}{\omega}$$

Model predictions



$$\omega_0^2 \approx \gamma_0^2 \left[H(4\pi M_s + H) - (4\pi M_s + 2H) \mathbf{b}_j \cdot \boldsymbol{\varepsilon} \right]$$

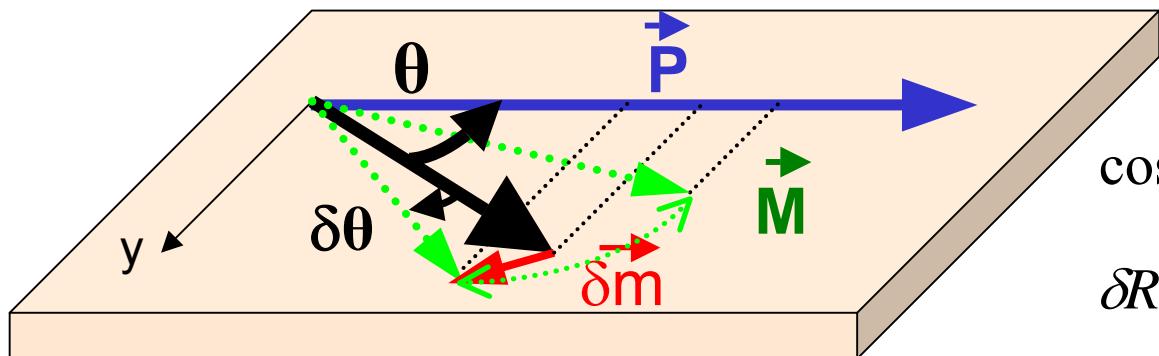
$$\Lambda = \gamma_0 \alpha (2H + 4\pi M_s) - \gamma_0 2 \mathbf{a}_j \cdot \boldsymbol{\varepsilon}$$

 $\varepsilon = 1 \text{ (P)}$
 $\varepsilon = -1 \text{ (AP)}$

From δm to δV

How to access experimentally to the magnetization fluctuations spectrum ?

$$\delta h_T \leftrightarrow \delta m \leftrightarrow \delta R \leftrightarrow \delta V$$



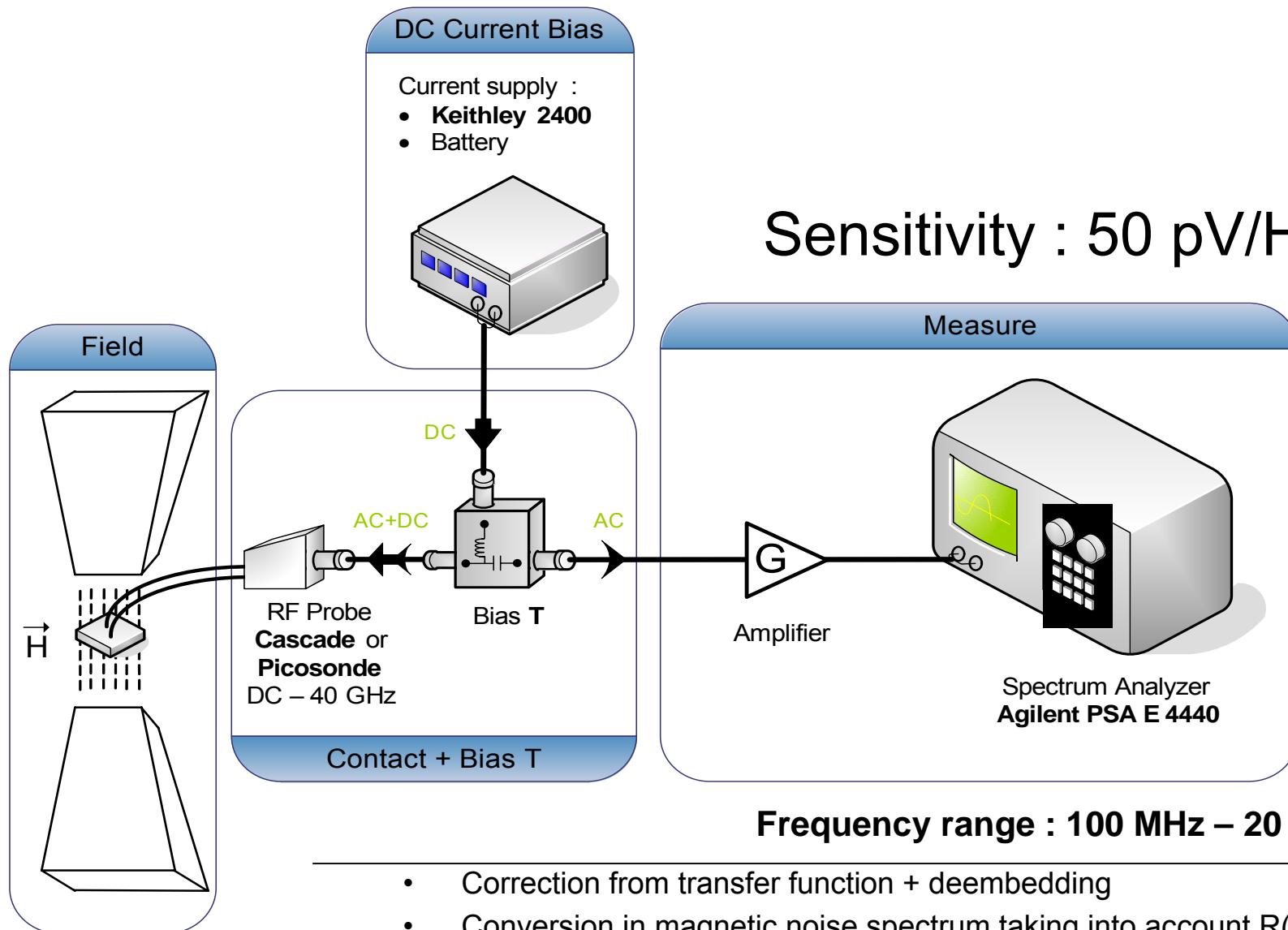
$$R = R_0 - \frac{\Delta R}{2} \cos \theta$$

