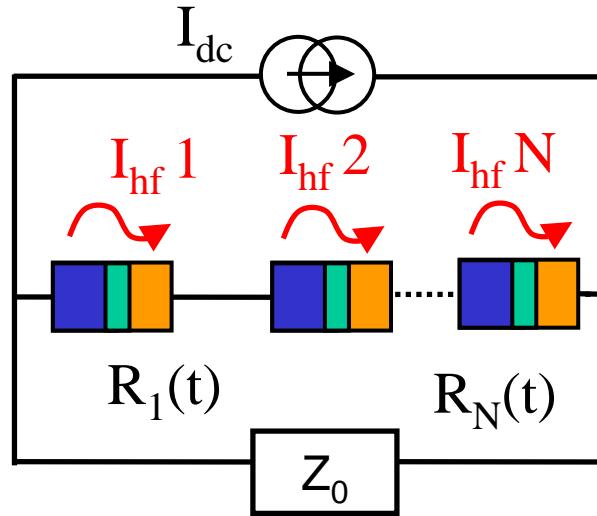
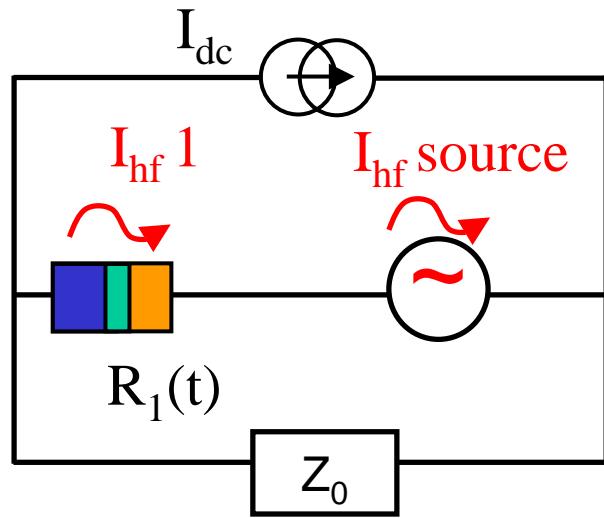


Phase locking of a Spin Transfer Nano-Oscillator to an external microwave current :



J. Grollier¹, B. Georges¹, M. Darques¹, V. Cros¹, G. Faini²,
C. Deranlot¹, B. Marcilhac¹, A. Fert¹

¹ Unité Mixte de Physique CNRS/Thales et Université Paris Sud, Palaiseau, France

² LPN CNRS, Marcoussis, France

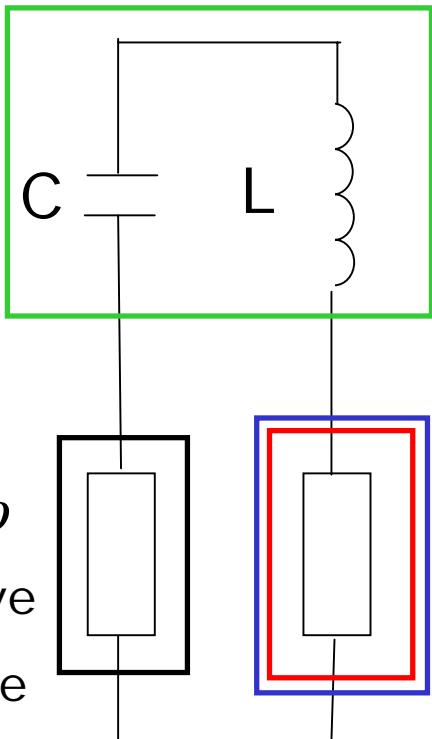


Magnetization dynamics with spin transfer



LLG equation

$$\omega = \gamma H_{eff} = \frac{1}{\sqrt{LC}}$$



$R_0 \propto \alpha_G \omega$
Dissipative
resistance

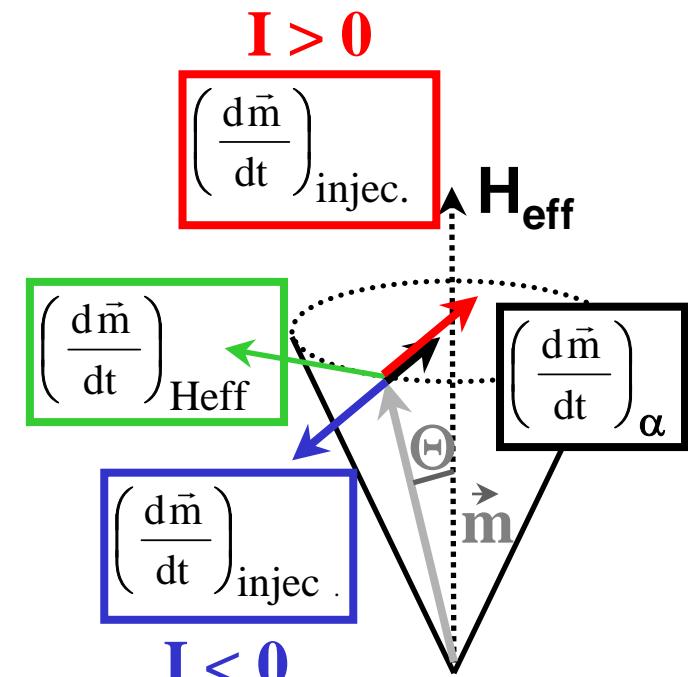
$$\frac{d\vec{m}}{dt} = -\gamma_0 \vec{m} \times \vec{H}_{eff} + \alpha \vec{m} \times \frac{d\vec{m}}{dt} + \gamma_0 J \vec{m} \times (\vec{m} \times \vec{M})$$

Precession if :

$$R_0 - R_s = 0$$

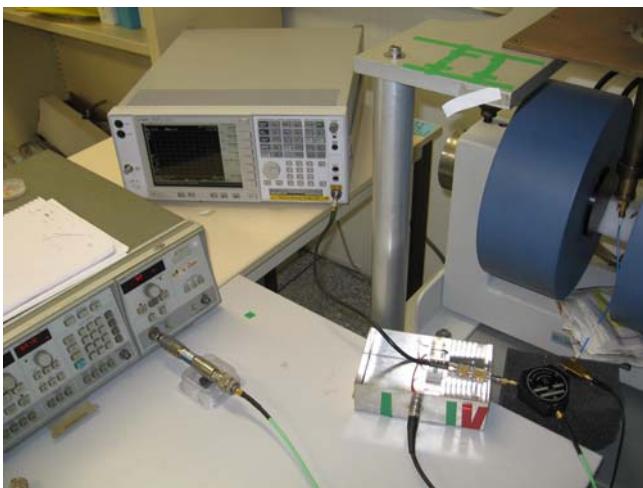
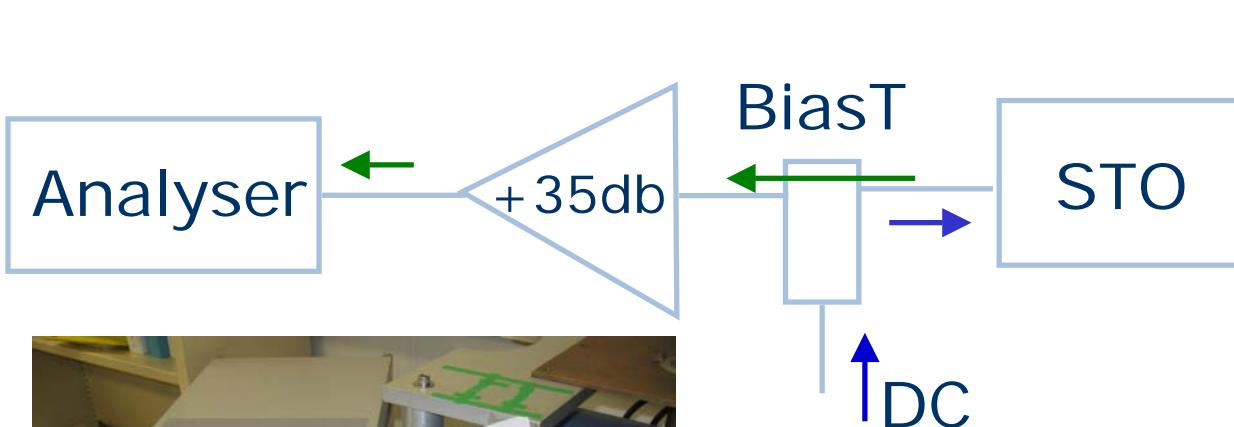
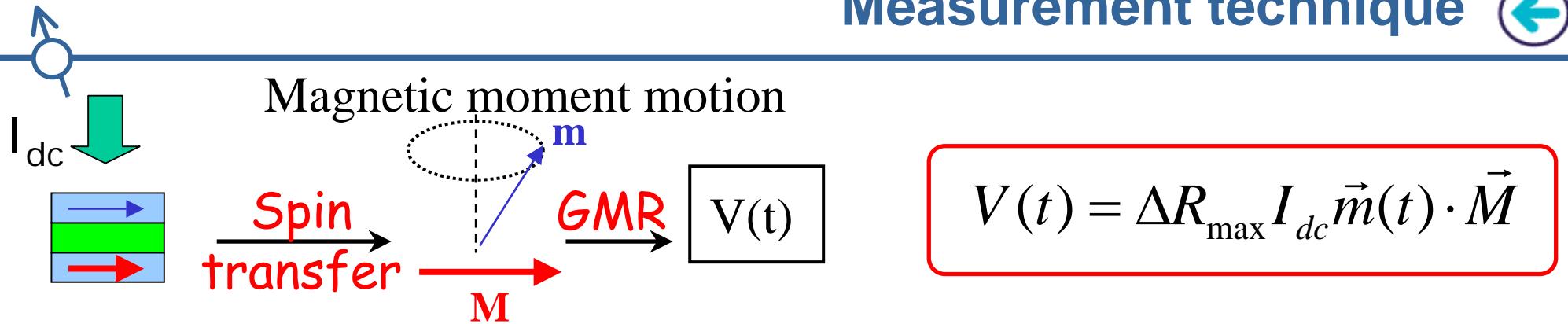
Spin torque
compensates
the damping

$-R_s \propto \gamma_0 J$
Negative
resistance

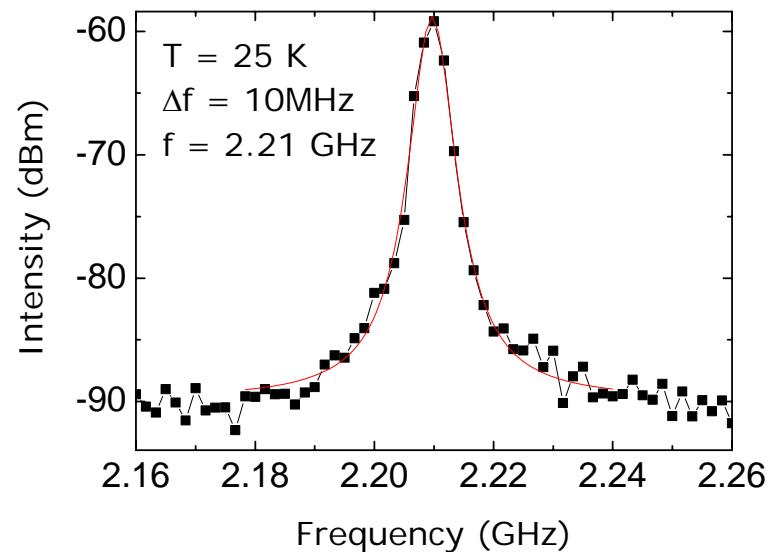


A. Slavin model *IEEE 41*, 4 (2005)

Measurement technique



$$P \approx \frac{\Delta R_{osc}^2 I_{dc}^2}{Z_{LOAD}}$$





- *non-linear ($I_{dc} \rightarrow V_{ac}$)*
- *agile (I_{dc}, H)*
- *direct emission up to 40 GHz*
- *high Q factor*
- *sub-micron size*



applications :
Telecommunications
RADAR



Problem : very weak emitted power

$P_{out} \sim -60 \text{ dBm}$

Too low for applications

Solution : synchronization

Coherent emission in frequency and phase of many oscillators

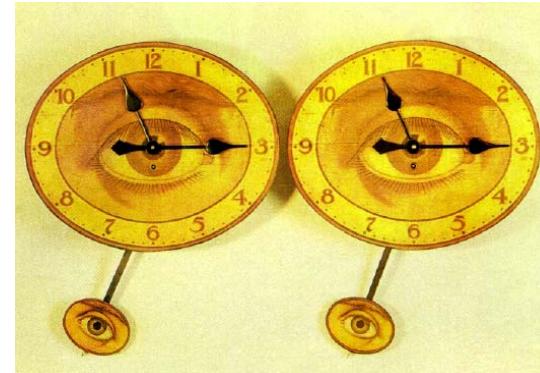


Synchronization : the oldest non-linear effect ever studied!!

Fireflies



Huygens
Pendulum
clocks (1665)



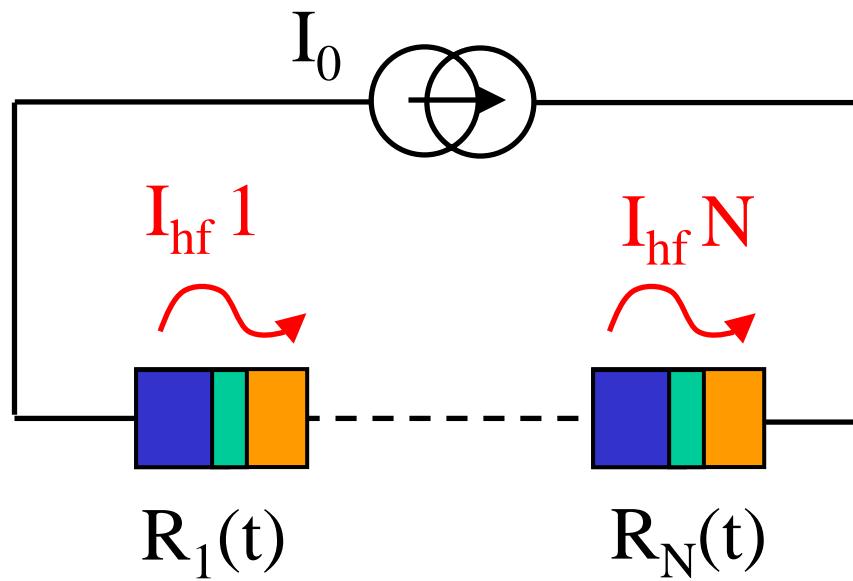
Applauding audience



- Other exemples : criquets, applauding audience, pace maker cells, intestine celluls, circadian rythm, Josephson junctions etc.



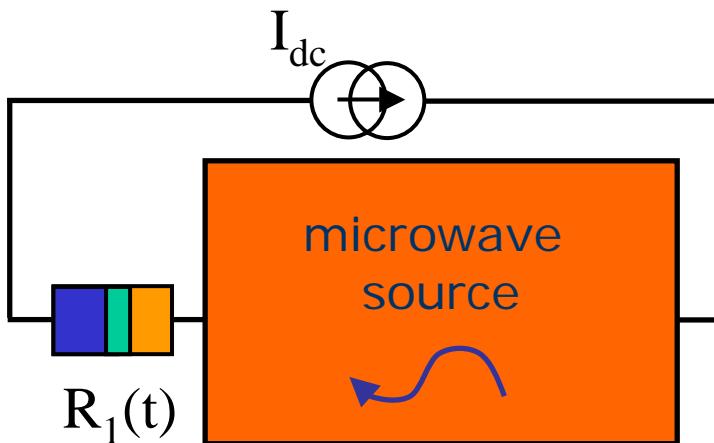
Non local coupling:
self-emitted
microwave currents



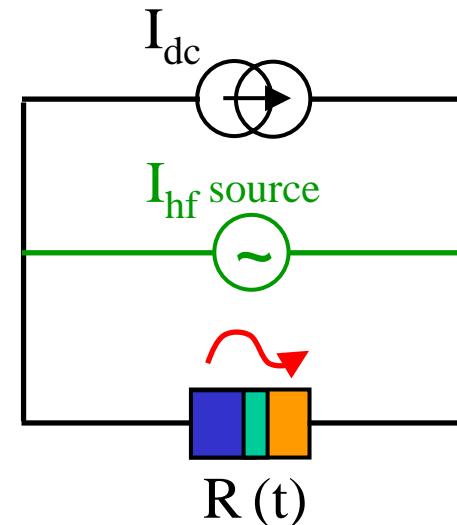
First step : phase locking to a microwave source



Electrically connected oscillators



Injection of a microwave signal



Current through ONE oscillator

$$I_{tot} = I_{dc} + \sum_{i=1}^N I_{hf}(i)$$

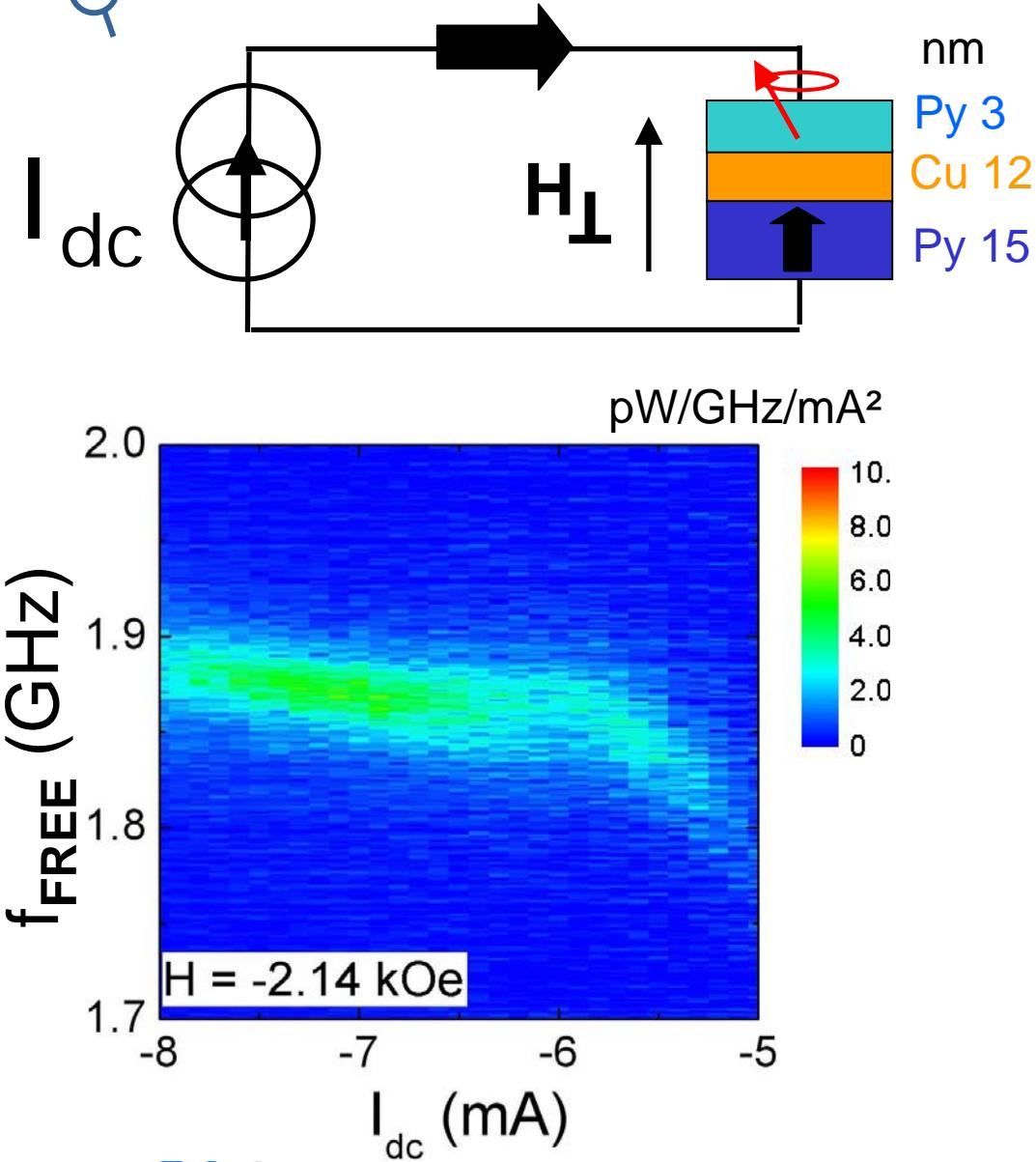
f_o $\Delta f, \Delta\varphi$

Current through THE oscillator

$$I_{tot} = I_{dc} + I_{hf\ source}$$

f_o $\Delta f, \Delta\varphi$

Features of a single STNO

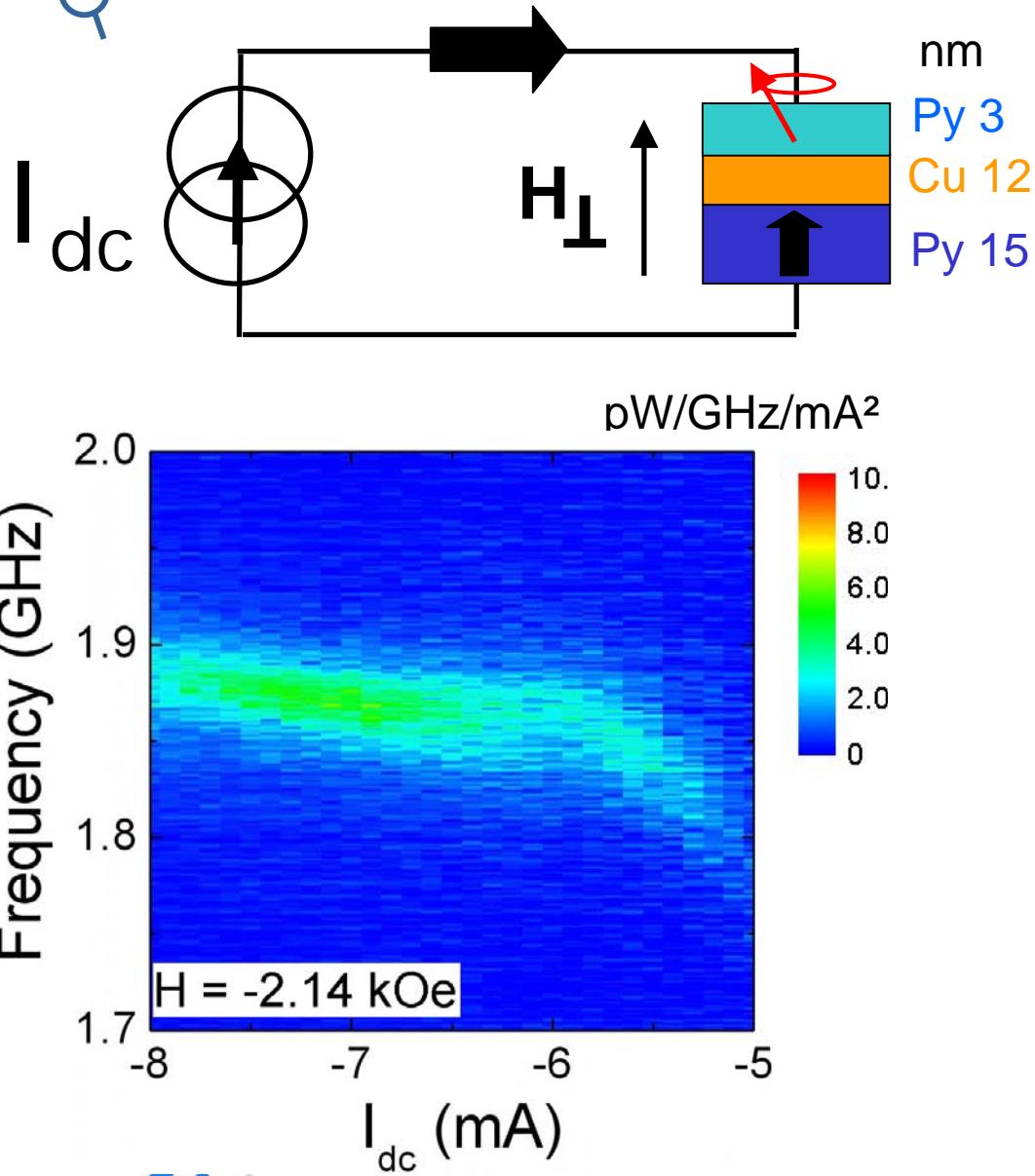


microwave characteristics of
the emission of an STNO:

- frequency
- linewidth
- power
- agility in current
(df_{FREE}/dI_{dc})

change with the experimental
conditions (H and I_{dc})

Features of a single STNO



microwave characteristics of
the emission of an STNO:

➤ frequency

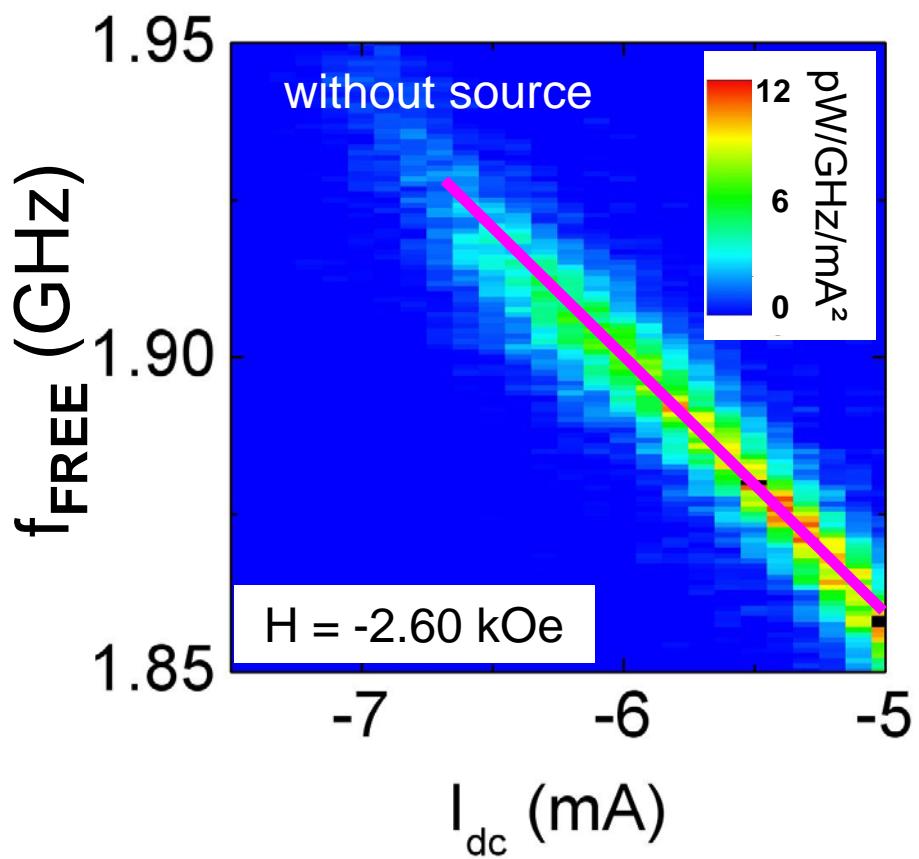
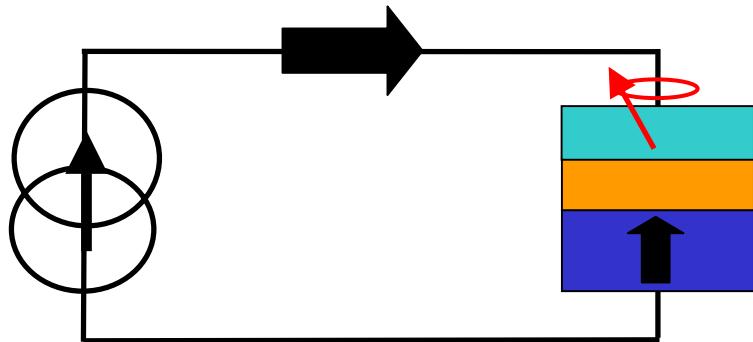
➤ linewidth