$$\cos(\theta + \delta\theta) \approx \cos \theta - \sin \theta * \delta\theta$$

$$\delta R \approx \frac{\Delta R}{2} \sin \theta * \delta\theta$$

$$\delta V = I * \delta R \approx I \frac{\Delta R}{2} \sin \theta \frac{\delta m}{M_s}$$

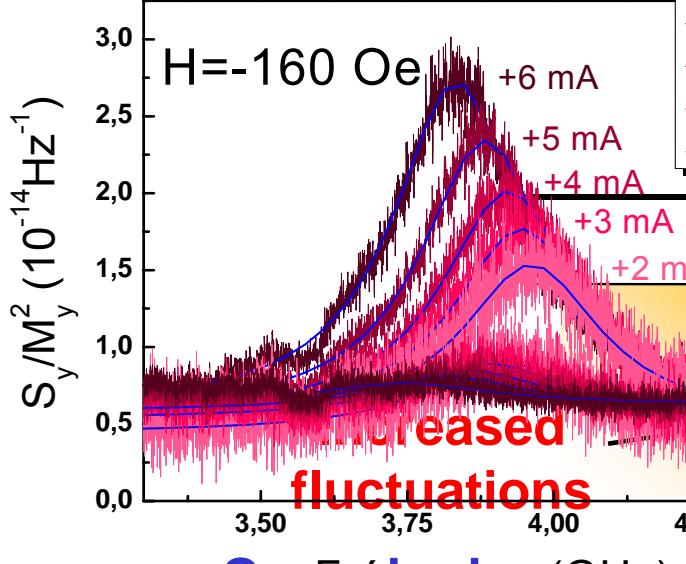
Experimental setup



- Correction from transfer function + deembedding
- Conversion in magnetic noise spectrum taking into account R(I)

Influence of STT on noise

Parallel state

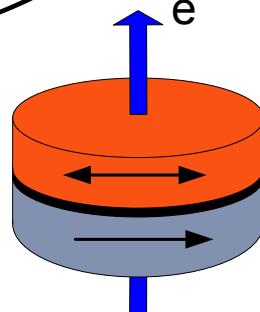


$I = \pm 6 \text{ mA}$
 $I = \pm 5 \text{ mA}$
 $I = \pm 4 \text{ mA}$
 $I = \pm 3 \text{ mA}$
 $I = \pm 2 \text{ mA}$