➤ power

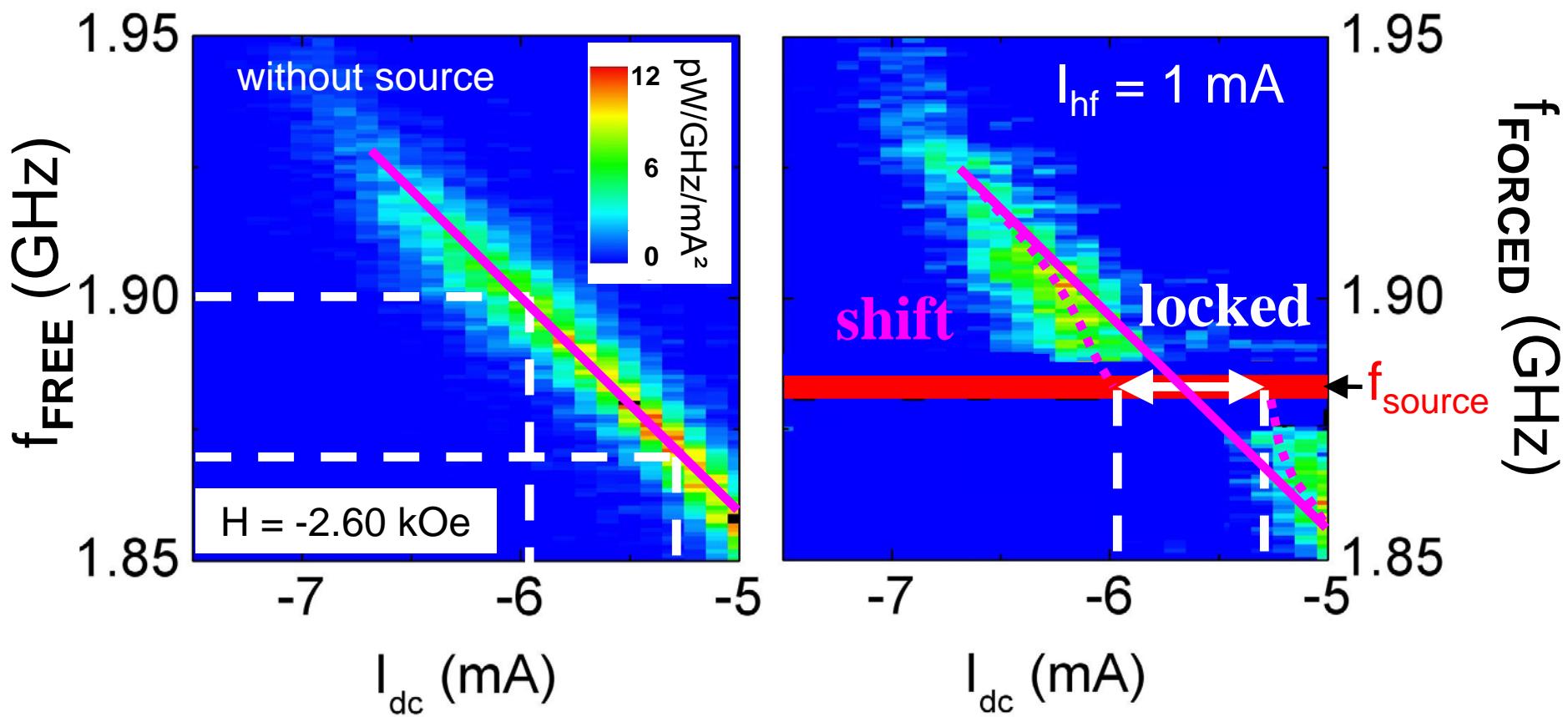
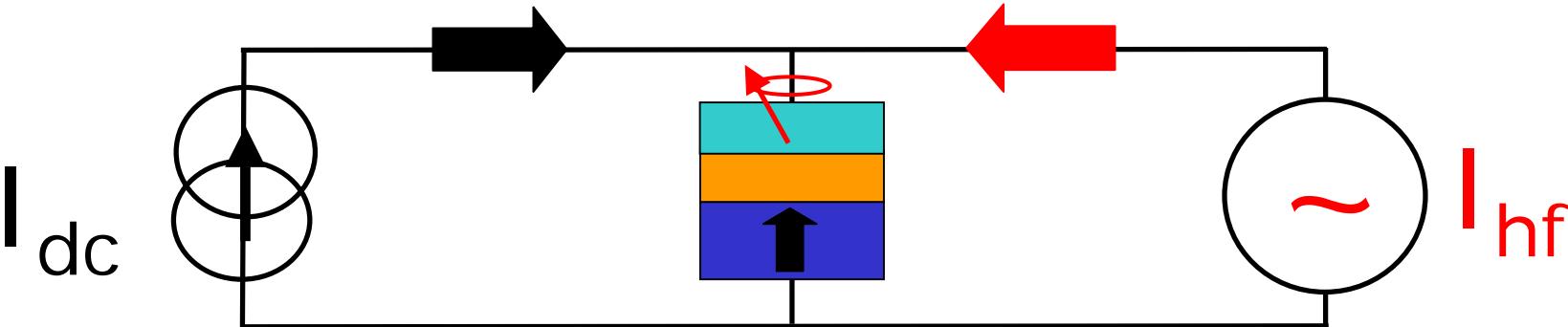
➤ agility in current
(df_{FREE}/dI_{dc})

change with the experimental
conditions (H and I_{dc})

Phase locking to an external source



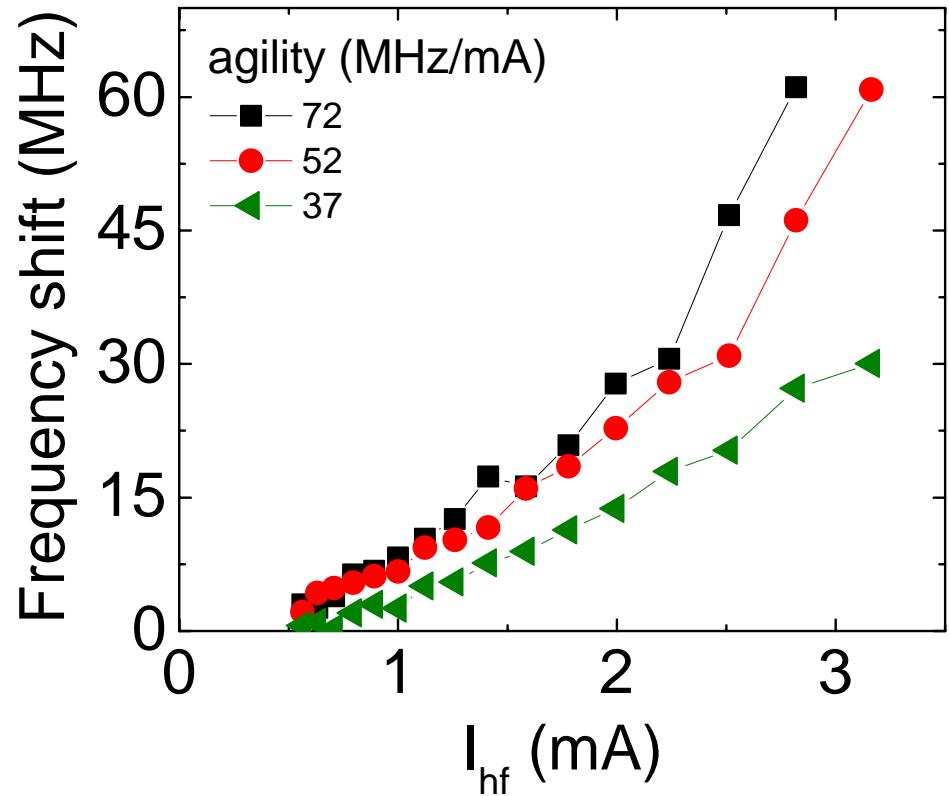
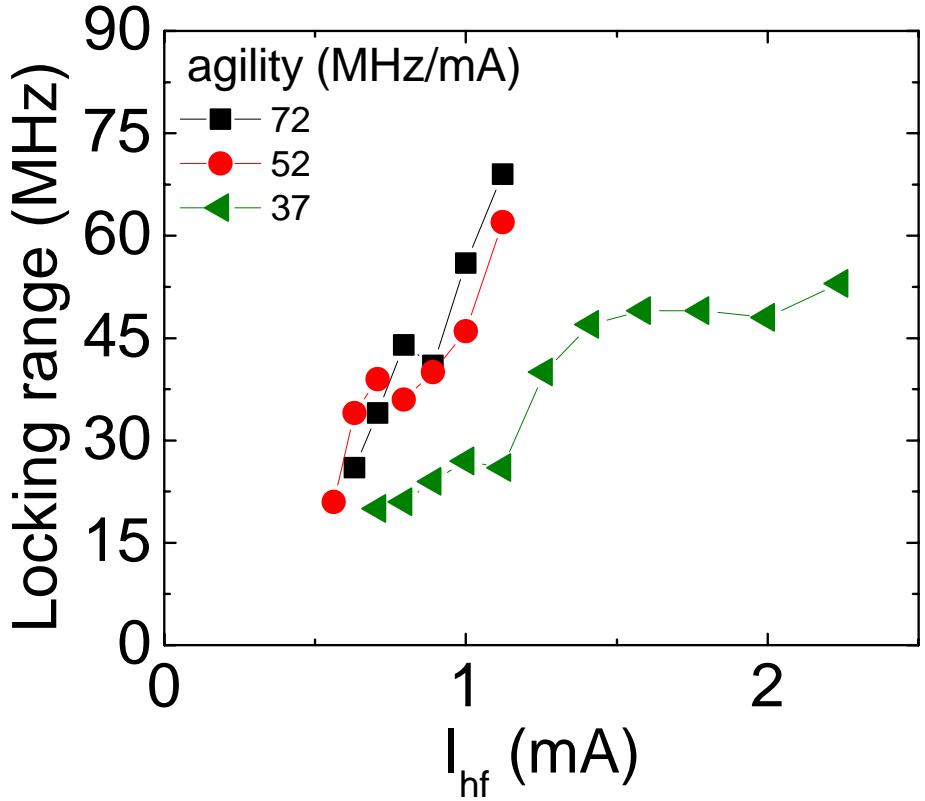
Phase locking to an external source



Locking range and frequency shift



(H, I_{dc}) chosen so that only the agility in current changes

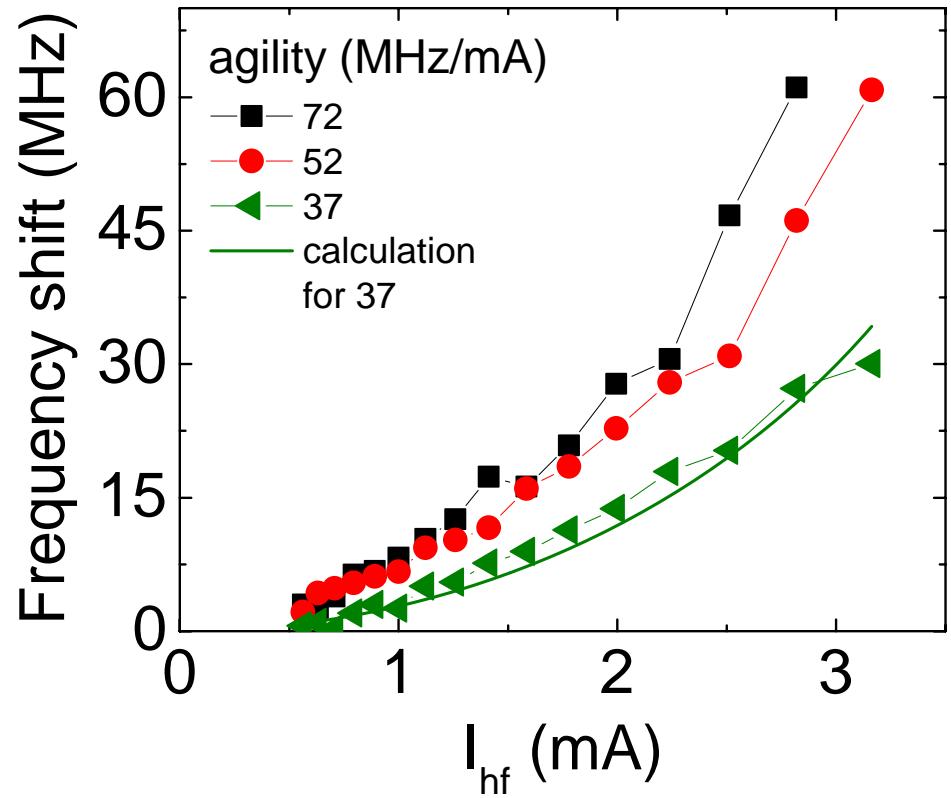
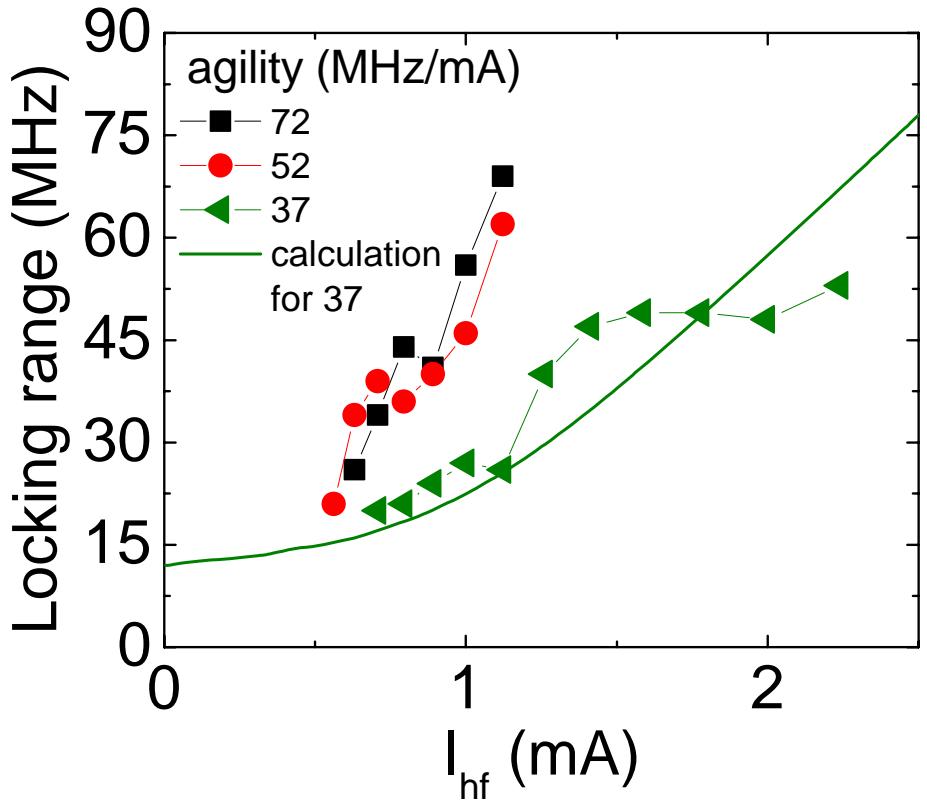


the microwave characteristics of the STNO
determine the coupling strength to the external microwave current

Locking range and frequency shift

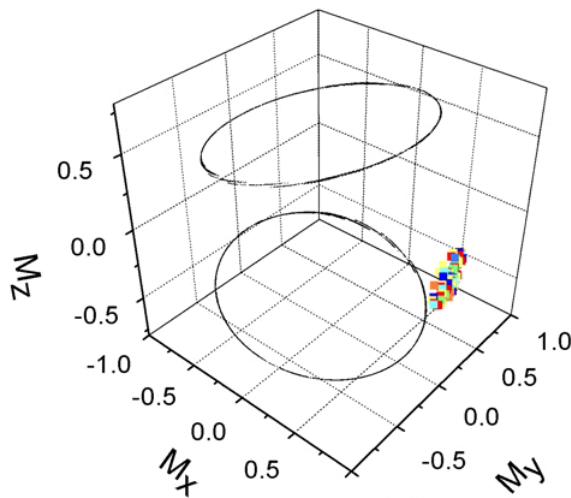


(H, I_{dc}) chosen so that only the agility in current changes

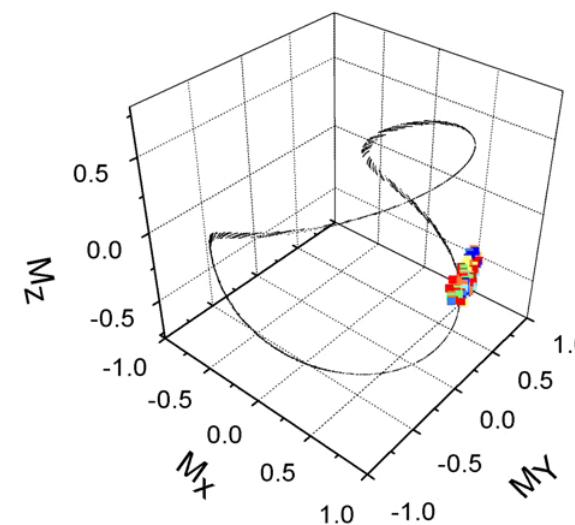
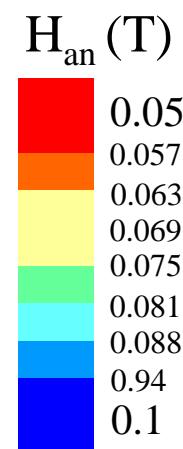


the microwave characteristics of the STNO
determine the coupling strength to the external microwave current

Synchronization / Phase dynamics (100 oscillators)

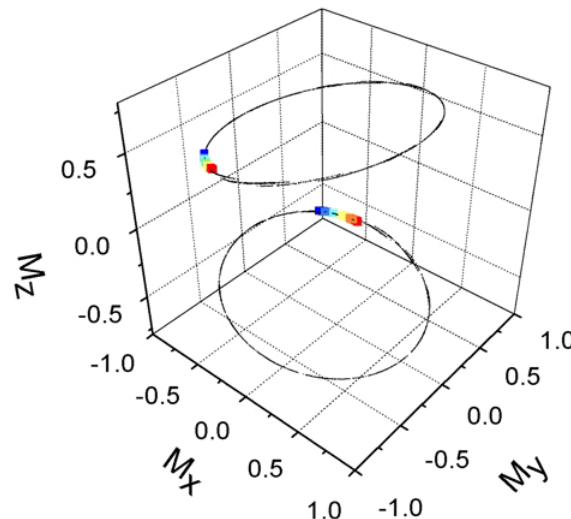


$J = 0.035 \text{ T}$
 $\tau = 5 \text{ ps}$
 $A_{GMR} = 0.03$

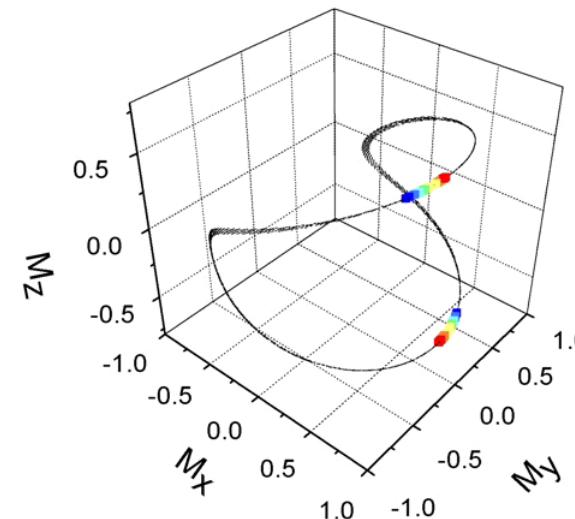


$J = 0.01 \text{ T}$
 $\tau = 300 \text{ ps}$
 $A_{GMR} = 0.4$

Macrospin simulations



**SYNC :
adjustment
of the
PHASE**



THALES



equation of the magnetization motion

$$\frac{db}{dt} = -i(\omega_{FMR} + Nb^2)b - \Gamma(1-Qb^2)b + \sigma I_{dc}(1-b^2)b$$

rotation damping spin transfer

$$b = ce^{i\varphi} \quad \text{Spin-wave : amplitude } c \text{ and phase } \varphi$$

A. Slavin *et al.* *IEEE 41*, 4 (2005)

uniformly rotating phase

A. Pikovsky, *Synchronization, A universal concept in nonlinear sciences*, Cambridge Nonlinear Science Series 12 (2001)

$$\phi = \varphi + \frac{N}{\sigma I_{dc} + \Gamma Q} \ln(c) + \phi_0$$

valid even when slightly perturbed



$$\frac{db}{dt} = -i(\omega_{FMR} + Nb^2)b - \Gamma(1 - Qb^2)b + \sigma I_{dc}(1 - b^2)b + \frac{\sigma I_{hf}}{2\sqrt{2}} \tan(\gamma) e^{-i\varpi_s t}$$

external source term

phase dynamics:

$$\Delta\phi = \phi - \phi_{source}$$

Adler equation

$$\frac{d(\Delta\phi)}{dt} = 2\pi(f_{FREE} - f_{source}) + \varepsilon \sin(\Delta\phi) + \xi(t)$$

detuning
 coupling strength
 noise ω^2

R. Adler, *IEEE*, 61, 10 (1973)

coupling strength

$$\varepsilon = \frac{I_{hf}}{2\sqrt{2}} \sigma \tan(\gamma) \sqrt{\frac{I_{dc}}{I_{dc} - I_{th}}} \sqrt{1 + \left(\frac{2\pi}{\sigma} \frac{I_{dc}}{I_{th}} \frac{\partial f_{FREE}}{\partial I_{dc}} \right)^2}$$

A. Slavin *et al.*, Phys. Rev. B, 72, 092407 (2005)



CENTRE NATIONAL
DE LA RECHERCHE
SCIENTIFIQUE



Phase locking to an external source : principle



General equation of the phase dynamics of a forced oscillator:

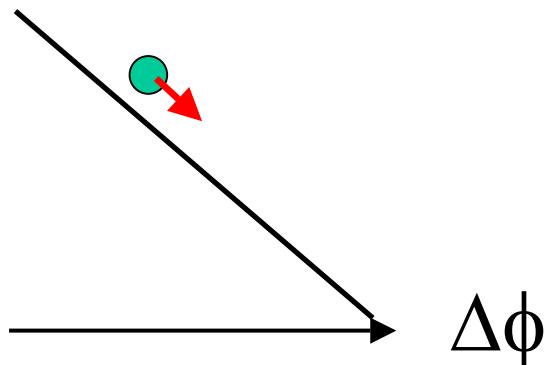
R. Adler, IEEE, 61, 10 (1973)

$$\frac{d(\Delta\Phi)}{dt} = (\omega_{free} - \omega_{source})$$

detuning

► if $\varepsilon = 0$ $\Delta\Phi = (\omega_{free} - \omega_{source}) t + \Phi_0$

the STNO and the source evolve independantly





Phase locking to an external source : principle

General equation of the phase dynamics of a forced oscillator:

R. Adler, *IEEE*, 61, 10 (1973)

$$\frac{d(\Delta\Phi)}{dt} = (\omega_{free} - \omega_{source}) + \varepsilon \sin(\Delta\Phi)$$

detuning

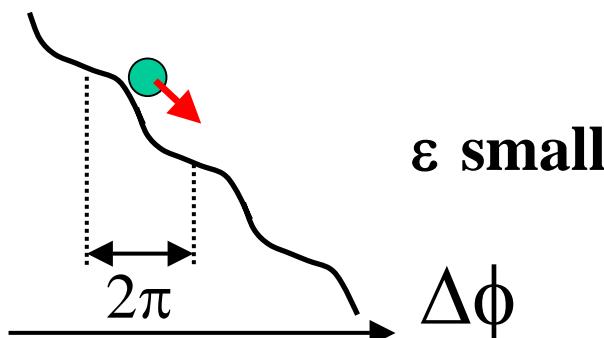
coupling strength

➤ if $\varepsilon \neq 0$

$\Delta\Phi$ is alternatively increased and decreased

Indeed : $I_{\text{tot}} = I_{\text{dc}} + I_{\text{hf}} \sin(\omega_{\text{source}} t)$

ω_{free} is alternatively increased and decreased





General equation of the phase dynamics of a **forced oscillator**:

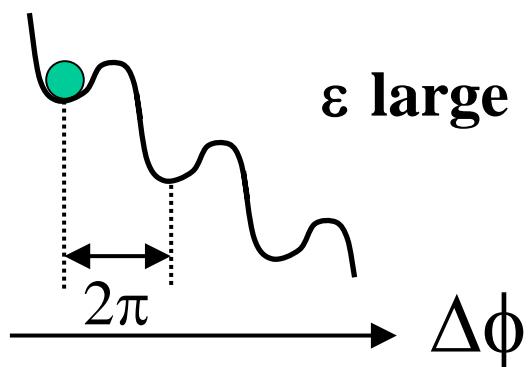
R. Adler, *IEEE*, 61, 10 (1973)

$$\frac{d(\Delta\Phi)}{dt} = (\omega_{free} - \omega_{source}) + \varepsilon \sin(\Delta\Phi)$$

detuning

coupling strength

- if $\varepsilon >$ detuning
 - there is a solution with constant $\Delta\Phi$
 - ➡ locking for large couplings





Phase locking to an external source : principle



General equation of the phase dynamics of a forced oscillator:

R. Adler, IEEE, 61, 10 (1973)

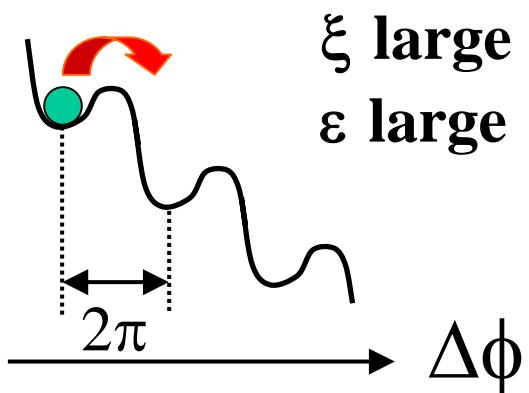
$$\frac{d(\Delta\Phi)}{dt} = (\omega_{free} - \omega_{source}) + \varepsilon \sin(\Delta\Phi) + \xi(t)$$

detuning

coupling strength

noise σ^2

- the noise ξ accounts for frequency fluctuations (linewidth)