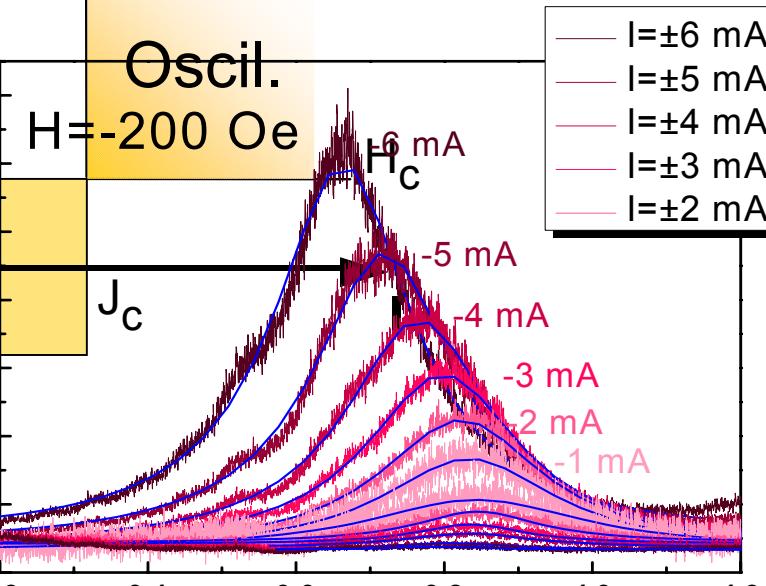
$S_y/M_y^2 (10^{-14} \text{Hz}^{-1})$
 $I < 0$