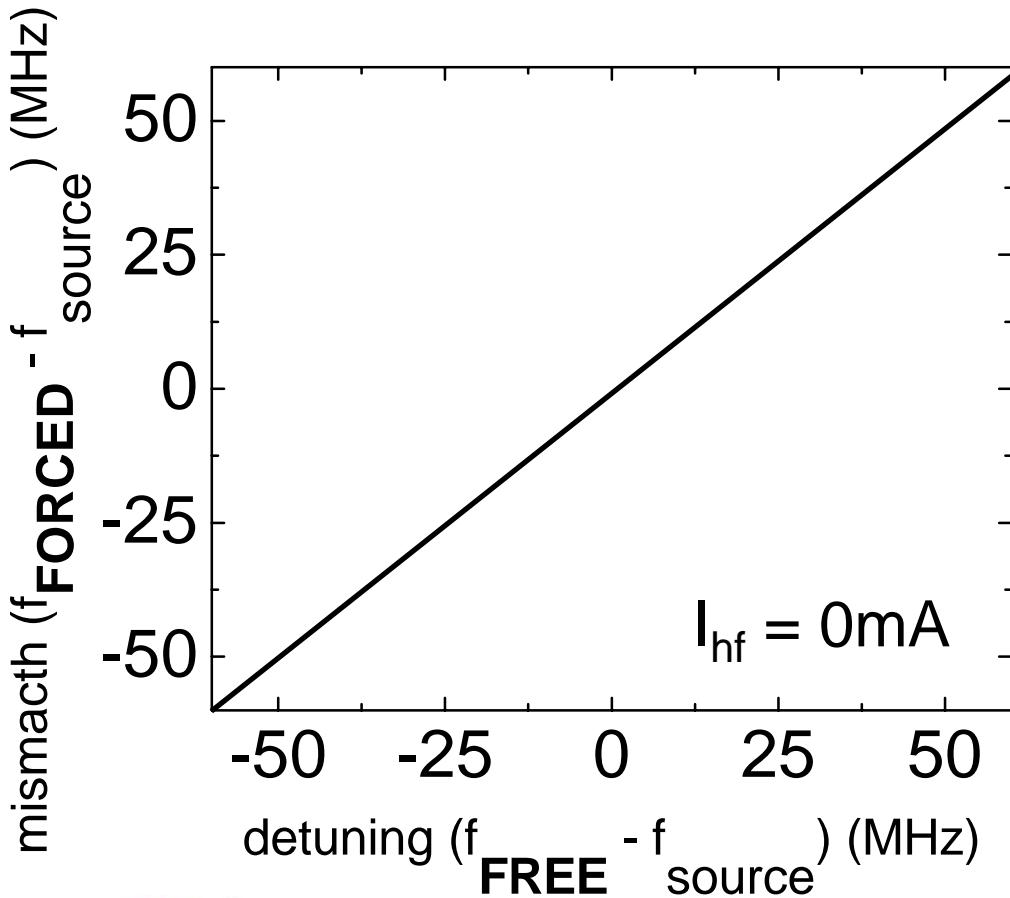


Influence of phase noise: experiment vs model



$$\text{mismatch} = f_{\text{FORCED}} - f_{\text{source}}$$

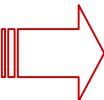
$$\text{detuning} = f_{\text{FREE}} - f_{\text{source}}$$



Influence of phase noise: experiment vs model



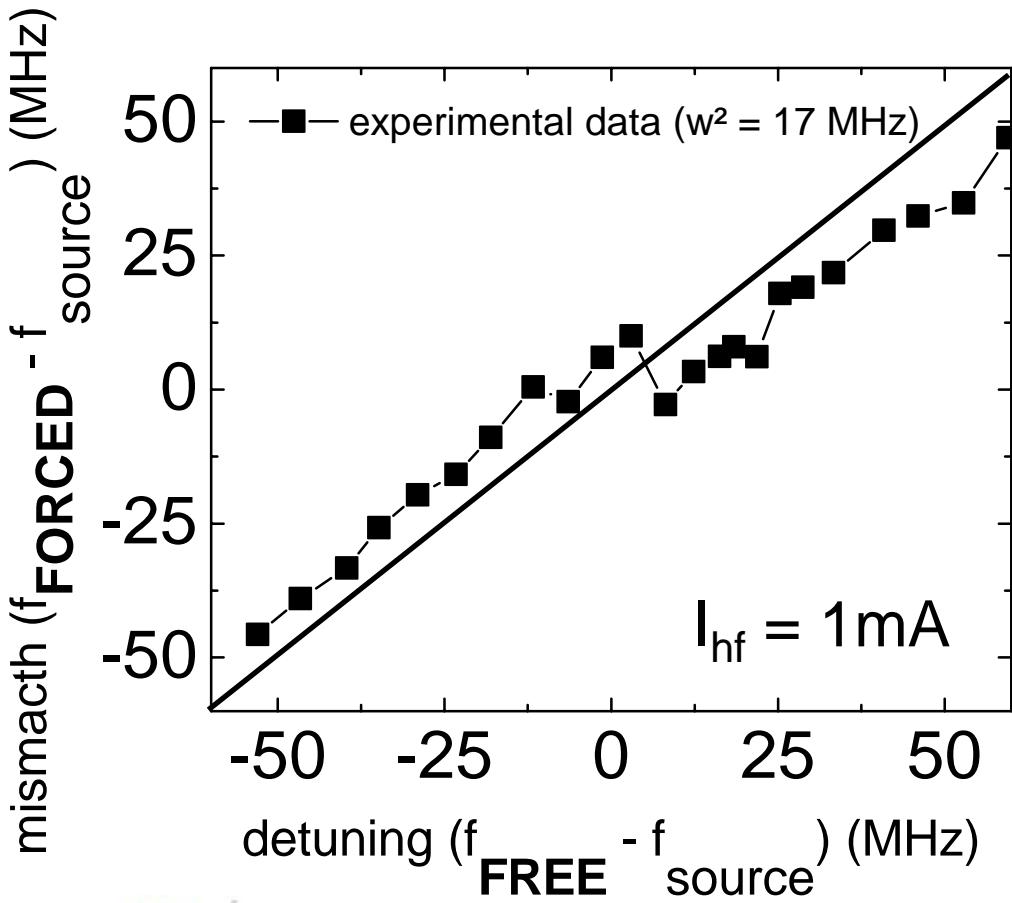
$$\frac{d(\Delta\phi)}{dt} = 2\pi(f_{FREE} - f_{source}) + \varepsilon \sin(\Delta\phi) + \xi(t)$$



calculation of the mismatch

depends on:

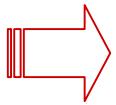
- detuning ($f_{FREE} - f_{source}$)
- noise w^2 (linewidth)
- coupling strength ε (free parameter)



Influence of phase noise: experiment vs model



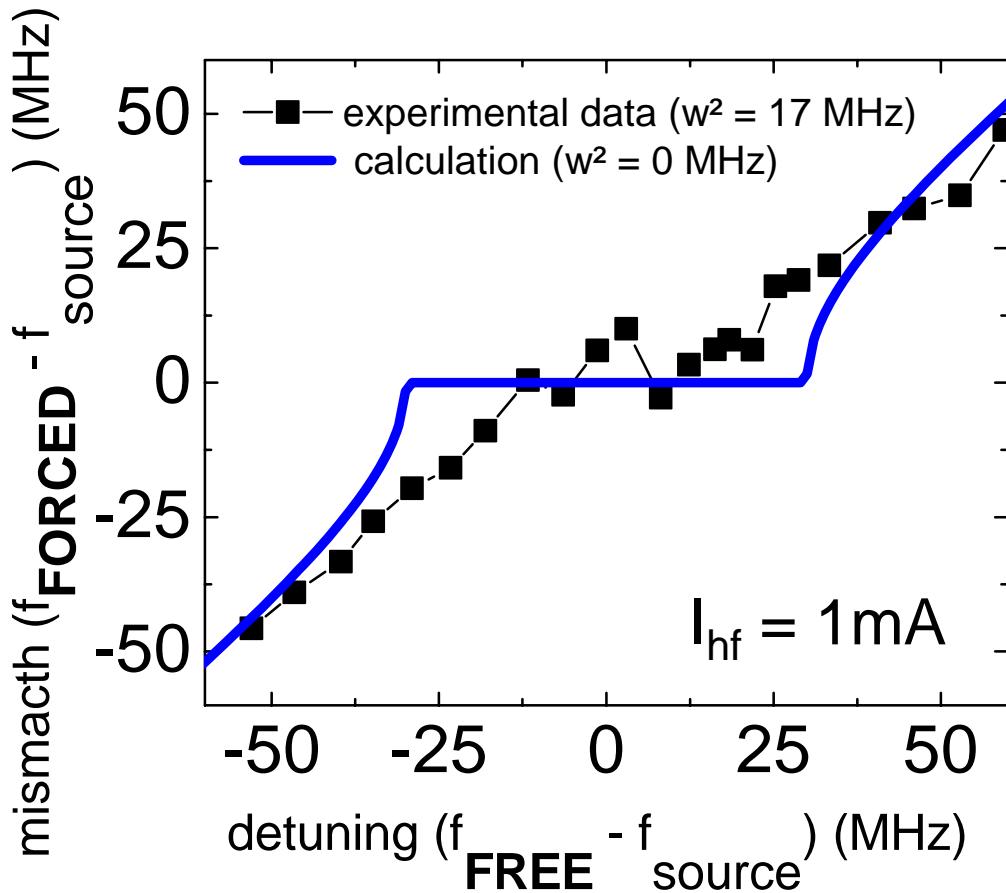
$$\frac{d(\Delta\phi)}{dt} = 2\pi(f_{FREE} - f_{source}) + \varepsilon \sin(\Delta\phi) + \xi(t)$$



calculation of the mismatch

depends on:

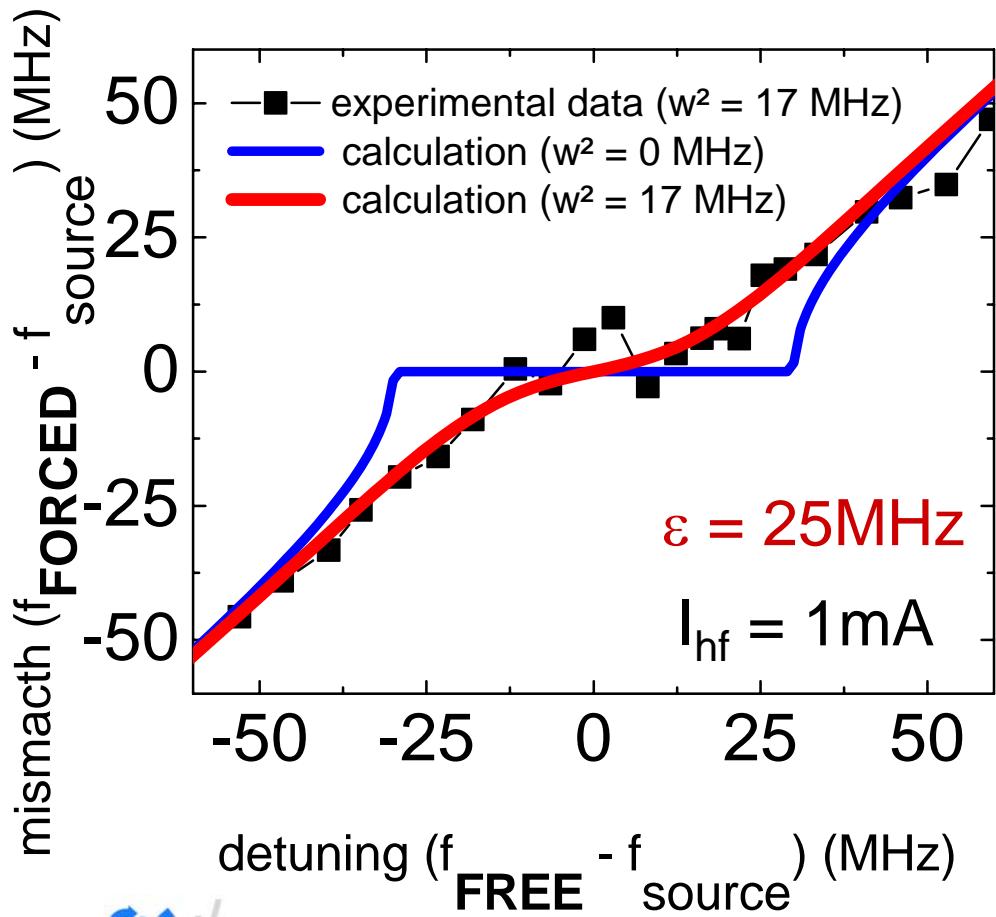
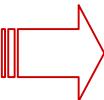
- detuning ($f_{FREE} - f_{source}$)
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Influence of phase noise: experiment vs model



$$\frac{d(\Delta\phi)}{dt} = 2\pi(f_{FREE} - f_{source}) + \varepsilon \sin(\Delta\phi) + \xi(t)$$



calculation of the mismatch

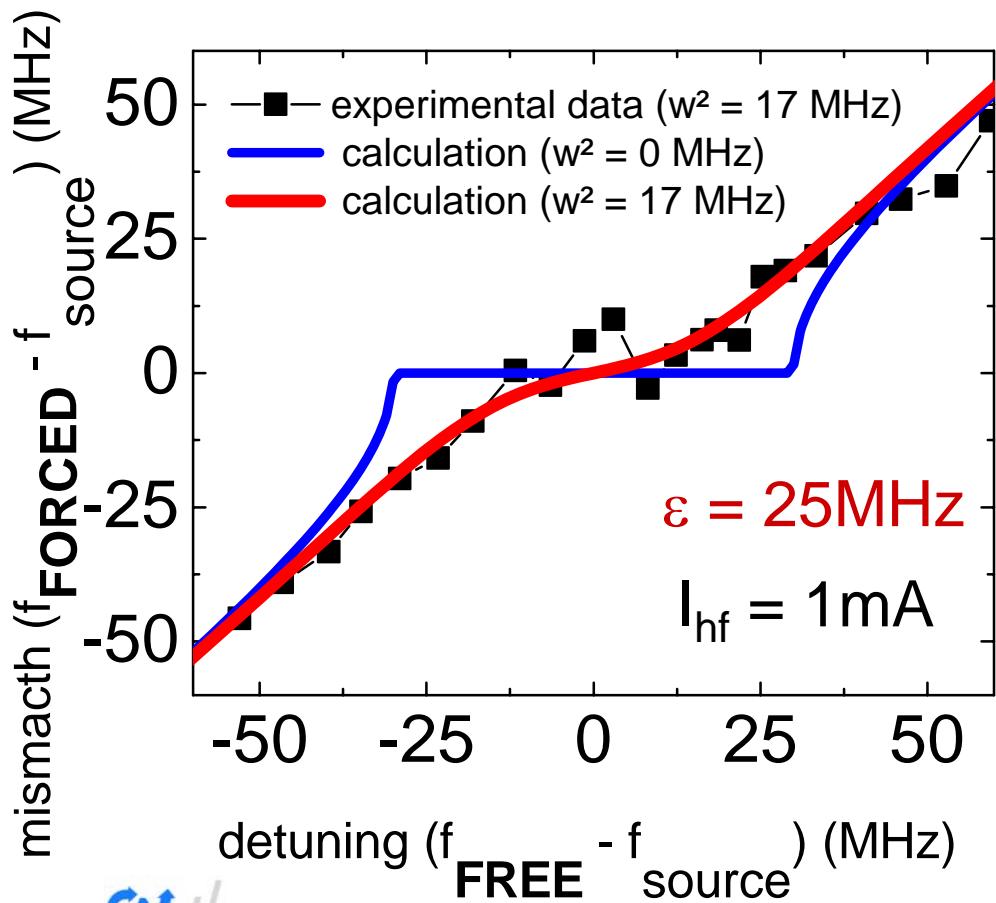
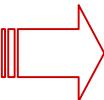
depends on:

- detuning ($f_{FREE} - f_{source}$)
- noise w^2 (linewidth)
- coupling strength ε (free parameter)

Influence of phase noise: experiment vs model



$$\frac{d(\Delta\phi)}{dt} = 2\pi(f_{FREE} - f_{source}) + \varepsilon \sin(\Delta\phi) + \xi(t)$$



calculation of the mismatch

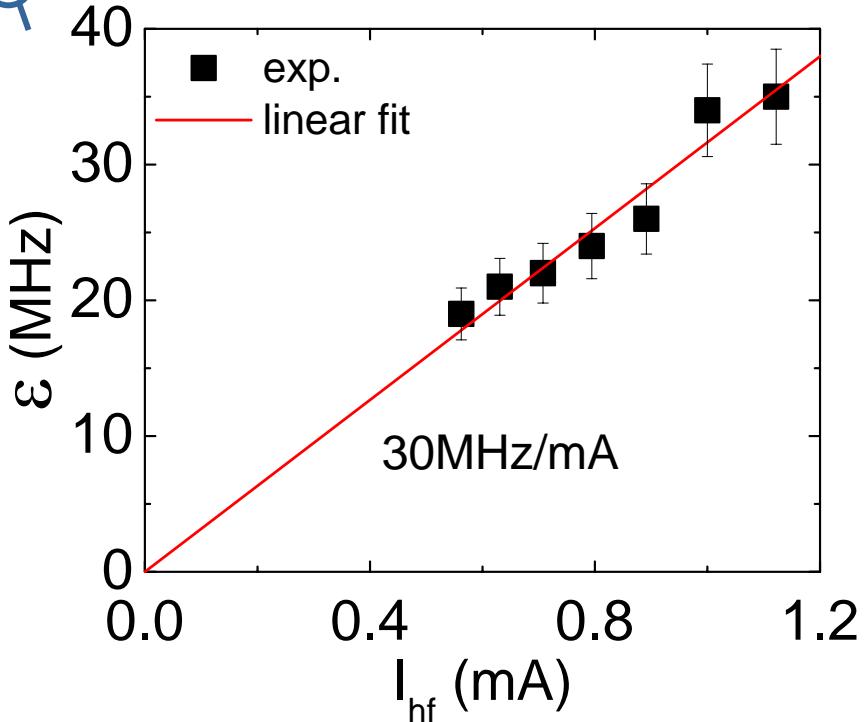
depends on:

- detuning ($f_{FREE} - f_{source}$)
- noise w^2 (linewidth)
- coupling strength ε (free parameter)

➤ forced STNOs can be described with the Adler model of phase locking

➤ critical influence of phase noise

Experimental test of the coupling calculation

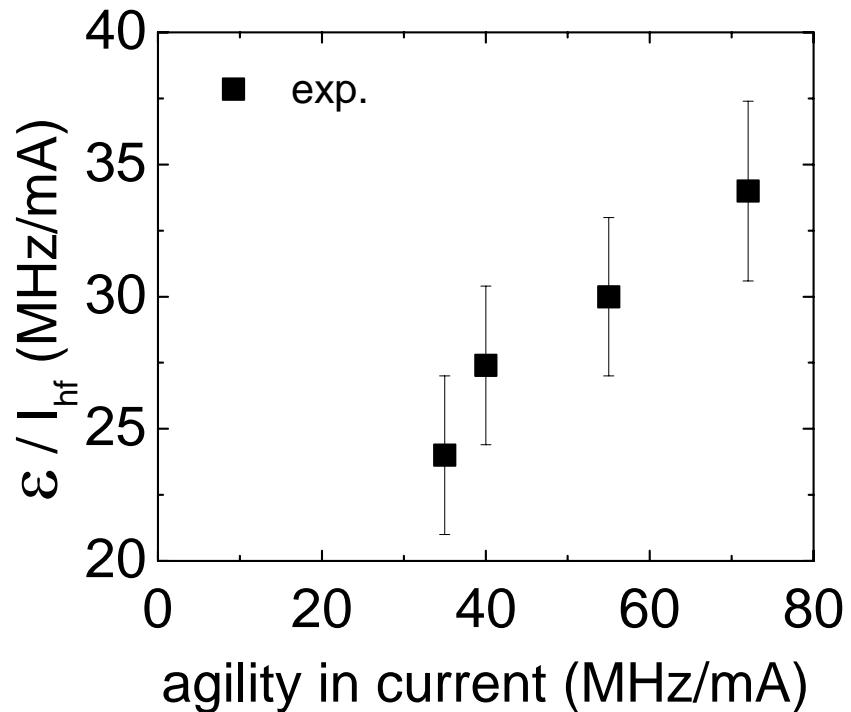
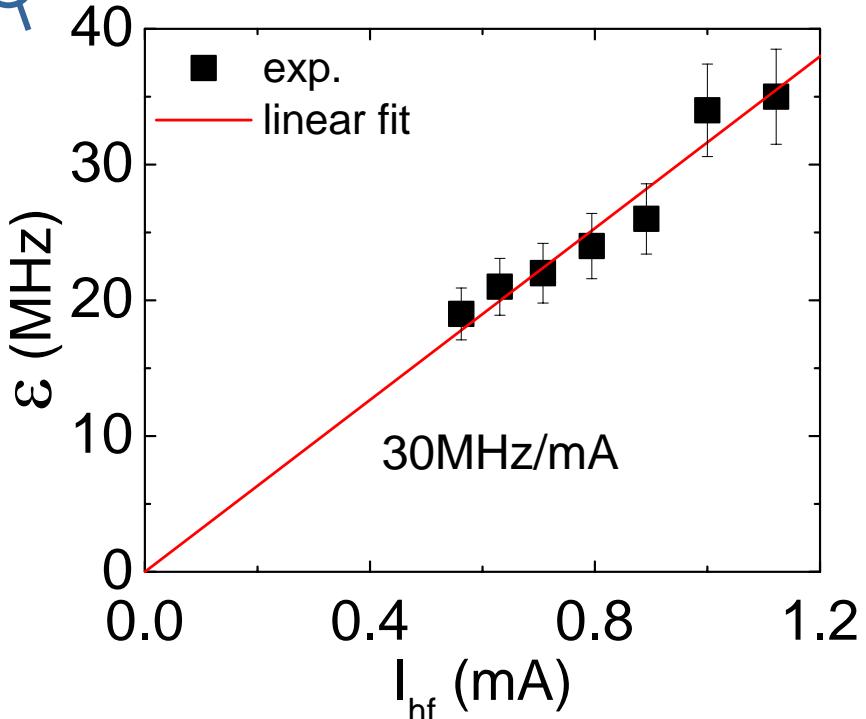


$$\epsilon = I_{hf} \frac{\sigma \tan(\gamma)}{2\sqrt{2}} \sqrt{\frac{I_{dc}}{I_{dc} - I_{th}}} \sqrt{1 + \left(\frac{2\pi}{\sigma} \frac{I_{dc}}{I_{th}} \frac{\partial f_{FREE}}{\partial I_{dc}} \right)^2}$$

varying slope

with $\sigma = 825$ MHz/mA
=> $\gamma = 2.75^\circ$

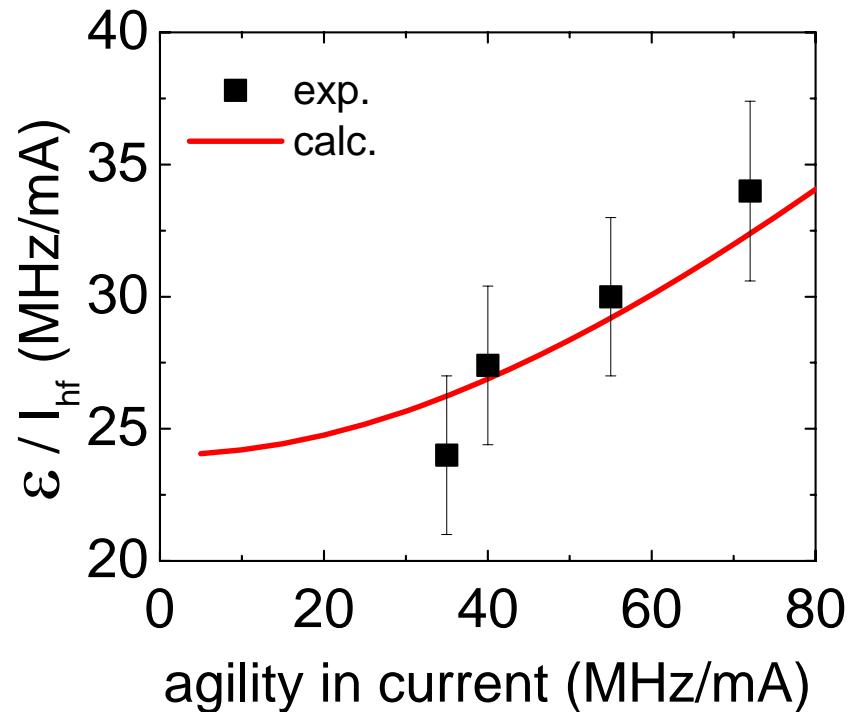
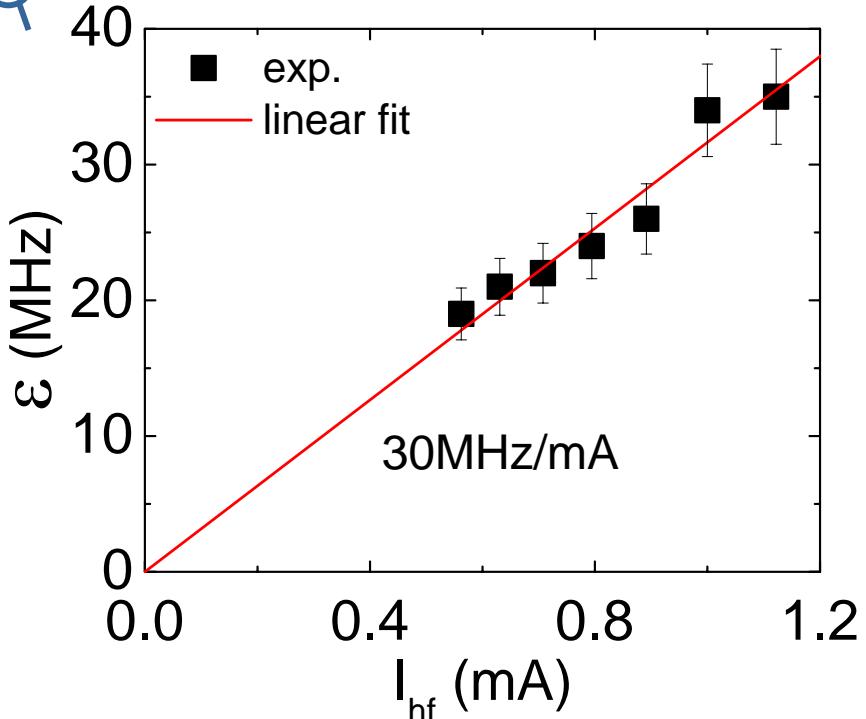
Experimental test of the coupling calculation



$$\epsilon = I_{hf} \frac{\sigma \tan(\gamma)}{2\sqrt{2}} \sqrt{\frac{I_{dc}}{I_{dc} - I_{th}}} \sqrt{1 + \left(\frac{2\pi}{\sigma} \frac{I_{dc}}{I_{th}} \frac{\partial f_{FREE}}{\partial I_{dc}} \right)^2}$$

constant varying

Experimental test of the coupling calculation



with $\sigma = 825 \text{ MHz/mA}$ ($s=0.31$)
 $\Rightarrow \gamma = 2.75^\circ$

$$\epsilon = I_{hf} \frac{\sigma \tan(\gamma)}{2\sqrt{2}} \sqrt{\frac{I_{dc}}{I_{dc} - I_{th}}} \sqrt{1 + \left(\frac{2\pi}{\sigma} \frac{I_{dc}}{I_{th}} \frac{\partial f_{FREE}}{\partial I_{dc}} \right)^2}$$

constant varying

high **agility** enhances the coupling



qualitative and quantitative understanding of injection locking experiments :

- description with **Adler model** of forced oscillators
- expression of the **coupling strength**

main trends :

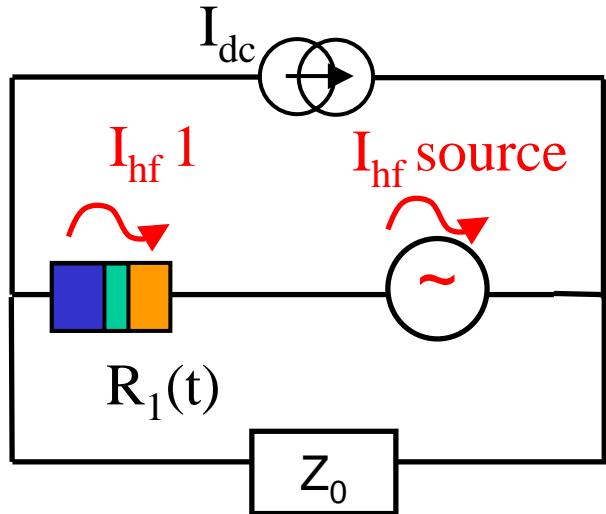
- ✓ bad influence of **the noise** on phase locking
- ✓ high **agility** enhances the coupling