Fréquence (GHz)
Quenched fluctuations

Couche libre
Couche de référence



Antiparallel state

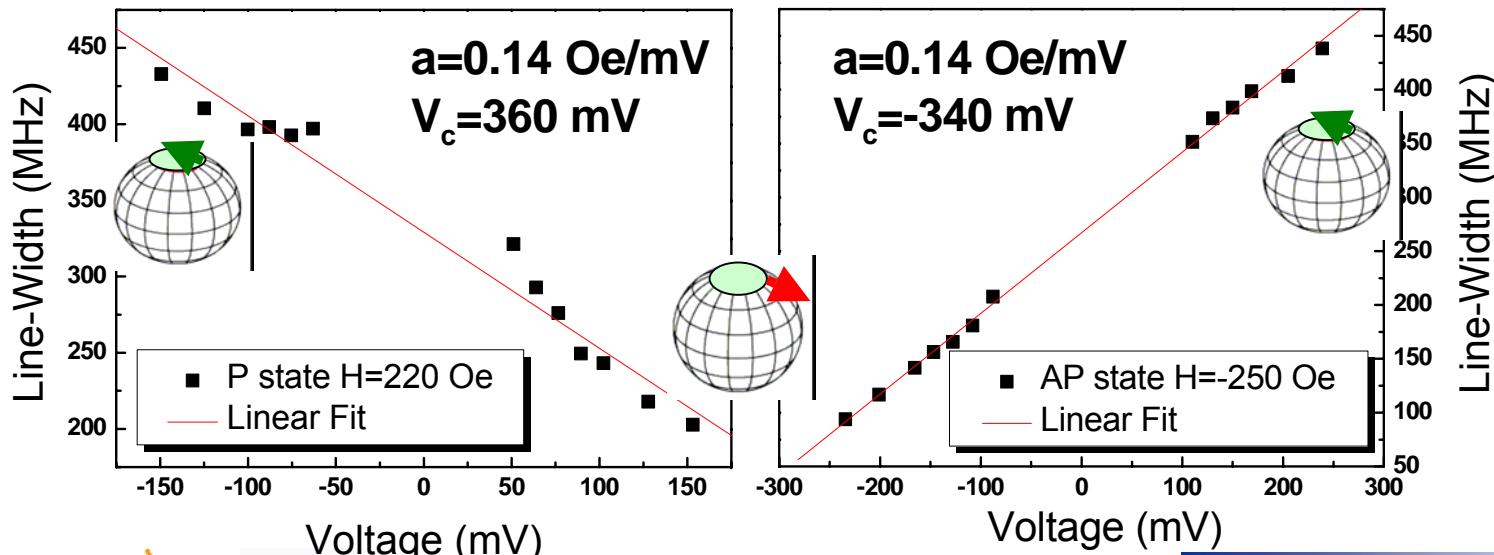
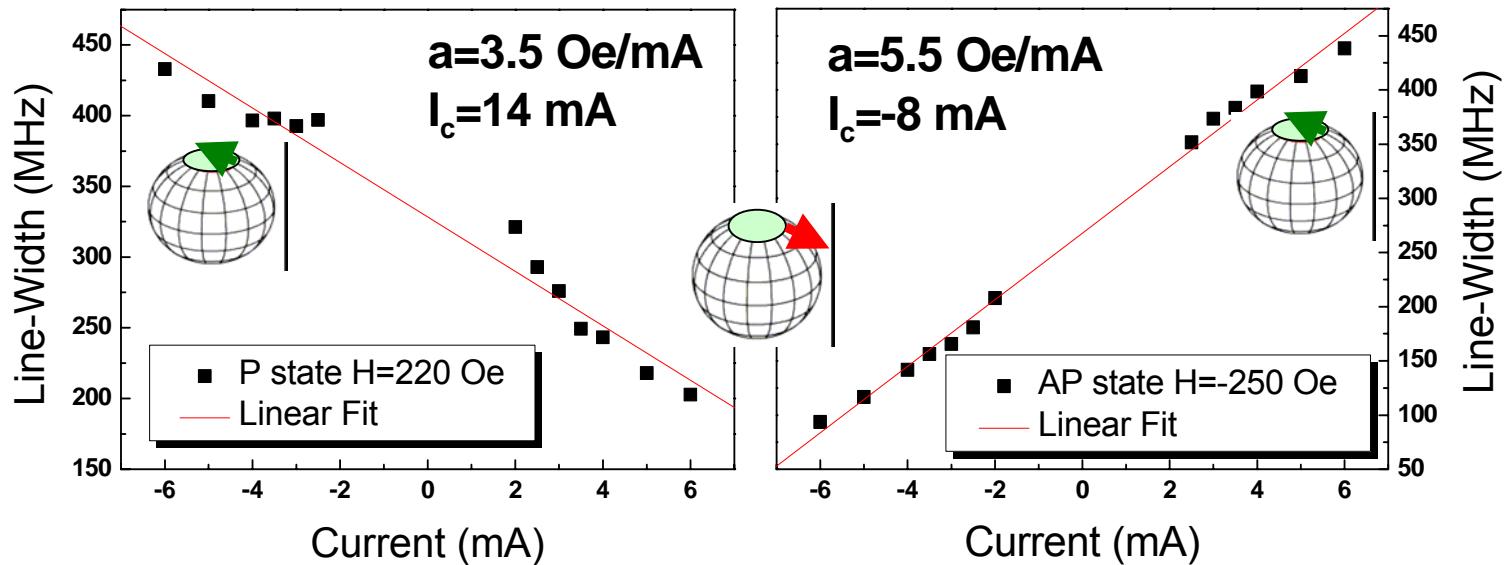


$I = \pm 6 \text{ mA}$
 $I = \pm 5 \text{ mA}$
 $I = \pm 4 \text{ mA}$
 $I = \pm 3 \text{ mA}$
 $I = \pm 2 \text{ mA}$

$S_y/M_y^2 (10^{-14} \text{Hz}^{-1})$
 $I < 0$

Fréquence (GHz)

Peak Linewidth → longitudinal torque a_j



Summary

- **Longitudinal (in-plane) STT:**
direct transfer of electron momentum → a_j
- **Transverse (out-of-plane) STT:**
field-like term (exchange coupling) → b_j
- **STT at equilibrium:**
effect on magnetization fluctuations (noise)
- Stabilizing and destabilizing torque can be measured:
extract a_j and b_j
- Voltage more relevant than current for STT in magnetic tunnel junctions.
Polarized current directly related to the voltage.

Petit et al., PRL **98**, 077203 (2007)