B. Georges *et al.*, PRL **101**, 017201 (2008)

From Adler equation to the Kuramoto model



forced STNO

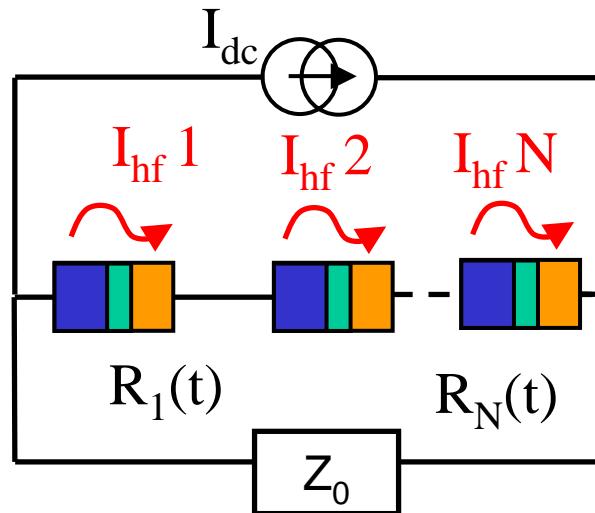


$$\frac{d(\Delta\phi)}{dt} = 2\pi(f_{FREE} - f_{source}) + \varepsilon \sin(\Delta\phi) + \xi(t)$$

Adler equation

$$\frac{\varepsilon}{I_{hf}} = \frac{\sigma \tan(\gamma)}{2\sqrt{2}} \sqrt{\frac{I_{dc}}{I_{dc} - I_{th}}} \sqrt{1 + \left(\frac{2\pi}{\sigma} \frac{I_{dc}}{I_{th}} \frac{\partial f_{FREE}}{\partial I_{dc}} \right)^2}$$

N STNOs connected in series



$$\frac{d\Delta\phi_n}{dt} = 2\pi\Delta f_n + \frac{K}{N} \sum_n \cos \Delta\phi_n + \xi_n(t)$$

Kuramoto model

$$K = \left(\frac{\varepsilon}{I_{hf}} \right) \cdot \frac{N}{Z_0 + NR} \Delta R_{osc} I_{dc}$$



Synchronization threshold $K_c = 2(w^2 + D)$

$$K = \left(\frac{\varepsilon}{I_{hf}} \right) \cdot \frac{N}{Z_0 + NR} \Delta R_{osc} I_{dc}$$

typically :
 linewidth $w^2 = 10$ MHz, dispersion $D = 100$ MHz, $I_{dc} = 5$ mA

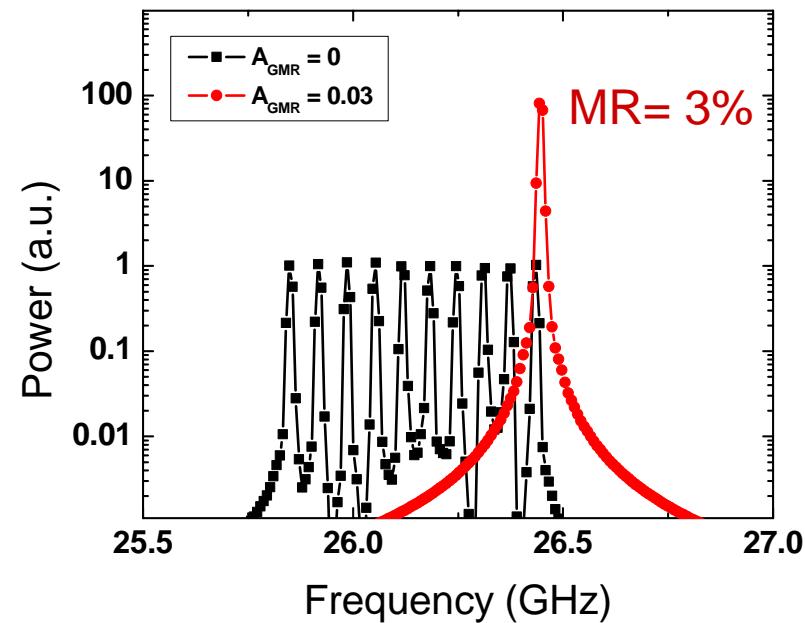
agility $df/dI = 1$ GHz/mA

condition $\frac{\Delta R_{osc}}{R} = 15\%$

agility $df/dI = 10$ GHz/mA (macrospin)

condition $\frac{\Delta R_{osc}}{R} = 1.6\%$

macrospin simulations



J. Grollier et al. PRB 73 060409 (R) 2006



N STNOs connected in series

$$P = \frac{Z_0 N^2}{(Z_0 + NR)^2} \Delta R^2 I^2_{dc}$$

Case 1 : $Z_0 = 50\Omega$

if $NR \ll Z_0$

$$P \approx \frac{N^2}{Z_0} \Delta R^2 I^2_{dc}$$

Increases as N^2

if $NR \gg Z_0$

$$P \approx \frac{Z_0}{R^2} \Delta R^2 I^2_{dc}$$

Independant of N

Case 2 : $Z_0 = 10NR$

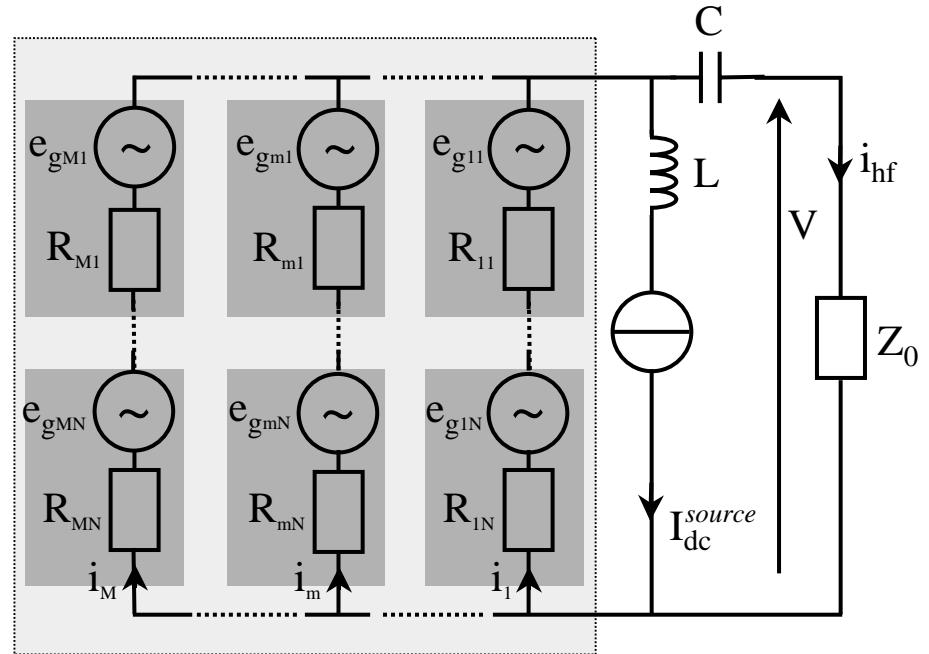
$$P \approx \frac{N}{10R} \Delta R^2 I^2_{dc}$$

Increases as N



Best solution : hybrid systems = parallel + series connection

For impedance adaptation



In all cases P increases as N



- Spin transfer oscillators : interesting features/physics
Spintronics / Non-linear dynamics
- low power : need to **synchronize** arrays
- conditions for synchronization using the results of phase locking experiments
 - B. Georges *et al.*, PRL **101**, 017201 (2008)
 - B. Georges *et al.*, APL **92**, 232504 (2008)



General equation of the phase dynamics of a forced oscillator:

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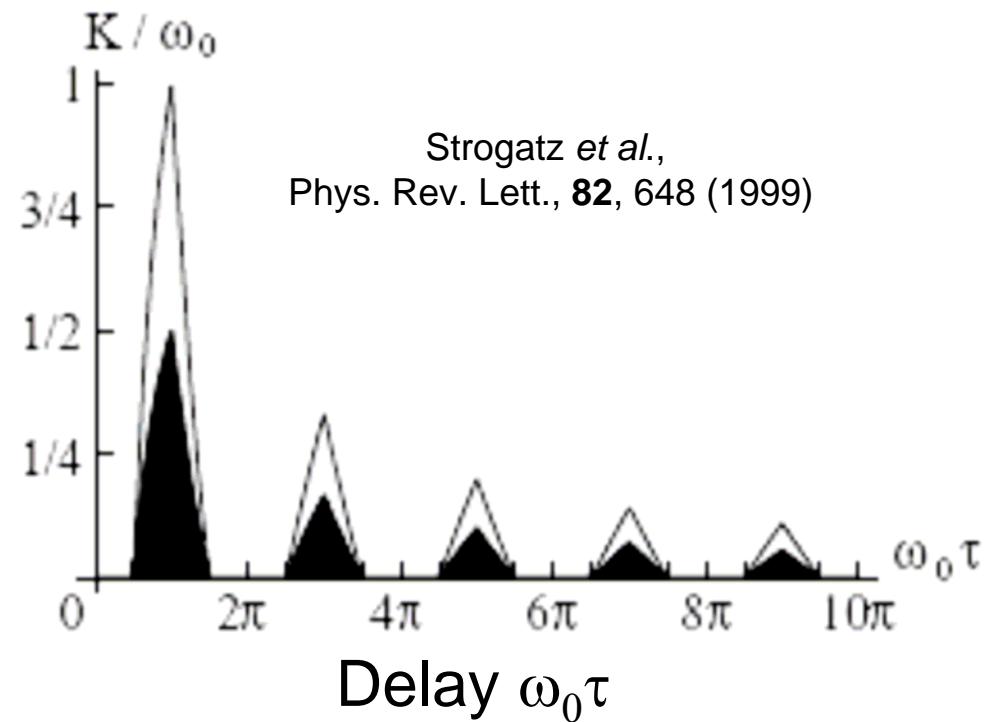
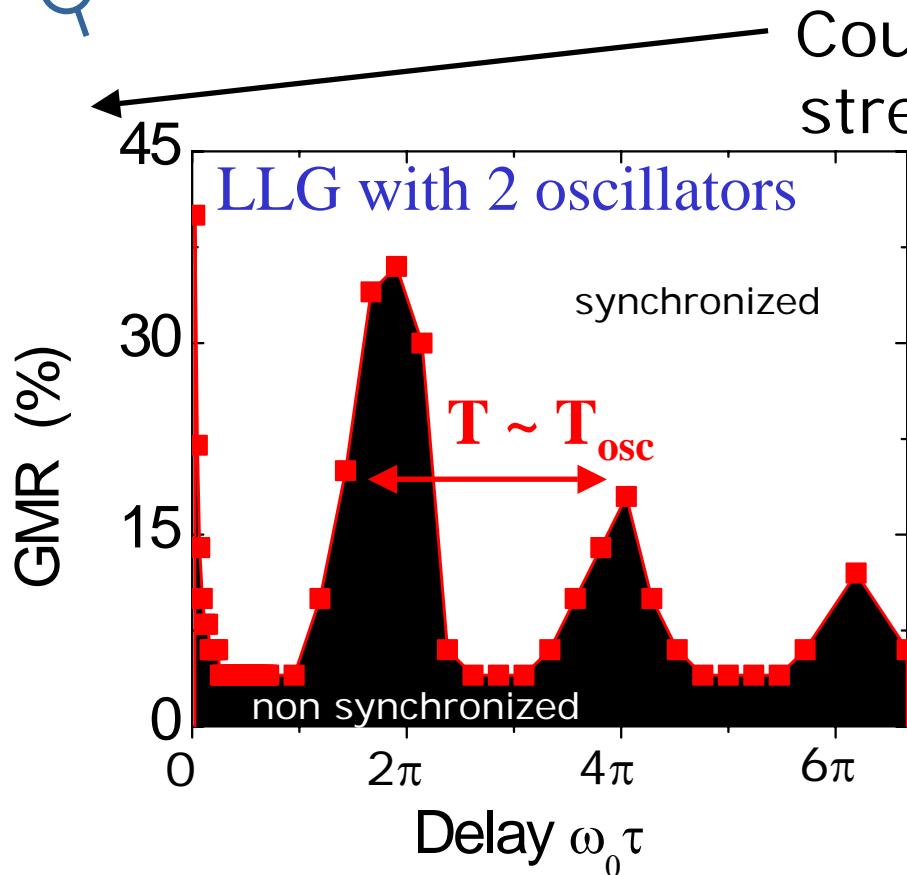
detuning

coupling strength

noise σ^2

- if $\varepsilon = 0$, $\Delta\Phi = (\omega_{\text{free}} - \omega_{\text{source}})t + \Phi_0$
 - if $\varepsilon \neq 0$, $\Delta\Phi$ is alternatively increased and decreased due to the fact that $I_{\text{tot}} = I_{\text{dc}} + I_{\text{hf}} \sin(\omega_{\text{source}}t)$
 - if $\varepsilon >$ detuning there is a solution with constant $\Delta\Phi$
 → locking
 - the noise accounts for frequency fluctuations

Influence of delay on synchronization: simulation



- Coupled STOs can be understood in the frame of **classical synchronization theory**
- **Delay** is a key parameter in the design of arrays of STOs

Phase dynamics of a forced STNO (from A. Slavin model)

$$\frac{d(\phi - \omega_{source} t)}{dt} = \omega_{FREE} - \omega_{source} + \varepsilon \cos\left(\omega_{source} t - \phi + \phi_0 + \frac{\pi}{2\sigma} \frac{\partial f_{FREE}}{\partial I_{dc}} \left(\frac{I_{dc}}{I_{th}} - 1 \right)\right)$$

with

$$\varepsilon = \frac{I_{hf}}{2\sqrt{2}} \sigma \tan(\gamma) \sqrt{\frac{I_{dc}}{I_{dc} - I_{th}}} \sqrt{1 + \left(\frac{2\pi}{\sigma} \frac{I_{dc}}{I_{th}} \frac{\partial f_0}{\partial I_{dc}} \right)^2}$$

While locked to the source:

$$\phi = \omega_{source} t + \phi_0 + \frac{\pi}{2\sigma} \frac{\partial f_{FREE}}{\partial I_{dc}} \left(\frac{I_{dc}}{I_{th}} - 1 \right) - ar \cos\left(\frac{2\pi(f_{source} - f_{FREE})}{\varepsilon} \right)$$



While locked to the source :

$$\phi = \omega_{source} t + \phi_0 + \frac{\pi}{2\sigma} \frac{\partial f_{FREE}}{\partial I_{dc}} \left(\frac{I_{dc}}{I_{th}} - 1 \right) - ar \cos \left(\frac{2\pi(f_{source} - f_{FREE})}{\varepsilon} \right)$$

If detuning = 0 :

$$\phi - \omega_{source} t = \phi_0 + \frac{\pi}{2\sigma} \frac{\partial f_{FREE}}{\partial I_{dc}} \left(\frac{I_{dc}}{I_{th}} - 1 \right) - \frac{\pi}{2}$$

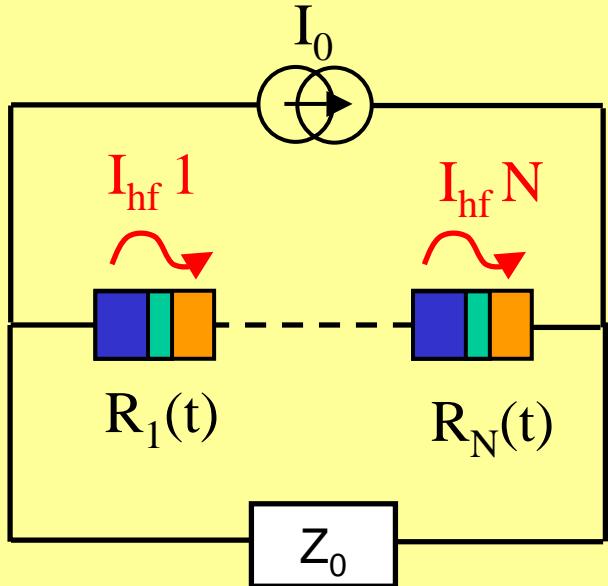
From the Adler equation to the Kuramoto model



1 STNO : Adler

$$\frac{d\Phi}{dt} = -2\pi\Delta f + \left(\frac{\varepsilon}{I_{hf}} \right) I_{hf} \cdot \cos \Phi + \xi(t)$$

N STNOs connected in series



Kuramoto model

$$\frac{d\Phi_n}{dt} = -2\pi\Delta f_n + \frac{K}{N} \sum_n \cos \Phi_n + \xi_n(t)$$

Typically :
linewidth $D = 10 \text{ MHz}$, $I_{dc} = 5 \text{ mA}$, agility $df/dI = 1\text{GHz}/\text{mA}$

+ frequency dispersion $\alpha\gamma\lambda\tau\psi$,

condition

$$\frac{\Delta R_{osc}}{R} = 15 \%$$

+ optimistic value $\alpha\gamma\lambda\tau\psi$

condition

$$\frac{\Delta R_{osc}}{R} = 2.6 \%$$

From the Adler equation to the Kuramoto model

1 STNO : Adler

$$\frac{d\Phi}{dt} = -2\pi\Delta f + \left(\frac{\varepsilon}{I_{hf}} \right) I_{hf} \cdot \cos \Phi + \xi(t)$$

N STNOs connected in series : Kuramoto model

$$i_{hf} = \frac{\Delta R_{osc} I_{dc}}{Z_0 + NR} \sum_n \cos \Phi_n$$

$$e_g i = \Delta R_{osc} I_{dc} \cos \Phi_i$$

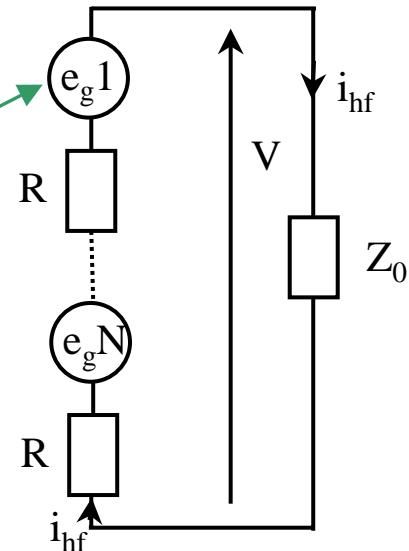
$$\frac{d\Phi_n}{dt} = -2\pi\Delta f_n + \frac{K}{N} \sum_n \cos \Phi_n + \xi_n(t)$$

Kuramoto model with :

$$K = \left(\frac{\varepsilon}{I_{hf}} \right) \cdot \frac{N}{Z_0 + NR} \Delta R_{osc} I_{dc}$$

- analytical resolution
- assumptions :
 - lorentzian frequency dispersion (γ)
 - white noise (linewidth D)

Synchronization threshold $K_c = 2(\gamma+D)$



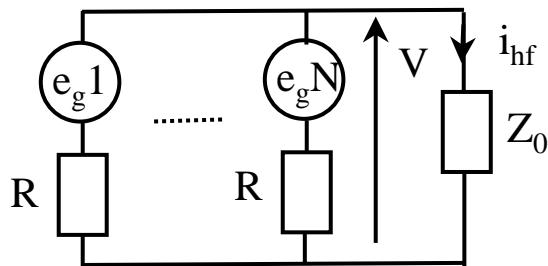


N STOs connected in series

$$P = Z_0 i_{hf}^2 = \frac{Z_0 N^2}{(Z_0 + NR)^2} \Delta R^2 I_{dc}^2$$

if all are synchronized

N STOs connected in // :



$$P = Z_0 i_{hf}^2 = \frac{Z_0 N^2}{(Z_0 N + R)^2} \Delta R^2 I_{dc}^2$$

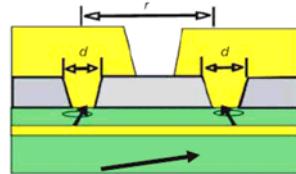
if all are synchronized

Best solution : hybrid systems // + series connection

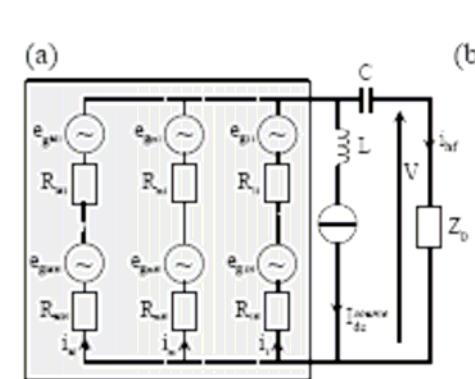
For impedance adaptation

In all cases P increases as N

NIST



Synchronization by SW : can work,
But pb with emitted power !





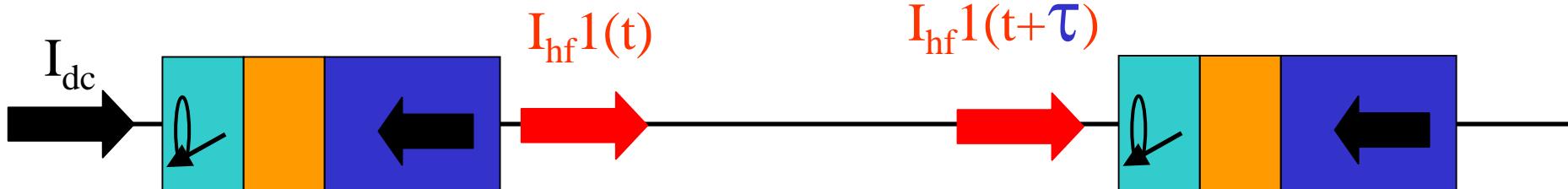
- Analytical expression for the phase dynamics of coupled STOs from LLG

$$\langle \dot{\varphi}_i \rangle = \omega_i - \frac{1}{2} \frac{\gamma_0}{1 + \alpha^2} J_{dc} \beta (GMR) \left[1 - \left(\frac{H}{H_d} \right)^2 \right] \sum_{j=1}^N \sin(\varphi_j - \varphi_i)$$

Kuramoto model

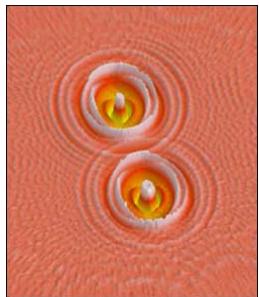
$$\dot{\varphi}_i = \omega_i + \frac{K}{N} \sum_{j=1}^N \sin(\varphi_j - \varphi_i)$$

Influence of the delay : length of the wire

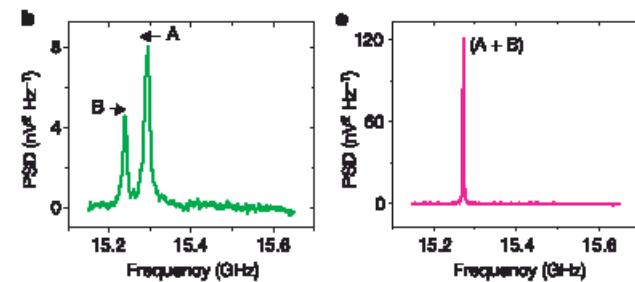
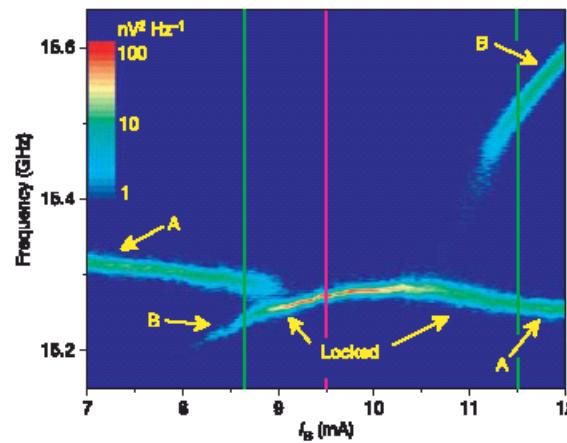
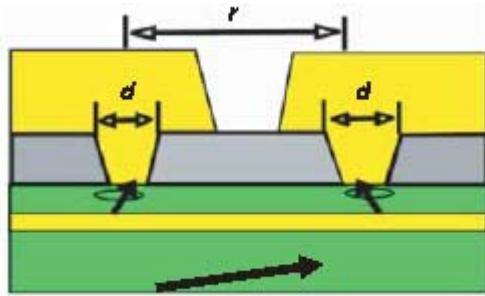




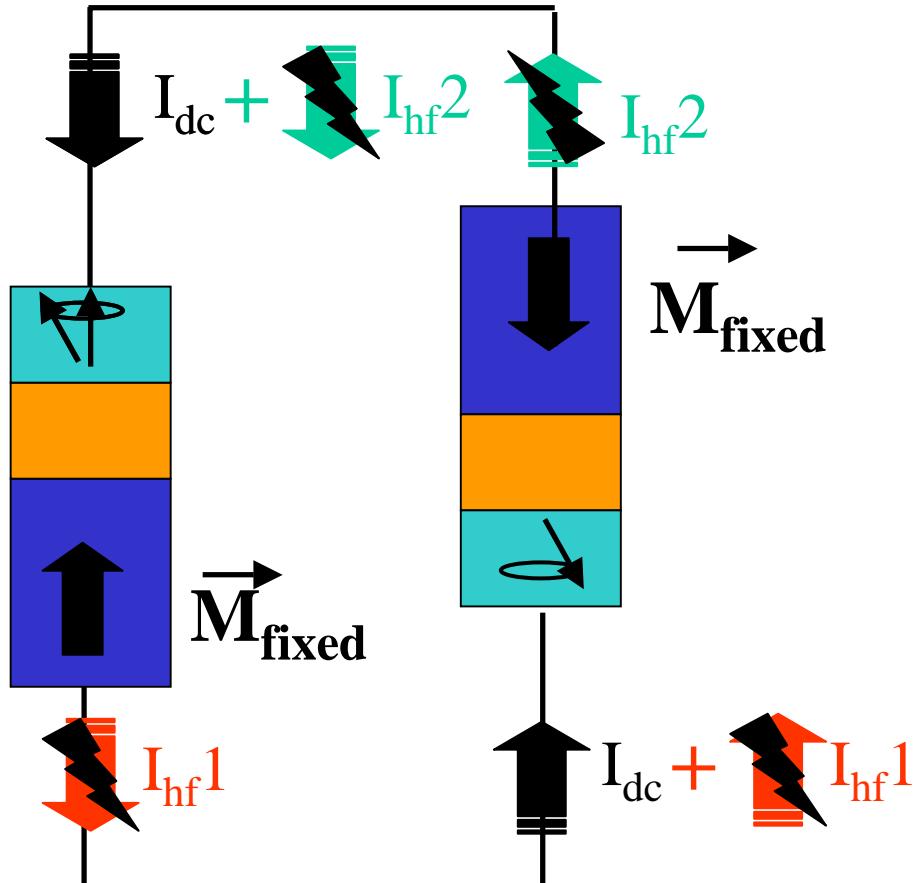
Different kinds of couplings in STOs (1)



Spin waves coupling (local)



Figures from Kaka *et al.*, Nature 2005



$I_{hf} \propto \Delta R$
coupling strength $\propto I_{hf}$

Coupling via self-emitted microwave current

Global coupling

Current through one oscillator

$$I_{tot} = I_{dc} + \sum_{i=1}^N I_{hf}(i)$$

$$f = f(I_{tot}) \quad f_0 \quad \Delta f$$

J. Grollier et al. PRB 73 060409 (R) 2006